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Reversible Actuation via Photoisomerization-Induced Melting of a Semi-Crystalline Poly(Azobenzene)

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ABSTRACT: Photoisomerization of azobenzene in polymer matrices is a powerful method to convert photon energy into mechanical work. While most previous studies have focused on incorporating azobenzene within amorphous or liquid crystalline materials, the limited extents of molecular ordering and correspondingly modest enthalpy changes upon switching in such systems has limited the achievable energy densities. In this work, we introduce a semi-crystalline main-chain poly(azobenzene) where photoisomerization is capable of reversibly triggering melting and re-crystallization under essentially isothermal conditions. These materials can be drawn into aligned fibers, yielding optically-driven two-way shape memory actuators capable of reversible bending.

Light is a powerful actuation stimulus due to the possibilities offered for specifying complex spatiotemporal responses through patterning of wavelength, mode, and polarization.^{1–3} To convert photon energy into mechanical work, photomechanical actuators are typically composed of a photoactive moiety embedded within a solid matrix, wherein a photochemical^{4,5} or photothermal⁶ effect triggers a macroscopic response. Among the most widely used are azobenzene-based systems, where molecular-scale conformational changes via *cis-trans* photoisomerization are harnessed to generate mechanical stress. Efforts in this area have largely focused on the incorporation of azobenzene within amorphous polymers^{7–9} and liquid-crystalline networks (LCNs),^{10–15} with the latter being particularly widely used due to the enhanced mechanical response afforded by the coupling of conformational changes of azobenzene to the orientational ordering of LC mesogens. While LCNs present a facile method for fabricating actuators in arbitrary geometries containing a variety of azo-based photoswitches,^{16–18} the low density of incorporated photochromes and modest enthalpy change of a liquid crystalline to isotropic transition ($\Delta H_{ni} \sim 1\text{--}10 \text{ J g}^{-1}$) limits their potential work output. Alternatively, crystallization of small molecule azobenzene derivatives can form photomechanical crystals wherein increased elastic moduli and cooperative isomerization offer the potential to dramatically enhance the photomechanical response.^{19–22} However, poor processability of such crystals limits the size and geometry of obtainable materials. Furthermore, these devices are prone to fracture due to the large internal photostresses generated during photoisomerization.²³

To advance photoactuator performance, novel material systems are needed that combine the processability of polymers with the high density of ordered chromophores found in

crystalline systems. One possible platform is semi-crystalline polymers, where assembly of high densities of polymerized chromophores into crystalline lamellae would modulate the crystalline unit cell via isomerization and generate a photoinduced stress. However, while reversible melting of semi-crystalline polymers in shape memory actuators demonstrates the potential of such systems to produce high work outputs driven by the large enthalpy change of melting ($\Delta H_m \sim 10\text{--}100 \text{ J g}^{-1}$),^{24–26} non-photothermal^{27–31} light-driven analogs have yet to be realized. Here, we demonstrate that photoisomerization of azobenzene incorporated in the backbone of semi-crystalline polymers can be harnessed to drive reversible photoactuation. While there exist reports of light-induced crystallization of supramolecular systems³² and melting of crystalline small-molecules^{33–38} and polymers^{39–41} due to photoisomerization of azobenzene, we show that this phenomenon can be harnessed for optical shape-memory materials capable of undergoing reversible deformation under cycles of illumination.

Though semi-crystalline polymers with azobenzene side-chains have been reported previously,^{41–43} we incorporate azo-

Scheme 1. Preparation of semi-crystalline poly(azobenzene)s

Figure 1. (A) ^1H NMR of equilibrated p(A6) in solution before (top) and after (bottom) 365 nm light exposure indicates efficient switching from $> 99\%$ *trans* to $\approx 97\%$ *cis*. (B) p(A6) transforms from a solid to a gooey melt following photoisomerization. (C) DSC of the *trans* polymer (grey trace) shows a melting endotherm at 100°C and a crystallization exotherm at 85°C . When isomerized to the *cis* form, a crystallization exotherm (orange trace) is observed due to thermal relaxation-driven in situ crystallization followed by a melting exotherm. The initial thermogram is recovered on the following thermal cycle (blue trace). (D) T_g switches from $\approx 34^\circ\text{C}$ to $\approx -35^\circ\text{C}$ upon photoisomerization. (E,F) The *trans* form shows crystalline ordering in WAXS and MAXS that is lost upon isomerization.

