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Manipulating the Aliovalent Magnetic Dopants in Ti(IV)-based Oxide Nanocrystals

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**Manipulating the Aliovalent Magnetic Dopants in Ti(IV)-based Oxide
Nanocrystals**

A Dissertation Presented

by

Muhammad Abdullah

Submitted to the Graduate School of the
University of Massachusetts Amherst in partial fulfillment
of the requirements for the degree of

DOCTOR OF PHILOSOPHY

September 2021

Chemistry

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Manipulating the Aliovalent Magnetic Dopants in Ti(IV)-based Oxide Nanocrystals

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Muhammad Abdullah

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Kevin R. Kittilstved, Chair

Scott Auerbach, Member

Dhandapani Venkataraman, Member

Wei Fan, Member

Ricardo Metz, Department Head
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DEDICATION

To those who inspired but may not read it – perhaps they can't!

[Translation in Urdu]

انتساب

ان کے نام جنہوں نے متاثر کیا مگر کبھی اسکو پڑھیں گے نہیں - شاید وہ پڑھ بھی نہیں سکتے !

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ABSTRACT
MANIPULATING THE ALIOVALENT MAGNETIC DOPANTS IN TI(IV)-BASED
OXIDE NANOCRYSTALS

SEPTEMBER 2021

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The intentional incorporation of impurities or dopants in semiconductors is fundamental to manipulate the properties that render them useful for spintronics, photocatalysis, and optoelectronics. One long-standing challenge in integrating the doped semiconductors in various applications is the design of materials with controlled individual dopant properties such as dopants speciation, valence state, and spin dynamics. Despite several elegant studies to circumvent these material challenges, the quest for new materials with tunable dopant properties to address the theoretical and experimental understanding continues. In this work, we combine synthetic chemistry and various spectroscopies to study a class of materials possessing both substitutional magnetic dopants, surface defects, and tunable charge carrier densities.

This thesis builds off a substantial body of work to understand the defect chemistry and dopant-carrier interactions in colloidal SrTiO₃, BaTiO₃, and TiO₂ nanocrystals (NCs). At the initial stages of this work, the formation of paramagnetic surface defects during hydrothermal synthesis of colloidal SrTiO₃ and BaTiO₃ NCs is

investigated. These defects are identified as surface adsorbed superoxide radicals formed by reducing molecular oxygen during synthesis. The critical roles of lactate ions as reducing agents, hydrazine as an oxygen scavenger, and choice of precursor are thoroughly explored. In a continuation of this work, the critical role of precursor and slow hydrolysis in dopant incorporation are discussed that has enabled us to prepare sub-10 nm colloidal NCs of BaTiO₃ and TiO₂ doped with an array of aliovalent magnetic dopants at substitutional Ti⁴⁺ sites. In a post-synthetic modification, the effect of excess electrons added by the photochemical method on dopant properties is explored. Using various spectroscopies, we presented a reversible control over Fe valence states between Fe³⁺ and Fe²⁺. This observation exists to at least in SrTiO₃, BaTiO₃, and TiO₂ NCs and contrasts with the acceleration of Cr³⁺ spin relaxation time in SrTiO₃ NCs found previously. This work demonstrates that dopant-carrier interactions are unique, which depends on the relative position of redox-level of dopants in the band structure and their Zeeman energies under applied magnetic field. This work helped us build the electronic structure of Ti(IV)-based oxide NCs where Fe^{3+/2+} redox-level located within the bandgap acts as trap center for photoexcited carriers while Cr^{3+/2+} is situated above the conduction band. The ability to refine and control these magnetic properties by simply using photons contributes to the fundamental understanding of carriers' interactions in oxide semiconductors and provides insights into the pathways for charge and spin-based applications for energy storage and quantum processing.

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CHAPTER 1

INTRODUCTION

1.1 Doped Semiconductors

Modulating the properties of semiconductors by introducing impurities or defects is a key to impart new functionalities which render them useful for various applications.¹ These impurities are typically foreign elements that are added unintentionally or deliberately during the growth of semiconducting materials. The typical example includes the silicon-based semiconductors which make the backbone of the entire electronic industry today. The intrinsic carrier concentration in silicon is relatively low which make it non-conducting material.² However, the addition of a small amount of impurity drastically changes the carrier concentration. As demonstrated in Figure 1.1, incorporation of phosphorous in silicon adds extra electrons while that boron introduces holes making silicon an n-type or p-type semiconductor, respectively. Such dopants are also known as aliovalent dopants as they have different valence states than the host site.

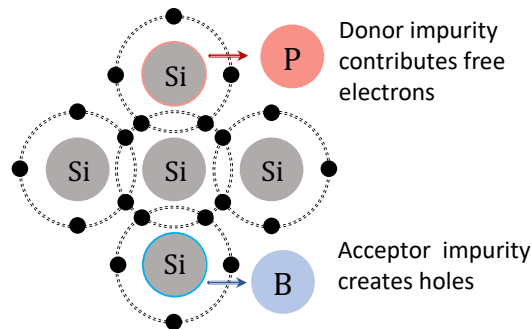


Figure 1.1 Pictorial representation of silicon doping with phosphorous (n-type dopant) and boron (p-type dopant) to add extra electrons and holes due to differences in the valence states.

This would be hard-pressed to find an electronic device that does not have silicon-based semiconductor in it. Many of the daily usage devices such as computers,

cellphones, automobiles, and all appliances in our kitchens, along with everything with a controlled circuit, has semiconductors integrated into them as a critical component. Aliovalent dopants not only change the carrier concentration and conductivity but also the optical properties of host lattices by introducing impurity levels within the bandgap. For example, Cr dopants in SrTiO_3 and TiO_2 create sub-bandgap energy levels, enhancing their visible light-harvesting properties involving these new levels.³ Similarly, Mn dopant in ZnSe and other perovskites imparts emission properties with the potential use in light-emitting diode and related display applications.^{4,5} Some of the aliovalent dopants are also known not to create the impurity levels within the bandgap yet alter the optical properties of host materials. The example includes substituting Nb^{5+} in TiO_2 and Al^{3+} dopants in ZnO-based lattices as an effective method to tune their optical response, carrier concentration and conductivity.^{6,7} Unlike many other transition metal dopants, Nb^{5+} dopant in TiO_2 does not create an impurity level within the bandgap but above the conduction band. It seems unrealistic for TiO_2 to change its optical response or carrier concentration by simply involving Nb impurity levels. Instead, the substitution of aliovalent Nb^{5+} at Ti^{4+} site in TiO_2 adds extra electrons as of charge compensation. These added electrons are typically free in the conduction band giving rise to surface plasmon in the IR region of wavelength. These n-type carriers are also explored for photocatalysis and optoelectronics.^{8,9}

Changes in the electronic structures of semiconductors are primarily defined by the individual dopant properties such as its oxidation state, energy level alignment and speciation within the host lattice. For instance, the substitution of Al^{3+} in ZnO acts

as an n-type dopant while the same dopant produces p-type carriers in SnO_2 .^{7,10} Similarly, some dopants do not change the carrier concentration or optical properties; they bring in new functionalities such as magnetism with an example of Mn^{2+} in ZnO .¹¹ Such dopants are typically isovalent as the oxidation state of dopant is the same as the host site.

Both the aliovalent dopants and isovalent magnetic dopants in semiconductors have been an integral part of technological applications. For example, in conventional silicon-based computing devices that rely on the flow of charge carriers under the influence of an external electrical field, storage is mainly accomplished on magnetic materials. In contrast, the information process is carried out in semiconducting chips containing billions of transistors.¹² As information technology evolves, these devices are expected to perform much faster in an energy-efficient way. However, there are several challenges down the road that includes energy inefficient and time-consuming exchange of information between active processing units and storage systems, the heat generated in microchips with the increasing number of transistors and quantum confinement effects with shrinking the size of technological devices. Therefore, seeking substitute electronic materials for both semiconducting devices and fundamental science is of great interest.

Alternatively, spin-based electronics, so-called “spintronics” have recently led to a new paradigm in technological advancement. Spintronics devices use spin in addition to the charge of electrons that requires far less energy to manipulate the spin states and produce less heat.¹³ The practical challenges associated with widespread applications of spintronics include *designing new materials* with tunable properties.