benzene into the polymer backbone to strongly couple isomerization with macromolecular conformation using the thiol-Michael addition of an azobenzene diacrylate (A6A) with 1,6-hexane dithiol (Scheme 1). Similar step-growth polymerizations have been previously employed for the synthesis of both main-chain azopolymers^{44,45} and semi-crystalline polymers,^{46–48} but to our knowledge, not one material containing both features. A diacrylate and dithiol (Scheme 1) are combined in a 1.025:1 molar ratio in chloroform with catalytic triethylamine and stirred for 24 h at room temperature. The stoichiometry is chosen to preserve acrylate end-groups for subsequent cross-linking (Figure S1). Following polymerization and purification, polymers with number average molecular weight of 20.1 kDa and dispersity of 1.65 (Figure S2) are obtained.

Following thermal equilibration for 24 h, ^1H nuclear magnetic resonance (NMR) spectroscopy (Figure 1A) of p(A6) in CDCl_3 shows that azobenzene units in p(A6) are in the *trans* form ($>99\%$). Upon illumination with 365 nm light, the *trans* population isomerizes to a photo-stationary state (PSS) composed of 97% *cis* isomers, as indicated by the shift of aromatic protons from 7.8 and 6.9 ppm (*trans*) to 6.8 and 6.7 ppm (*cis*) (Figure 1A). No additional isomerization is observed with further light exposure, confirming the PSS is reached. UV-visible (UV-vis) absorption spectra further confirm photo-switching with a dramatic decrease of the *trans* π,π^* peak and concomitant increase of the *cis* n,π^* peak centered at 360 nm and 530 nm, respectively (Figure S3). Though the slight overlap of the *trans* and *cis* absorption peaks (Figure S3) precludes quantitative switching, the significant change in isomer population

greatly affects the polymer properties. Before switching, *trans*-p(A6) is a hard yellow solid (Figure 1B); following isomerization, *cis*-p(A6) is a gooey orange melt. We employ differential scanning calorimetry (DSC) to investigate the cause of such dramatic changes in material properties. Solid-state *trans*-p(A6) films are initially semi-crystalline, with a melting endotherm upon heating at $T_m = 106^\circ\text{C}$ with $\Delta H_m = 67.0 \text{ J g}^{-1}$ and a crystallization exotherm at $T_c = 89^\circ\text{C}$ and $\Delta H_c = 60.4 \text{ J g}^{-1}$ during cooling (Figure 1C, grey). Assuming a purely crystalline sample has a typical value of $\Delta H_m \approx 150 - 300 \text{ J g}^{-1}$,⁴⁹ we estimate that *trans* p(A6) is $\sim 20 - 40\%$ crystalline. However, when the same sample is photoisomerized in solution, dried, and subjected to an identical thermal treatment, heating produces an exotherm (69.2 J g^{-1}) followed closely by a melting endotherm (64.0 J g^{-1}) at T_m of *trans*-p(A6) (Figure 1C, orange), both of comparable magnitude to the corresponding transitions of the *trans*-p(A6). We attribute the initial exotherm to in-situ crystallization of the polymer as it thermally isomerizes from *cis*-p(A6) to *trans*-p(A6). On subsequent thermal cycles, the initial thermogram of *trans*-p(A6) is recovered (Figure 1C, blue), confirming thermal isomerization of *cis*-p(A6). Furthermore, the glass transition temperature (T_g) changes from $\approx 34^\circ\text{C}$ in *trans*-p(A6) to $\approx -33^\circ\text{C}$ in *cis*-p(A6), similar to previous observations for side-chain azopolymers.^{38,50–54} Additionally, the T_g broadens upon crystallization (Figure 1D), likely due to restricted molecular relaxation near the crystalline regions.^{55,56} These thermograms suggest that *cis*-p(A6) is amorphous and crystallizes upon thermal relaxation to *trans*-p(A6) providing evidence that crystallinity is modulated via *trans*-*cis* photo-isomerization,

Figure 2. (A) UV-vis spectra of films irradiated at different temperatures. While a large fraction of *trans* isomer (black curve) switches to *cis* under 365 nm illumination at 22 °C (blue curve), a greater extent of switching is found under illumination at 40 °C (red curve), as evidenced by a decrease in the *trans* π,π^* peak at 360 nm. (B) GIWAXS data for thin films subject to different illumination conditions. Persistence of the high q peaks after photo-switching at 22 °C (blue curve) indicates that no melting occurs, while their disappearance after illumination at 40 °C indicates photomelting. Upon visible light exposure at 65 °C, recovery of the high q peaks indicates recrystallization. (C) POM images of thin films show an initially spherulitic morphology (top) that is lost upon exposure to UV light at 40 °C (middle), and significantly restored upon exposure to visible light at 65 °C (bottom).

similar to how changes in stereochemistry modulate crystallinity in some thermoplastics^{57–59} though rarely controlled using light.⁶⁰ These effects are corroborated by wide-angle (WAXS) and medium-angle X-ray scattering (MAXS) (Figure 1E,F). Prior to photoisomerization, *trans*-p(A6) shows scattering peaks at $q = 17.4 \text{ nm}^{-1}$, 12.9 nm^{-1} , and 1.6 nm^{-1} ; the highest- q peaks are presumably associated with inter-chain spacings between neighboring polymer chains while the peak at $q = 1.6 \text{ nm}^{-1}$ (corresponding to a d -spacing of 3.9 nm) likely reflects the periodicity along the chain axis. We attribute the scattering peak at $q = 0.49 \text{ nm}^{-1}$ ($d = 13 \text{ nm}$) to the characteristic spacing between crystalline lamellae. Upon isomerization, the low q peaks disappear while the high q peaks coalesce into one broad peak, indicating a crystalline-to-amorphous phase change.

We next investigate the effect of photoswitching properties on molecular ordering in solid-state thin films. UV-vis spectra of films on quartz confirm that *trans*-p(A6) (Figure 2A, black) switches to *cis*-p(A6) (Figure 2A, blue) with a PSS of $\approx 89\%$ *cis* isomer (see Supplementary Information) upon exposure to 365 nm light at 22 °C. Furthermore, *trans*-p(A6) is recovered upon exposure to 530 nm light. This process is fatigue-resistant over several cycles (Figure S5), confirming that switching is preserved in the solid state, albeit with a PSS under UV illumination that contains a higher residual *trans* content (11%) than in solution (3%). However, when samples are exposed to further UV light at 40 °C, a PSS of $\approx 96\%$ *cis* isomer is observed (Figure 2A, red), similar to the solution value. To investigate how solid-state isomerization influences crystalline ordering,

Figure 3 *Cis-trans* thermal relaxation (black stars) follow Arrhenius kinetics (red dotted line) for all temperatures measured using UV-vis spectroscopy. Isothermal crystallization measurements (blue circles) show that crystallization kinetics follow those of *cis-trans* isomerization at high temperatures but drop significantly faster as T_g of the *trans* polymer (yellow star) is approached.

we employ grazing incidence wide-angle scattering (GIWAXS; Figure 2B) and polarized optical microscopy (POM; Figure 2C) on thin films. While *trans*-p(A6) samples show crystalline ordering as evidenced by peaks at $q = 17.4 \text{ nm}^{-1}$ and 12.9 nm^{-1} , negligible change in the scattering profile is observed upon exposure to 365 nm light at 22 °C (Figure 2B, blue) despite the substantial degree of *trans-cis* isomerization observed by UV-vis. However, increased photoswitching upon UV exposure at 40 °C results in the coalescence of scattering peaks into one broad peak, indicating a loss of crystalline ordering (Figure 2B, red). While the $q = 1.6 \text{ nm}^{-1}$ peak observed in bulk samples is obscured in these scattering patterns due to the large low- q background, a broad peak in this region is observed in *trans*-p(A6) and disappears upon switching to *cis*-p(A6), consistent with photoisomerization-induced melting.