One of the most employed solutions to bring both semiconducting and magnetic properties in a single material is achieved by introducing paramagnetic dopants into semiconducting materials, so-called magnetic semiconductors.¹⁴ Convincing the magnetic semiconductors to exhibit room-temperature ferromagnetism (long-range spin ordering phenomenon) is crucial for their integration into devices for various applications such as spin-based information storage, creation, and transport of spin-polarized current, signal amplification and switching applications.¹⁵ Developing the field of spin-based applications, particularly quantum computing, has become a focused research area across several disciplines of chemistry, materials science and physics which aim to address materials' synthetic challenges, sensitive characterization tools to understand the behavior of magnetic centers and theoretical and experimental understanding of magnetic properties.¹⁶

1.2 ABO₃ oxides – a model host-lattice

SrTiO₃ and BaTiO₃ belong to a class of materials called perovskites with the general chemical formula of ABO₃. In the single crystal of the ABO₃ unit cell, the B-site cation is Ti⁴⁺ located at the center of the cube and coordinated to six O²⁻ anions, making a perfect octahedron. While those A-site cations, i.e., Sr²⁺ or Ba²⁺, are located at the corner of the cube as of charge compensation, as shown in Figure 1.2. The coexistence of covalent as well as ionic bonding and unfilled 3*d* shell of Ti⁴⁺ leads to unique crystal and electronic structure, which make ABO₃ a model host lattice for dopant incorporation.^{17,18}

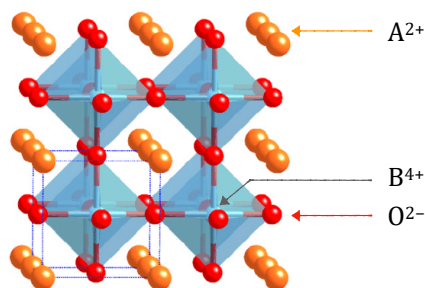


Figure 1.2 Crystal structure of ATiO_3 lattice displaying the position of atoms within the unit cells.

1.3 Applications of doped SrTiO_3 and associated challenges

SrTiO_3 has been proven to be among efficient photocatalytic materials for water splitting and CO_2 reduction owing to its proper band alignment with respect to the reaction being catalyzed.¹⁹⁻²¹ It has the conduction band located above the standard reduction potential of H^+/H_2 reaction and the valence band below the $\text{O}_2/\text{H}_2\text{O}$ reactions. This band position fulfills the fundamental requirement to qualify as a photocatalyst for complete water-splitting reaction. Despite having the appropriate band position, the practical applications of pristine SrTiO_3 have been limited due to its optical transparency in the visible region owing to its large bandgap energy (~ 3.25 eV). It means it can absorb only less than 4% of the available solar spectrum. Therefore, SrTiO_3 requires a strategy to make it absorb visible light. The introduction of cationic or anionic dopants into SrTiO_3 has been an effective strategy for tuning its bandgap and optical properties. For example, Cr, Ru and Rh doping introduce sub-bandgap states which increases the optical density of SrTiO_3 in the visible region making it an efficient light photocatalyst.²²⁻²⁶ Although these dopants have been known to enhance photocatalytic properties, they can also impart undesired properties. For instance, higher oxidation states such as Cr^{4+} and Cr^{6+} have been known to act as trap centers for photoexcited carriers in SrTiO_3 bulk powder, which

can ultimately deteriorate the photocatalytic properties.^{27,28} Similarly, aliovalent dopants and multiple oxidation states also manipulate the defect chemistry and produce various crystal disorders that can damage the intrinsic catalytic properties.²⁹ Therefore, controlled incorporation of dopants with desired oxidation state and speciation is of utmost importance for advancing the applicability of SrTiO₃ as a photocatalyst.

SrTiO₃ has also drawn attention as a multifunctional material by doping various transition metal ions.³⁰⁻³² As an example, Mn-doped SrTiO₃ has been reported to exhibit a high dielectric and magnetism simultaneously.³³ Similarly, the substitution of magnetic dopants in SrTiO₃ has led to making it dilute magnetic semiconductor (DMS), although the concrete efforts are yet in progress to achieve room temperature ferromagnetism.^{34,35} The example includes Pr-doped SrTiO₃, where room temperature magnetic ordering has been realized with increasing the dopant concentration.³⁶ The origin of many of these magnetic properties has been ambiguous and suspected as dopants clustering, the interactions between dopants and free or trapped carriers, presence of interstitial impurities and mixed-valence states of the dopants.³⁷⁻⁴⁰ Integration of SrTiO₃ based DMS into technological devices such as spintronics or any related applications demands the theoretical and experimental understanding to prepare high-quality doped SrTiO₃. Briefly, there is a need to develop synthetic methodologies that allow the dopant incorporation without clustering or losing control over the oxidation state of dopants and defect formation. There is also a need to understand the interplay between defects and dopants using sophisticated structural and spectroscopic tools.

1.4 Specific challenges addressed in the thesis

Our group has been actively studying SrTiO_3 bulk powders and colloidal NCs possessing both substitutional dopants and n-type carriers. We combine solid-state to hydrothermal syntheses and various dopant-specific spectroscopies to prepare and characterize doped SrTiO_3 . There has been a dramatic variance in their properties, such as dopant valence state, speciation, and defect chemistries across the length scale that remain elusive. For example, the substitution of aliovalent Cr^{3+} dopants at Ti^{4+} site in SrTiO_3 is compensated by the formation of oxygen vacancies (V_o) in bulk powder, while there are no spectroscopically detectable V_o in colloidal NCs.^{41,42} Instead, there is the formation of oxygen-related defects on the surface of colloidal NCs which are very sensitive to the presence of additives during the synthesis. Similarly, there is also the existence of higher oxidation states of dopants such as Cr^{4+} and Cr^{5+} in bulk powder which can be reduced to Cr^{3+} by introducing n-type carriers.²⁷ This observation has also been consistent with Mn dopants which also exist in multiple oxidation states and can be reduced to Mn^{2+} by annealing the powder samples.⁴³ Conversely, there is no evidence of multiple oxidation states of dopants in colloidal NCs. The introduction of n-type carriers in colloidal NCs instead shows the spin interactions with substitutional Cr^{3+} dopants.⁴⁴ The underlying mechanism of such spin interactions, the existence of such phenomenon in other dopants and the effect of host lattices and site symmetries are many of the unanswered questions attempted to address in the current thesis.

In chapter 2, we investigate the formation of paramagnetic surface defects during hydrothermal synthesis of ATiO_3 ($A = \text{Sr}$ and Ba) colloidal NCs by combining the

synthetic chemistry and a set of spectroscopic tools. We identified these defects as surface adsorbed superoxide radicals formed by reducing molecular oxygen during synthesis. The critical roles of lactate ions as reducing agents and hydrazine as oxygen scavengers have also been thoroughly investigated to establish a correlation between the formation and the control of these surface defects. Chapter 3 discusses the syntheses and photodoping of Fe^{3+} -doped ATiO_3 colloidal NCs. This work provides a detailed investigation of Fe^{3+} dopants speciation in as-prepared NCs, and reversible control over its valence state ($\text{Fe}^{3+}/\text{Fe}^{2+}$) by post-synthetic manipulation of Fermi level. We also propose the mechanistic details of this interplay between excess electrons and dopants in electronic structure. In chapter 4, a synthetic method to make aliovalent doped TiO_2 colloidal NCs is developed. This method allows the incorporation of an array of dopants such as Fe^{3+} , Nb^{5+} , and Cr^{3+} at the substitutional Ti^{4+} sites in the TiO_2 NCs. Following the synthesis, the effect of excess carriers introduced by (a) photodoping and (b) aliovalent Nb^{5+} doping on Fe^{3+} is explored. This works also opens the avenue for the long-standing challenge of spectroscopic investigation of “air-sensitive” colloidal NCs. Chapter 5 builds off the spin cross-relaxation found previously in SrTiO_3 NCs by our group to explore the existence of this phenomenon in structurally similar lattices. A synthetic method for making Cr^{3+} -doped BaTiO_3 and Cr^{3+} -doped TiO_2 colloidal NCs has been developed. In this method, we employed slow-hydrolyzing Ti(IV) precursor that allows the successful dopant incorporation at the Ti^{4+} site in BaTiO_3 during hydrothermal synthesis. In the second half of this chapter, the introduction of extra charge carriers using photodoping method has been discussed. Our spectroscopic studies showed that these

photochemically introduced electrons are localized at Ti^{3+} defects creating small polarons. The coexistence of both paramagnetic Cr^{3+} and Ti^{3+} in a single lattice shows the spin cross-relaxation phenomenon at room temperature, where fast-relaxing Ti^{3+} accelerates the Cr^{3+} spin relaxation rate.

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CHAPTER 2

ON THE FORMATION OF SUPEROXIDE RADICALS ON COLLOIDAL STRONTIUM AND BARIUM TITANATE NANOCRYSTAL SURFACES

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2.1 Introduction

SrTiO₃ and BaTiO₃ are wide-gap semiconductors that have received much attention due to their rich defect chemistries that enable a wide range of electrical properties from insulating to metallic.¹ In addition to exhibiting prototypical quantum paraelectric behavior,² these semiconductors are also promising host lattices for visible-light photocatalytic H₂O splitting and CO₂ reduction,³⁻⁹ sensors,¹⁰ and memristors.^{11,12} Recent advances in the colloidal synthesis¹³⁻²⁰ of ternary metal oxide nanocrystals (NCs), including SrTiO₃ and BaTiO₃, have generated further interest in emergent phenomena arise in these technologically relevant materials when prepared with nanometer dimensions.

Most of the interesting applications of bulk SrTiO₃ result from the presence of two native n-type defects: oxygen vacancies (V_O) and self-trapped Ti³⁺ ions. These defects increase the carrier density and mid-gap trapped states which give rise to tunable blue light emission²¹ but also greatly influence light absorption in the visible region that has been shown to enhance photocatalytic performance in bulk powders.²² For example, superoxide radicals ($O_2^{\bullet-}$) present on the surface of Ti(IV)-based metal oxides play an essential role in photocatalysis and degradation of organic pollutants.²³⁻²⁵ These $O_2^{\bullet-}$ are generally produced by post-synthetic treatments of

reduced titanium oxide-based materials with molecular oxygen (O_2) or H_2O_2 . Yu *et al.* generated surface $O_2^{\bullet-}$ by injecting electrons from photoinduced excited states into adsorbed O_2 on the surface of colloidal TiO_2 nanoparticles.²⁶ Another potential use of $O_2^{\bullet-}$ is the catalytical decomposition of H_2O_2 on the surface of titanium oxides²⁷ and oxidative coupling of alkyl radicals.²⁸ In most of these cases, the metal oxide surface needs to be activated prior to $O_2^{\bullet-}$ formation by irradiation or suitable chemical treatments.

We recently reported the synthesis²⁹ of colloidal $SrTiO_3$ NCs by modified hydrothermal methods^{15,17,18} that are readily suspended in non-polar solvents. Electron paramagnetic resonance (EPR) and optical spectroscopies showed that the $SrTiO_3$ NCs can display spectroscopic signatures that were assigned to an oxygen-related surface defect that is sensitive to the presence of hydrazine hydrate during synthesis. However, the role of hydrazine in the formation of this surface defect remained speculative. Hydrazine has been previously used as an additive during the hydrothermal synthesis of $SrTiO_3$ and $BaTiO_3$ NCs. Sun and co-workers have studied the absorption of hydrazine on the surface of $BaTiO_3$ nanoparticles.³⁰ Experimental and simulation analysis indicated that only amorphous $BaTiO_3$ particles are obtained without the addition of hydrazine. Similarly, Fujinami *et al.* claimed that $SrTiO_3$ NCs without cubic shape were obtained when synthesis was carried out in the absence of hydrazine.¹⁷ However, our recent report presented evidence that near-cubic colloidal $SrTiO_3$ NCs can be prepared with or without the addition of hydrazine.²⁹

Herein we report a systemic investigation of this correlation between hydrazine and the presence of this surface defect in colloidal $SrTiO_3$ NCs. We hypothesize that

the hydrazine may scavenge molecular O_2 inside the closed reaction vessel during synthesis, thus providing an inert atmosphere that would inhibit surface $O_2^{\bullet-}$ formation on the NC surface. To test the role of hydrazine, we performed the hydrothermal synthesis both with and without hydrazine prepared in ambient conditions, and also without hydrazine but under anaerobic conditions after purging the reaction solution with argon gas. Reactions performed with hydrazine or under anaerobic conditions produced the same observations and strongly suggest that hydrazine is indeed removing O_2 inside the closed reaction vessel. These experiments also strengthen our assignment of the surface-related defect species as $O_2^{\bullet-}$ ions for the $SrTiO_3$ NCs samples prepared aerobically without hydrazine. We propose that lactate ligands of the Ti(IV) precursor reduce O_2 to $O_2^{\bullet-}$ ions during the hydrothermal synthesis (200 °C and basic conditions). To support this mechanism and generality of $O_2^{\bullet-}$ formation and the importance of lactate ions, we prepared $BaTiO_3$ NCs from the same lactate-containing Ti(IV) precursors as well as a lactate-free precursor. $O_2^{\bullet-}$ defects are observed only in NCs prepared from a precursor that contains lactate ligands under aerobic conditions. These results provide convincing evidence that the choice of precursor can lead to unique surface chemistry and the formation of $O_2^{\bullet-}$ defects on titanate-based NCs.

2.2 Experimental

2.2.1 Materials

Strontium hydroxide octahydrate ($Sr(OH)_2 \cdot 8H_2O$, 99%, Alfa Aesar), barium hydroxide octahydrate ($Ba(OH)_2 \cdot 8H_2O$, 99.95%, Acros Organics), titanium(IV) bis(ammonium lactate) dihydroxide (TALH, 50% in water, Alfa Aesar), titanium(IV)

tert-butoxide ($\text{Ti}(\text{OBu})_4$, 99%, Acros Organics), 1-butanol (99%, Across Organics), tetramethylammonium hydroxide (NMe_4OH , Acros Organics), sodium hydroxide (NaOH , Certified ACS, Fisher Chemical), hydrazine hydrate (99%, Acros Organics), oleic acid (90%, Fisher Chemical), oleylamine (50%, TCI America), ethanol (200 proof, PHARMCO-AAPER) and hexanes (optima grade, Fisher Chemicals) were all used as received.

2.2.2 Synthesis of Nanocrystals

Synthesis of colloidal SrTiO_3 NCs from TALH. Synthesis of colloidal SrTiO_3 NCs was carried out by the hydrothermal method as reported previously by our group.²⁹ In a typical procedure, 1.25 mmol of each of TALH and $\text{Sr}(\text{OH})_2$ were dissolved in 30 mL of distilled water in a 45-mL Teflon-lined autoclave. The pH of the solution was then adjusted to 12.1 with a 10 M NMe_4OH solution followed by the addition of oleic acid (2.5 mmol). The reaction vessel was then sealed and heated to 200 °C in the oven for 24 hours. The resulting NCs were collected, washed with ethanol three times, and suspended in hexanes. This procedure described above is referred to herein as method-I.

Method-I was modified to prepare two different sets of colloidal NCs. In the first modification, referred to as method-II, hydrazine hydrate (5 mmol) was added to reaction mixture before adding oleic acid. The pH of the solution remained at 12.1 after the addition of hydrazine hydrate. The last modification, method-III, was identical to method-I but the reaction solution was purged with argon gas for 60 minutes prior to closing the Teflon-lined reaction vessel.

Synthesis of colloidal BaTiO₃ NCs from TALH. Synthesis of colloidal BaTiO₃ NCs was also carried out by a similar hydrothermal method and TALH precursor but included constant magnetic stirring. In a typical preparation, 1.5 mmol of each TALH and Ba(OH)₂ were dissolved in 24 mL distilled water followed by the addition of 6 mL of NaOH (5 M). The reaction solution was then transferred to a 45-mL Teflon-lined autoclave and oleylamine (6 mmol) and oleic acid (6 mmol) were added. The pH of the reaction solution was 12. The sealed autoclave was placed in custom-made aluminum block housing heated to 215 °C and constantly stirred for 24 h using a stirring hotplate. After the synthesis, the autoclave was cooled to room temperature and the solid product was collected, washed with EtOH, and dissolved in nonpolar solvents such as hexanes to produce a transparent solution. This procedure described above is referred to herein as method-A. In addition, Method-A was modified to prepare another set of BaTiO₃ NCs referred to as method-B. In this modification, hydrazine hydrate (5 mmol) was added before adding the oleylamine and oleic acid. Method-A and -B for BaTiO₃ mimic method-I and -II for the SrTiO₃ NCs.

Synthesis of BaTiO₃ NCs from Ti(OBu)₄. Synthesis of lactate-free BaTiO₃ NCs was done using an aerobic hydrothermal method adopted from the literature.³¹ In a typical experiment, aqueous solutions containing 1 mmol of Ba(OH)₂ dissolved in 5 mL water and 12.5 mmol of NaOH dissolved in 5 mL water were mixed with 10 mL of 1-butanol containing 1 mmol Ti(OBu)₄ and 2.5 mL of oleic acid. The pH of the solution was 12.2. The resulting precursor solution was transferred to a 45-mL autoclave, closed under ambient conditions, and heated to 180 °C for 18 hours. The resulting

product was collected, washed with EtOH several times, and dispersed in non-polar solvents such as hexanes.