POM provides further insight into the interplay between isomer population and crystallinity. Initially, *trans*-p(A6) displays banded spherulites, an optical signature of semi-crystalline polymers (Figure 2C, top).⁶¹ After 365 nm light exposure at $T = 40 \text{ °C}$, the spherulites disappear (Figure 2C, middle), indicating photomelting. Notably, this change is due to photoisomerization and not photothermal heating, as IR images show an illumination-induced temperature change of $< 5 \text{ °C}$, and thus the sample temperature remains well below T_m (Figure S4A). However, no morphological change is observed upon UV light exposure at 22 °C (Figure S4B), consistent with the result from GIWAXS showing persistent crystalline ordering despite photoswitching. Collectively, these data suggest that at room temperature, molecular mobility within the crystalline domains is sufficiently restricted such that isomerization is suppressed and crystalline lamellae are preserved, while the looser packing of polymers in amorphous regions enables these regions to photoisomerize to a PSS similar to the solution-state value. As the temperature is increased, there is sufficient mobility for the azobenzene units within the crystalline domains to isomerize,⁶² as evidenced by an increase in the PSS obtained for the solid films,

Figure 4. (A) (Top) Fibers are formed by drawing from the melt and photo-crosslinking acrylate end-groups. (Bottom) POM images of fibers at 45° (left) and 0° (right) with respect the analyzer indicate a high degree of alignment. (B) Two-dimensional wide-angle X-ray scattering pattern and (C) corresponding plot of intensity vs azimuthal angle in the low- q and high- q regimes indicates chain alignment along the drawing axis (denoted by a white arrow).

and correspondingly for the lamellae to melt as observed using GIWAXS and POM. While birefringence does not re-appear upon exposure to 530 nm light at 40 °C, indicating a lack of crystallization despite photoswitching (Figure S5), spherulites re-form under visible light exposure at higher temperatures (Figure 2C, bottom).

To study the effect of temperature on re-crystallization, we monitor *cis-trans* thermal relaxation rates in thin films via UV-vis and re-crystallization in *cis-p*(A6) samples via isothermal DSC experiments (Figure S6). While thermal relaxation is consistent with Arrhenius kinetics across the range of temperatures studied (Figure 3, red dotted line),⁶³ crystallization kinetics (black circles) are Arrhenius at high temperatures but decrease significantly faster as the T_g of *trans-p*(A6) is approached (yellow star). This suggests that polymer viscosity becomes increasingly important with decreasing temperature. We hypothesize that re-crystallization is suppressed under visible light illumination at low temperatures because *cis-trans* isomerization increases T_g , thus kinetically trapping the amorphous phase upon vitrification. We note that although the rates of thermal *cis-trans* isomerization and recrystallization accelerate as temperature is increased, the light exposure time (≈ 1 min) needed to induce re-crystallization is much less than that needed for thermally-induced re-crystallization below 80 °C. Thus photoswitching, not thermal relaxation, is predominantly responsible for re-crystallization under illumination at sufficient temperatures. At even higher temperatures, we expect that the rate of thermal *cis-trans* relaxation will exceed the inherent rates of nucleation and growth, though the relative rates of these two phenomena can likely be tuned through modulation of the

thermal isomerization kinetics of azobenzene via ortho- or para-substitutions.⁶⁴

Photoactuators are fabricated by drawing polymer fibers from melts containing 2.5 wt% photoinitiator (Figure 4A). Crosslinking of acrylate end-groups using 405 nm light at 80 °C topologically fixes the crystal alignment and biases melting and re-crystallization along the drawing direction.^{65,66} Weak trans absorption at 405 nm and increased thermal relaxation limits photoswitching, preserving crystalline orientation during crosslinking (Figure S7). The fibers are insoluble in organic solvents with a broad T_g of ≈ 35 °C, a Young's modulus of $E = 120$ MPa (Figure S8), and are highly oriented as evidenced by the observed anisotropic birefringence using POM (Figure 4A, bottom). Additionally, transmission WAXS measurements reveal that lamellar normal are highly aligned along the drawing direction, as evidenced from the integrated intensity as a function of azimuthal angle with a FWHM in the low- and high- q regimes of 21° and 17°, respectively (Figure 4B,C). While processing conditions and crosslinking may both influence the degree of crystallinity and crystallite size,⁶⁷ these data reveal that cross-linked materials retain the crystalline characteristics of the linear polymers. Further, POM data on crosslinked films (Figure S9) confirm that light-induced melting and recrystallization still occur following crosslinking. Interestingly, aligned fibers show a four-spot scattering pattern at low q which suggests a tilting of the chain axis such as that often found in poly(olefin)s and poly(ester)s⁶⁸; work is on-going to better understand the crystal structure.