2.2.3 Physical Characterization

All measurements were collected at room temperature unless specified otherwise. Electronic absorption spectra were collected on colloidal suspensions in air-tight quartz cuvettes with 1-cm pathlengths (Cary 50 Bio). EPR spectra were measured at X-band frequency (9.6 GHz) with a Bruker Elexsys-500 equipped with a Super High QE (ER4123SHQE) cavity. Transmission electron microscopy (TEM) images of NCs deposited onto copper grids (CF400-CU-50, Electron Microscopy Sciences) with a 3 nm carbon coating (JEOL TEM-2000FX). Powder X-ray diffraction patterns were collected in the Bragg-Brentano configuration with a Cu K- α source (Rigaku SmartLab SE). FTIR spectra were collected using Bruker Alpha-P equipped with a diamond attenuated total reflectance (ATR) crystal. High-resolution electrospray ionization mass spectra (ESI-MS) were collected in negative ion mode with Bruker MicroTOF-II. In a typical sample preparation for ESI-MS, the aqueous reaction mixture from the hydrothermal synthesis was centrifuged to separate NCs from the rest of the water-soluble side-products. The aqueous supernatant was used for ESI-MS measurements with no further purification.

2.3 Results and Discussion

The crystallinity and phase purity of the SrTiO₃ NCs was confirmed by powder X-ray diffraction (see Supporting Information (SI) Figure 2.7 and Table 2.1). All samples matched to cubic SrTiO₃ in the Pm $\bar{3}$ m space group.³² The similar powder patterns among the samples suggest that modification of hydrothermal synthesis either by

adding hydrazine hydrate or degassing the precursor solution has no effect on the crystalline phase or phase purity. TEM images also demonstrate cuboidal morphology of the SrTiO₃ NCs prepared by method-III in agreement with previously reported NC shapes made by method-I or method-II (see SI Figure 2.8).²⁹

FTIR absorption measurements collected on SrTiO₃ NCs prepared from each method are shown in Figure 2.1. The prominent peaks at 1540 cm⁻¹ and 1450 cm⁻¹ correspond to the symmetric and asymmetric stretches of the bound carboxylate head group of the surface oleate ligand and are observed in each sample. Aliphatic bands from the oleate surface ligands also appear *ca.* 2900 cm⁻¹. All samples thus indicate the NC surfaces are passivated by oleate ligands within the sensitivity of the instrument. These spectra are similar to prior reports of SrTiO₃ made in a similar method¹⁷ but at pH 13.5. There is no evidence of N-H stretches from hydrazine in the NCs prepared by method-II that would appear around 3400 cm⁻¹.³³ This observation suggests that hydrazine is not being adsorbed on the surface of SrTiO₃ NCs. Furthermore, hydrazine adsorption is not expected as it is known to decompose rapidly at this reaction temperature (200 °C) in the presence of oxygen to possible products of N₂ and H₂O.³⁴ Fujinami *et al.* also observed essentially no change in pH of the reaction solution after SrTiO₃ NC synthesis regardless of whether hydrazine was present, which was speculated to play a different role than decomposing to NH₃.¹⁷

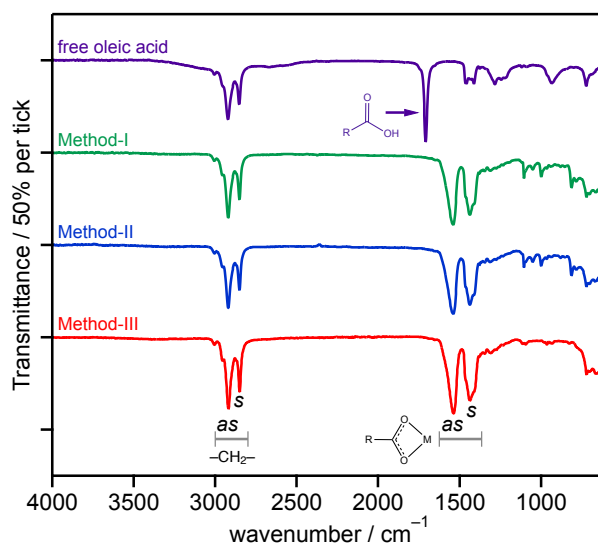


Figure 2.1 FTIR spectra of solvent-evaporated SrTiO₃ NCs prepared by method-I (green), method-II (blue) and method-III (red). The top spectrum belongs to oleic acid (purple). The *s* and *as* next to the peaks refer to the symmetric and asymmetric stretches associated with the aliphatic C-H or carboxylate vibrations.

The electronic absorption spectrum of all colloidal SrTiO₃ NCs dispersed in hexanes is dominated by the band edge absorption at *ca.* 3.25 eV consistent with bulk SrTiO₃ (see Figure 2.2).³⁵ However, NCs prepared by the aerobic hydrothermal method-I also exhibit broad absorption throughout the entire visible region and appear reddish-brown (see color photographs shown in Figure 2.2 inset). The broad sub-bandgap absorption is absent in SrTiO₃ samples prepared using either method-II or method-III.

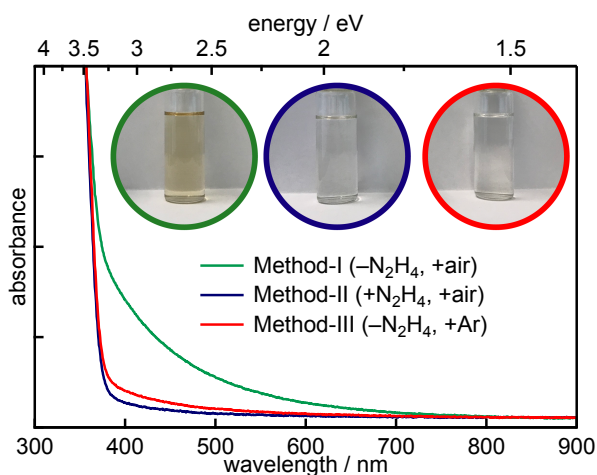


Figure 2.2 Electronic absorption spectra of colloidal SrTiO₃ NCs dispersed in hexanes prepared by method-I (green), method-II (red) and method-III (blue). Optical densities were matched at ~350 nm. Color photographs of NCs dispersed in hexanes are shown in insets.

The origin of the visible absorption in the method-I NCs suggests the presence of defects with sub-bandgap levels. Native point defects such as V_O or Ti^{3+} are also known to give rise to visible light absorption in bulk SrTiO₃. Electronic transitions from the valance band to different V_O charge states have been observed between 450 and 550 nm.³⁶⁻³⁸ Rice *et al.* recently attributed a transition involving the V_O to an absorption band at *ca.* 405 nm in reduced SrTiO₃ crystals.³⁹ Similar V_O defects have also been described as polarons. However, the electronic transition of the localized electron trapped on Ti^{3+} sites to the conduction band minimum is broad and centered in the near-IR region.^{40,41} To explore the origin of the brownish coloration of the SrTiO₃ NCs prepared by method-I, we carried out EPR spectroscopy measurements on all three types of SrTiO₃ NCs.

The room temperature, X-band EPR spectrum of the SrTiO₃ NCs prepared by method-I is shown in Figure 2.3. These NCs exhibit a narrow resonance near 351 mT corresponding to a g -value of 2.003. Upon addition of hydrazine (method-II) or purging the precursors with argon (method-III), this EPR signal disappears. Carter *et al.* have identified four types of oxygen-related defects on the surface of TiO₂ (anatase-type) using EPR spectroscopy and the signal at $g=2.003$ was assigned to surface-adsorbed paramagnetic superoxide radicals ($O_2^{\bullet-}$, $S = 1/2$).⁴² Those authors also propose the formation involves electron transfer from trapped electrons at Ti^{3+} sites to surface-O₂ to produce $O_2^{\bullet-}$. We also tentatively assigned the $g=2.003$ signal to $O_2^{\bullet-}$ in our initial report on the synthesis and characterization of SrTiO₃ NCs.²⁹ We also demonstrated in a more recent report that this $g = 2.003$ signal can be increased upon

exposing the photochemically reduced (Ti^{3+} -rich) Cr^{3+} -doped SrTiO_3 NCs to air.⁴¹ Others have also reported a TiO_2 -based NCs treated with H_2O_2 that display similar absorption in the visible region originating from surface-adsorbed $\text{O}_2^{\bullet-}$ defects.^{43,44} Based on the similarity to previous reports and observations,^{41,43} we attribute the reddish-brown coloration and EPR resonance at $g=2.003$ observed in Figures 2 and 3, respectively, to the presence of $\text{O}_2^{\bullet-}$ radicals on the surface of SrTiO_3 NCs prepared under aerobic hydrothermal conditions (method-I). We further assign the visible absorption to defect-induced sub-bandgap states that give rise to charge transfer-like transitions.

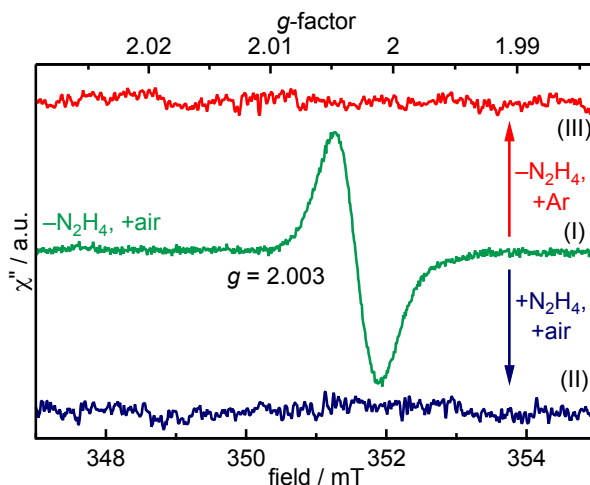


Figure 2.3 Room temperature EPR spectra of NCs prepared by method-I (middle/green), method-II (bottom/blue), and method-III (top/red).

Our working hypothesis is that the $\text{O}_2^{\bullet-}$ ions are formed via reduction of dissolved O_2 during the hydrothermal synthesis. The lack of $\text{O}_2^{\bullet-}$ formation when prepared by either method-II or method-III is thus presumably due to a very low O_2 content in the reaction solution during the hydrothermal synthesis. Despite these results, we do not have a direct mechanistic understanding of how the $\text{O}_2^{\bullet-}$ is formed during synthesis. One outstanding question that we sought to answer is the identity of the electron donor. The $\text{O}_2^{\bullet-}$ formation at TiO_2 surfaces is believed to proceed via electron transfer

from Ti^{3+} sites to adsorbed O_2 .^{45,46} In bulk SrTiO_3 , oxygen vacancies are thought to exist between adjacent localized Ti^{3+} sites,^{47,48} which could serve as an electron donor. However, no spectroscopic evidence for the existence of V_{O} and any localized Ti^{3+} defects were found in SrTiO_3 colloidal NCs. It is possible that these species are only formed under hydrothermal conditions during synthesis but suggests the presence of alternative electron donors such as lactate ions.

The hydrothermal synthesis of SrTiO_3 NCs involves the hydrolysis of titanium(IV) bis(ammoniumlactato)dihydroxide (TALH) as the titanium source under basic conditions. Möckel *et al.* have proposed that the hydrolysis of TALH to form TiO_2 NCs produces free ammonium lactate in neutral or slightly acidic conditions.⁴⁹ Seisenbaeva *et al.* have proposed hydrolysis of TALH forms a stable, tetrameric $(\text{NH}_4)_8\text{Ti}_4\text{O}_4(\text{lactate})_8$ cluster that serves as a precursor to the formation of TiO_2 NCs under basic conditions.⁵⁰ In both of above-mentioned references, lactate ions are commonly observed as hydrolysis by-product. Lactate ions have been shown to serve as electron donors in various chemical and biological systems to produce O_2 -based radicals.^{51,52} Therefore, we hypothesize that the lactate ions reduce dissolved O_2 under extremely high pH to produce $\text{O}_2^{\bullet-}$ and the corresponding by-product pyruvate anions as described by Figure 2.4 shown below.

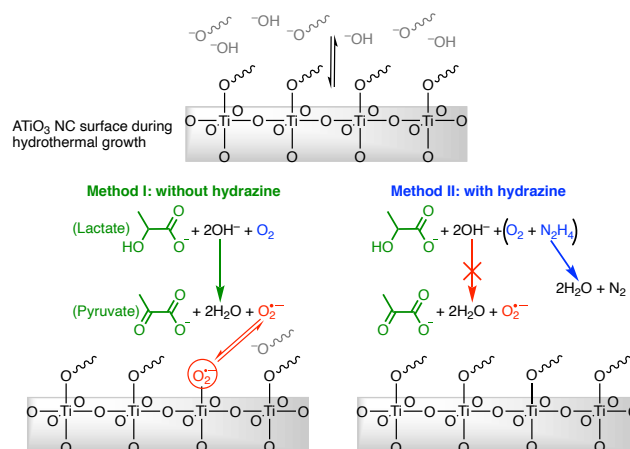


Figure 2.4 Proposed mechanism of superoxide formation and adsorption onto the dynamic ATiO₃ surface during hydrothermal synthesis.

In Figure 2.1, we depict ATiO₃ surface that is stabilized by oleate ligands under hydrothermal conditions (pH 12 and 200 °C). At this high temperature, we expect the surface ligands to rapidly exchange between bound and free. This dynamics surface would allow for O₂^{•-} to bind to the surface in Method-I and become stabilized upon cooling to room temperature. In the anaerobic cases, Method-II and Method-III, there is no direct way to produce O₂^{•-} in solution and therefore rule out the possibility of O₂^{•-} adsorption on NC surface.

To confirm the presence of lactate and pyruvate anions, we performed mass spectrometry analysis of reaction supernatants from the various syntheses without purification. Figure 2.5 shows a zoomed-in region of the negative ion mode ESI-MS spectra collected from all three methods. The dominant peak at *m/z* 89 is attributed to lactate ions which are formed in all three methods consistent with our hypothesis. Despite being weak, the presence of a minority peak at *m/z* 87 is detected in each of the samples and corresponds to pyruvate ions. When normalized to the lactate signals for the three methods, we detect about a three-fold higher relative intensity of pyruvate anion signal from the supernatant from method-I. While not absolutely

quantitative, the ESI-MS results suggest this conversion yield is <1%. Despite this limitation, the observation is consistent with the proposed mechanism that at least some fraction of the lactate ions undergo the conversion shown in Figure 2.4. The lower O₂ content in method-III precludes the O₂^{•-} formation which in turn can be supported by observing relatively a lower pyruvate concentration. However, the absence of any O₂^{•-} defects together with much lower pyruvate content in method-II questions the role of hydrazine in controlling the O₂^{•-} formation during hydrothermal synthesis.

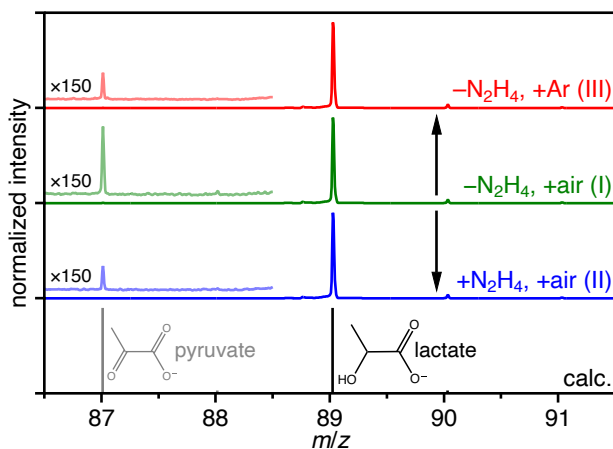


Figure 2.5 Negative-ion mode ESI-MS of supernatants from the various reactions in the region of the lactate $[\text{C}_3\text{H}_5\text{O}_3]^-$ and pyruvate $[\text{C}_3\text{H}_3\text{O}_3]^-$ anions. The spectra are normalized to the major peak from the lactate anion at m/z 89.

Based on the results above, hydrazine appears to control the formation of O₂^{•-} defects by limiting the available O₂ in the synthesis. Hydrazine has been used to remove O₂ in industrial-scale reactors for decades.^{53,54} The reaction between hydrazine and O₂ produces N₂ and 2H₂O, which provides the inert atmosphere during synthesis as shown in Figure 2.4.³⁴ Post-synthetic attempts to control the concentration of defects on the surface of the SrTiO₃ NCs prepared by Method-I were also performed. When an excess of hydrazine hydrate is added to an aqueous mixture of defect-rich NCs and heated under aerobic hydrothermal conditions, no change in

the appearance of the NCs was observed suggesting no change in surface defects (data not shown).

To test the generality of this approach, we extended the study to BaTiO₃ NCs using a similar hydrothermal-based synthesis with TALH and a literature solvothermal method (see SI Figure 2.9-2.12 and Table 2.2). The room temperature EPR spectra of BaTiO₃ NCs prepared from all methods is shown in Figure 2.6. Similar to the results for SrTiO₃, the appearance of a O₂^{•-} resonance at $g=2.003$ is observed in BaTiO₃ samples prepared using TALH and air. The formation of these O₂^{•-} defects can be inhibited by adding hydrazine to remove O₂ from the reaction that is consistent with our proposed mechanism.