The actuation properties of these devices are studied by monitoring the change in radius of curvature upon illumination in

Figure 5. (A) The change in radius of curvature upon illumination increases with increasing temperature. (B) Schematic illustration of a cross-section of a fiber. Below 40 °C, isomerization drives softening via a change in T_g . Photomelting at higher temperatures results in enhanced actuation. (C) Shape-memory cycles for photo-melting and re-crystallization upon exposure to UV and green light at 60 °C and (D) corresponding photographs. Scale bar = 5 mm.

fibers with a diameter of $D = 50 \mu\text{m}$ (Figure 5A). Samples are held isothermally on a heated block and exposed to 385 nm light, chosen to maximize the penetration depth of illumination ($< 6 \mu\text{m}$ for *trans*-p(A6); see Supplementary Information) while preserving sufficient photoisomerization to induce melting (Figure S10). Illumination with 100 mW cm^{-2} of 385 nm light at $T = 25 \text{ °C}$ results in a change in curvature of $\Delta R^{-1} \approx 0.05 \text{ mm}^{-1}$ (Figure 5A, light grey), though no melting is observed at this temperature. We note that similar effects have been attributed to photosoftening due to an isomerization-induced reduction in T_g in several previous cases,^{69–72} but the specific mechanism at play here requires further study. However, when the fiber is exposed to 385 nm light at 40 °C, the bending dramatically increases due to photoinduced melting of semi-crystalline *trans*-p(A6) at the top surface of the fiber (Figure 5B, bottom), with a nearly three-fold increase in the radius of curvature compared to that generated from a change in T_g alone (Figure 5A, dark grey). Thermal images (Figure S11) reveal a modest temperature change significantly below that required for photothermal melting. Further, as heat transfer is expected to minimize the temperature difference maintained across such a thin fiber, photothermal heating is unlikely to be responsible for the observed bending behavior. A slight increase in the radius of curvature is observed upon UV irradiation at 55 °C, presumably due to additional entropic recovery of the polymer chains. Finally, reversible actuation is observed upon sequential illumination with

UV light (Figure 5C, purple) and green light (Figure 5C, green) for 1 min at 60 °C; this cycle can be repeated several times, demonstrating that isomerization-induced melting and recrystallization can be harnessed for shape-memory in an analogous way to thermally-actuated systems. We estimate that the work output for photoinduced melting is $W \approx 0.2 \text{ kJ m}^{-3}$ based on a photo-induced strain of $\epsilon \approx 1\%$ and a modulus of $E \approx 120 \text{ MPa}$ (see Supporting Information). This is comparable to LC-based photoactuators, which exhibit specific work values of $\sim 0.1 - 1 \text{ kJ m}^{-3}$ and where light penetration over $\sim 100\text{--}1000$ seconds typically dictates deformation kinetics.^{3,13,14,73} However, actuation in azobenzene systems is dependent on many factors, most notably non-negligible photothermal effects.^{74,75} Here, the work output is predominantly photochemically-mediated and limited by the small penetration depth relative to the fiber diameter. We expect that photomechanical performance could be substantially enhanced with uniform light penetration through-thickness, for example by the use of negative chromophores.⁷⁶

In conclusion, we have developed a facile method to synthesize semi-crystalline main-chain azopolymers and demonstrated that these materials can reversibly melt and re-crystallize upon nearly athermal photoisomerization. Assembly of these materials into aligned fibers yields reversible photoactuators capable of generating mechanical work via photomelting. This approach to light-responsive materials presents new opportunities in the field of photoaddressable devices.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website.

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Author Contributions

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Notes

Any additional relevant notes should be placed here.

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