To confirm the critical role of lactate ions in the formation of these surface O₂^{•-} defects, we synthesized BaTiO₃ NCs from a lactate-free titanium precursor, Ti(OBu)₄, using similar hydrothermal conditions as used above with TALH.³¹ The resulting lactate-free BaTiO₃ NCs are transparent below the band gap and do not display any detectable EPR signal as shown in Figure 2.6 and in the SI. This result is entirely consistent with the critical role that lactate ions play in the formation of O₂^{•-} defects during hydrothermal synthesis of ATiO₃ and likely also TiO₂ NCs.

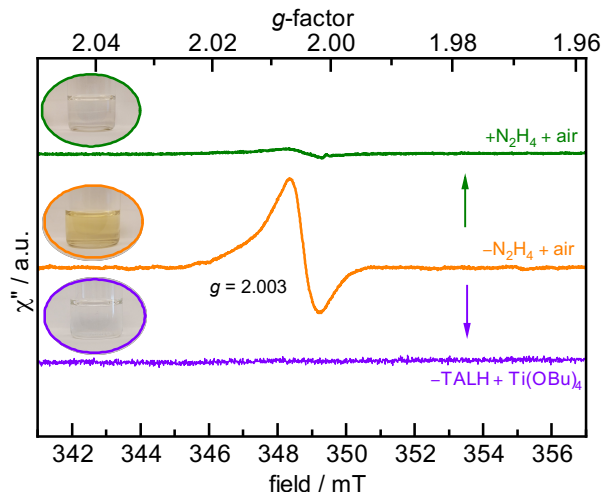


Figure 2.6 Room temperature EPR spectra of BaTiO₃ NCs prepared without hydrazine (orange/middle), with hydrazine (top/green) and from lactate-free precursor (bottom/purple). The color photographs of colloidal solution in hexanes for NCs prepared from all three methods are shown in insets.

2.4 Supporting Information

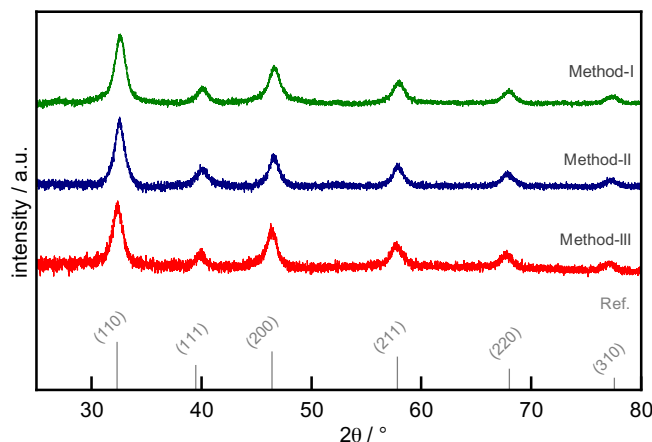


Figure 2.7 Powder X-ray diffraction patterns of SrTiO₃ NCs prepared from all three methods indexed to the cubic phase of bulk SrTiO₃ with Pm $\bar{3}$ m space group.⁵⁵

Table 2.1 The average lattice parameter and crystallite size calculated using Rigaku SmartLab Studio II for all three samples.

Preparation Method	Lattice parameter (Å)	Crystallite size (nm) [†]
I (aerobic without hydrazine)	3.9069	7.3 ± 0.3
II (aerobic with hydrazine)	3.9162	7.5 ± 0.4
III (anaerobic without hydrazine)	3.9163	7.6 ± 0.7

[†]Crystallite size is the average size calculated from the three most intense reflections: (110), (200) and (211).

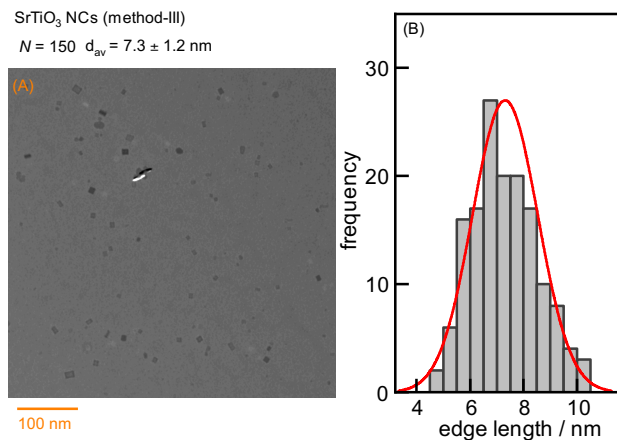


Figure 2.8 Representative transmission electron microscopy (TEM) image of SrTiO₃ NCs prepared from method III (degassed with argon) showing a clear cubic morphology of NCs. The size distribution plot from analyzing over hundred different particles in raw TEM images using ImageJ software displays an average size of 7.3 ± 1.2 nm in good agreement with crystallite size calculated from the diffraction pattern in Figure S1. The cubic morphology of NCs prepared from the method I and method II is already confirmed in our previous report.⁵⁶

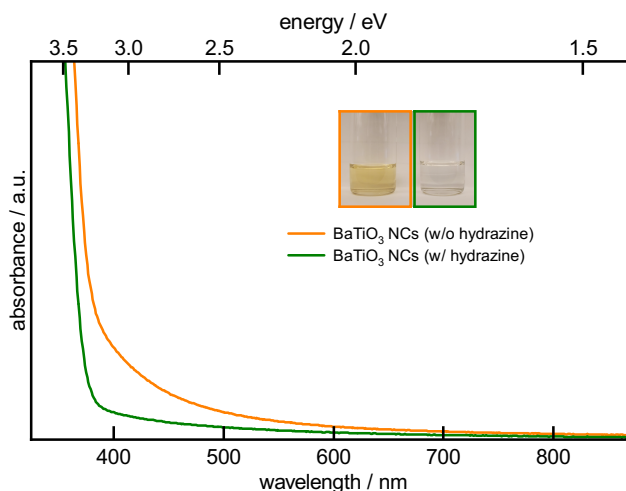


Figure 2.9 Room temperature electronic absorption spectra of BaTiO₃ NCs prepared from method-A (without hydrazine/orange) and method-B (with hydrazine/green). The color photographs of both samples are shown in inset.

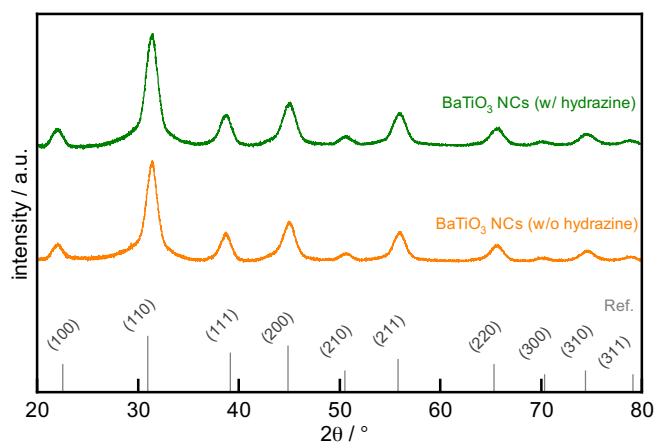


Figure 2.10 Powder X-ray diffraction pattern of BaTiO₃ NCs prepared from method-A (without hydrazine/orange) and method-B (with hydrazine/green) indexed to the cubic phase of BaTiO₃.⁵⁷

Table 2.2 The average lattice parameter and crystallite size calculated using Rigaku SmartLab Studio II for both samples.

Sample NCs	Lattice parameter (Å)	Crystallite size (nm)†
With hydrazine	4.0230	5.66 ± 0.50
Without hydrazine	4.0335	6.54 ± 0.53

†Crystallite size is the average size calculated from the three most intense reflections: (110), (200) and (211).

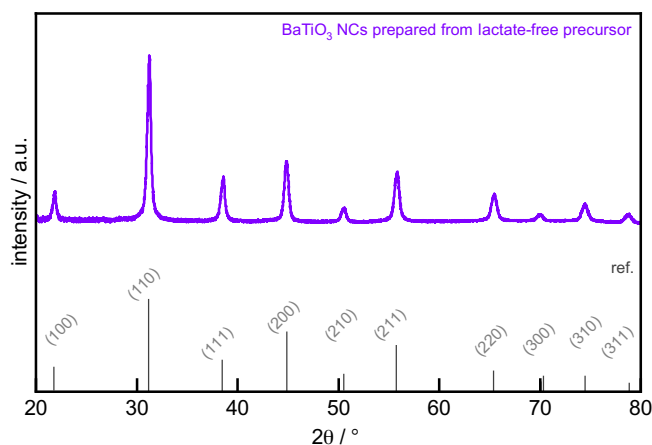


Figure 2.11 Powder X-ray diffraction pattern of BaTiO₃ NCs prepared from lactate-free method indexed to the cubic phase of single-crystal BaTiO₃.

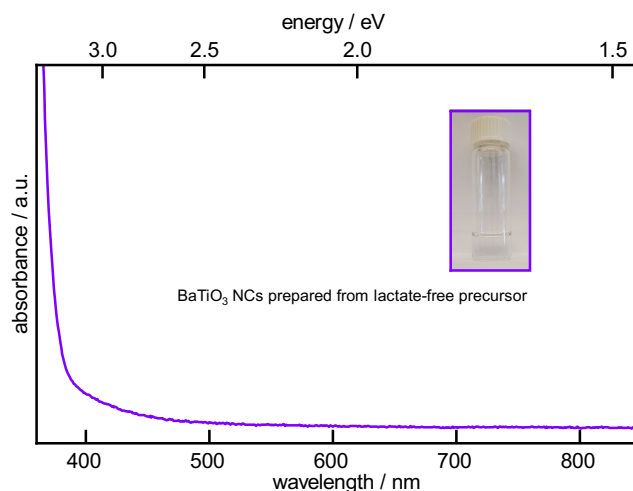


Figure 2.12 Room temperature electronic absorption spectrum of BaTiO₃ NCs prepared from the lactate-free precursor. The color photograph of hexanes solutions of as-prepared NCs is shown in inset.

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CHAPTER 3

TUNABLE REDOX ACTIVITY AT IRON CENTERS IN COLLOIDAL STRONTIUM AND BARIUM TITANATE NANOCRYSTALS

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3.1 Introduction

Exploiting the properties of semiconductor nanomaterials by incorporating impurities constitutes a long-standing challenge to materials chemistry that has hindered their implementation in various technologies.¹⁻⁴ Numerous studies have focused on the design of high-quality colloidal nanocrystals (NCs) with controlled incorporation of either aliovalent dopants to tune the carrier concentration or magnetic dopants for spin-based electronic or quantum information processing applications.⁵⁻⁸ Significant progress has occurred in recent years with respect to introducing excess charge carriers into metal oxide semiconductors through aliovalent doping,^{9,10} control over material stoichiometry,^{11,12} photochemical (also known as photodoping),^{13,14} doping via electron transfer from a chemical reductant,¹⁵ and electrochemical doping.¹⁶

An emerging class of substitutional dopants that function simultaneously as aliovalent dopants and paramagnets has received much attention due to their multifunctionality. Examples of such dopants in oxide semiconductors are Fe³⁺-doped ZnO (Fe³⁺:ZnO),^{10,17} where Fe³⁺ dopant acts as an electron donor, and Cr³⁺:SrTiO₃ where Cr³⁺ substitutes at the Ti⁴⁺ site and behaves as an acceptor.^{18,19} The aliovalent magnetic dopants of relevance to our studies are the 3d transition metal ions in Ti-based oxide semiconductors. SrTiO₃ is a wide-gap semiconductor with the bottom of

the conduction band having primarily empty Ti-3d character. This electronic structure feature is common with so-called d^0 semiconductors and can be exploited to produce quasi-delocalized Ti^{3+} defect centers (self-trapped electrons) after an increase in carrier concentration.^{15,20} Magnetic ordering has also been reported in Ti-based oxide semiconductors containing aliovalent magnetic dopants including $Cr^{3+}:TiO_2$, $Co^{2+}:TiO_2$, $Fe:SrTiO_{3-\delta}$, and $Co^{2+}:(La,Sr)TiO_3$.²¹⁻²⁴

We recently reported that the spin relaxation time of Cr^{3+} dopants in colloidal $SrTiO_3$ NCs is accelerated by the presence of photochemically introduced Ti^{3+} defects.²⁵ We demonstrated an apparent enhancement of this effect specifically in the $Cr^{3+}:SrTiO_3$ NCs, relative to the Cr^{3+} -doped bulk powders.¹⁵ The discovery of this efficient and reversible control over the spin-relaxation time of localized Cr^{3+} motivated our current work to explore the behavior of additional aliovalent magnetic dopants in d^0 semiconductor NCs. Herein we report the preparation of sub-10 nm Fe-doped $SrTiO_3$ and $BaTiO_3$ colloidal NCs using hydrothermal methods described previously by our group.^{26,27} We confirm that the Fe is in Fe^{3+} state and is located at the octahedral Ti^{4+} sites in as-prepared NCs by electron paramagnetic resonance (EPR) spectroscopy. Upon anaerobic photodoping of colloidal NCs, the EPR signal of Fe^{3+} disappears with increasing Ti^{3+} concentration. The lack of Fe^{3+} EPR signal in the photodoped $SrTiO_3$ and $BaTiO_3$ persists down to 100 K, where the Ti^{3+} EPR signal is observed. These results are consistent with the electrons trapping at the Fe^{3+} sites prior to accumulating at Ti^{3+} defects in these NCs. This electron trapping phenomenon reduces Fe^{3+} to EPR-silent Fe^{2+} centers and can be effectively manipulated as a

function of irradiation time. All of the spectroscopic changes induced by photodoping are fully reversible upon aerobic oxidation of the photodoped sample.

3.2 Experimental

3.2.1 Materials

Strontium hydroxide octahydrate ($\text{Sr}(\text{OH})_2 \cdot 8\text{H}_2\text{O}$, 99%, Alfa Aesar), barium hydroxide octahydrate ($\text{Ba}(\text{OH})_2 \cdot 8\text{H}_2\text{O}$, 99%, Alfa Aesar), titanium(IV) bis(ammonium lactate) hydroxide (TALH, 50% in water, Alfa Aesar), tetramethylammonium hydroxide (NMe_4OH , Acros Organics), sodium hydroxide (NaOH , Certified ACS, Fisher Chemical), hydrazine hydrate ($\text{N}_2\text{H}_4 \cdot \text{H}_2\text{O}$, 99%, Acros Organics), oleylamine (>50%, TCI America), oleic acid (>95%, Fisher Chemical), iron nitrate nonahydrate ($\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$, 99.99% Sigma Aldrich), ethanol (200 proof, PHARMCO-AAPER), QC-21 elements standard (PerkinElmer, 5% HNO_3), nitric acid (Certified ACS Plus, Fisher Chemical), hydrochloric acid (Certified ACS Plus, Fisher Chemical), toluene (Optima, Fisher Chemicals) and hexanes (Optima, Fisher Chemicals) were all used as received.

3.2.2 Synthesis of Nanocrystals

Pure and Fe-doped SrTiO_3 colloidal NCs. Synthesis of colloidal SrTiO_3 NCs was carried out by a modified hydrothermal method reported earlier by Harrigan *et al.*²⁶ In a typical synthesis, 1.25 mmol of TALH and 1.25 mmol of $\text{Sr}(\text{OH})_2 \cdot 8\text{H}_2\text{O}$ were dissolved in 30 mL of distilled water. The pH of the solution was then adjusted to 12.1 with an aqueous solution of NMe_4OH (10 M) followed by the addition of oleic acid (2.5 mmol) and hydrazine (5 mmol). The resulting solution was transferred to a 45 mL Teflon-lined autoclave (4744 General Purpose Acid Digestion Vessel, Parr Instrument Company) and heated to 200 °C in an oven for 24 hours. The resulting NCs were

collected, washed with ethanol several times, suspended in hexanes, and sonicated for about 30 minutes to produce a cloudy suspension. Finally, a transparent hexane layer containing NCs was extracted by centrifuging the cloudy suspension (5 min at 4000 rpm). These NCs can be further purified by precipitation with ethanol. Synthesis of nominally 0.5% Fe-doped SrTiO₃ NCs was carried out by adding 625 μ L of 0.01 M aqueous solution of Fe(NO₃)₃·9H₂O right after adjusting the pH of the precursor solution. Similarly, nominally 1% Fe-doped SrTiO₃ NCs were prepared by adding 1.25 mL of 0.01 M aqueous solution of Fe(NO₃)₃·9H₂O.

Pure and Fe-doped BaTiO₃ colloidal NCs. Synthesis of BaTiO₃ NCs was carried out by a similar hydrothermal method but included constant magnetic stirring.²⁷ In a typical synthesis, 1.5 mmol of each of TALH and Ba(OH)₂ were dissolved in 24 mL distilled water followed by 6 mL of 5 M NaOH aqueous solution. The reaction mixture was then transferred to a 45 mL Teflon-lined autoclave, and oleylamine (6 mmol), oleic acid (6 mmol), and hydrazine (6 mmol) were added. The sealed autoclave was placed in custom-made aluminum block housing heated to 215 °C and stirred continuously for 24 hours using a stirring hotplate. After the synthesis, the autoclave was cooled down to room temperature, and the crude product was collected along with water into two separate test tubes and washed with ethanol (33% by volume) twice. The solid product was then dissolved in hexanes to produce a cloudy suspension with sonication (15 minutes). Finally, the hexane layer containing the NCs was extracted by centrifuging the suspension. Synthesis of nominally 0.5% Fe-doped BaTiO₃ NCs was carried out by adding 750 μ L of 0.01 M aqueous solution of Fe(NO₃)₃·9H₂O before adding NaOH.

3.2.3 Physical characterization

All measurements presented below were collected on either pure or 0.5% Fe-doped SrTiO_3 NCs (referred to as Fe:SrTiO_3 hereafter) and pure or 0.5% Fe-doped BaTiO_3 NCs (referred to as Fe:BaTiO_3 hereafter) unless specified otherwise. Electronic absorption and emission spectra were collected at room temperature on colloidal NCs in air-tight quartz cuvette on Cary 50 Bio and Cary Eclipse, respectively. Near-IR absorption measurements were collected on Varian 670 FT spectrometer equipped with an InGaAs detector and near-IR quartz beam splitter using a 1-mm pathlength quartz cuvette. The functional groups and surface chemistry analysis of solvent-free NCs were carried out on a Bruker Alpha-P FTIR spectrophotometer equipped with an attenuated total reflectance (ATR) platinum diamond optic. EPR measurements were carried out on continuous wave (CW) Bruker Eleksys-500 fitted with Super High QE X-band cavity (ER4123SHQE) and cryostat (ESR-900, Oxford) for low-temperature measurements. For the photodoping experiments, a 75 W xenon lamp was used (60000 Q series, Oriel Corporation). Transmission electron microscopy (TEM, JEOL 2000FX) images were measured from freshly prepared NCs drop-casted onto carbon-coated (3nm) copper grids (CF400-CU-50, Electron Microscopy Sciences). The average size distribution and standard deviation were generated by analyzing a few hundred particles from TEM images using ImageJ software. Powder X-ray diffraction patterns were collected in a Bragg-Brentano configuration and Cu K-source (Rigaku Smart Lab SE). Metal ion stoichiometry was measured on dissolved ensembles of NCs and measured by an inductively coupled plasma - optical emission spectrometer (ICP-OES, Perkin Elmer Optima 4300 DV).

3.2.4 Photodoping

Photodoping experiments were performed to introduce excess electrons in as-prepared NCs. Sample preparation was completed in the argon-filled glove box where NCs suspended in hexanes were transferred to a 1 cm air-tight cuvette before removal. Samples were photodoped by prolonged exposure to unfocused irradiation from 75 W xenon lamp in the presence of EtOH (0.5 mL, >100 eq/NC) as a sacrificial reductant. The electronic absorption spectra were taken periodically during the entire photodoping process. It was estimated that the NCs were photodoped to their saturation limit when no further spectroscopic changes could be observed following prolonged exposure. For in-situ EPR measurements, NCs were transferred to a 4-mm air-free EPR tube, and photodoping procedure inside the EPR resonator was carried out by UV irradiation via fiber optics. The lamp was warmed up 30 min before each photodoping experiment in order to achieve uniform temperature and irradiation intensity.

3.3 Results and Discussion

The crystallinity and phase purity of as-prepared undoped and Fe-doped SrTiO₃ NCs were confirmed by powder X-ray diffraction measurements (see SI Figure 3.6). The diffraction patterns of both samples are indexed to cubic SrTiO₃ (space group: Pm $\bar{3}$ m) with no additional secondary phases.²⁸ A slight increase of the lattice parameter was observed in the Fe:SrTiO₃ NCs (see Table 3.1). This typical lattice expansion is attributed to the substitution of a relatively larger Fe³⁺ ion (0.79 Å) at the Ti⁴⁺ (0.74 Å) sites and is consistent with other spectroscopic evidence of Fe³⁺ substitution shown below.^{29,30} TEM images also confirmed that the as-prepared

product is nanocrystalline and exhibits cubic morphology with an average edge length of $7.7 \text{ nm} \pm 1.9 \text{ nm}$ in good agreement with the size calculated from the diffraction pattern (see Figure 3.7). The elemental analysis showed a total of 1.03 % Fe content in the nominally 0.5% Fe-doped SrTiO_3 NCs (see Figure 3.8 and Table 3.2). The titanium precursor (TALH) also contains small amounts of TiO_2 impurities due to slow hydrolysis at pH 8.5. The presence of such impurities thus overestimates the concentration of Ti^{4+} in the reaction solution.

Figure 3.1A shows the electronic absorption spectra and the corresponding color photographs of concentrated solutions of SrTiO_3 and $\text{Fe}:\text{SrTiO}_3$ NCs in hexanes. The absorbance at *ca.* 3.25 eV dominates the electronic absorption spectra of both the samples, consistent with the bandgap transition energy of SrTiO_3 .³¹ $\text{Fe}:\text{SrTiO}_3$ NCs exhibit an additional broad sub-bandgap absorption that tails into the visible region that is responsible for the yellowish-brown color of the concentrated suspensions (see photographs in Figure 3.1A insets). This absorption tail was previously observed in bulk $\text{Fe}:\text{SrTiO}_3$ involving deep Fe^{3+} impurity levels and explicitly attributed to a valence band-to-metal charge transfer transition.^{32,33}

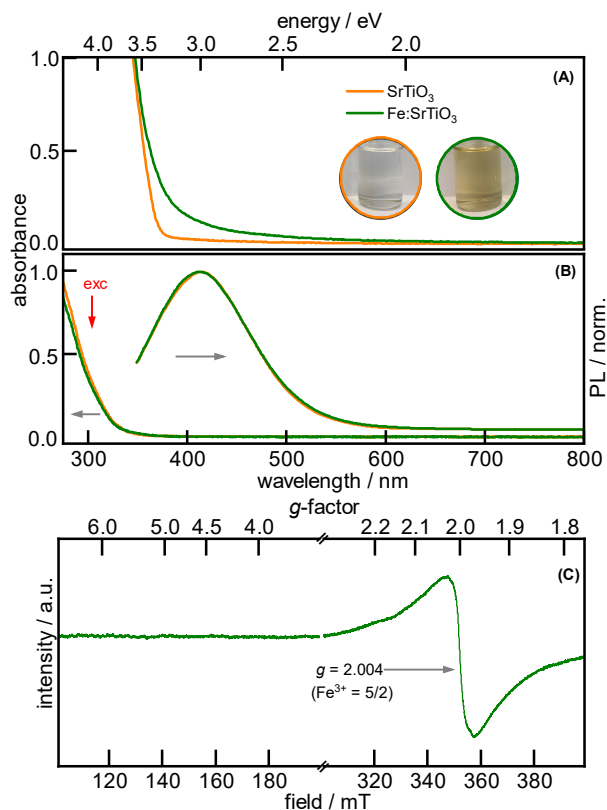


Figure 3.1 (A) Room temperature electronic absorption spectra and color photographs of concentrated solutions of SrTiO_3 (orange) and Fe:SrTiO_3 (green) in hexanes. (B) Absorption and PL spectra of dilute solutions of both samples excited at 300 nm. The left axis represents the optical density of dilute solutions that is ≤ 0.5 at the excitation wavelength, while that normalized PL intensity is displayed on the right axis. (C) Room temperature EPR spectrum of as-prepared Fe:SrTiO_3 NCs showing a single, broad resonance at ~ 350 mT ($g = 2.004$).

Figure 3.1B displays the electronic absorption and photoluminescence (PL) spectra of dilute suspensions of SrTiO_3 and Fe:SrTiO_3 NCs excited at 300 nm. The normalized PL spectra of both the samples appear identical, with a broad emission feature centered around ~ 3.0 eV. This PL energy is in agreement with the previously reported emission described as self-trapped excitons' (STE) recombination at the surface of SrTiO_3 .^{34,35} There is no discernable variance between the PL energy of SrTiO_3 and Fe:SrTiO_3 NCs. This observation suggests that both the samples possess similar surface chemistry and that radiative recombination from the STE state is faster than the energy transfer to Fe-centered excited states. FTIR measurements

were also utilized on solvent evaporated NCs (see Figure 3.9) to understand the surface coordination environment. The FTIR spectra display the characteristic symmetric and asymmetric carboxylate stretches (doublet at 1540 cm^{-1} and 1450 cm^{-1}) and aliphatic bands around 2900 cm^{-1} that are consistent with oleate passivation of the NC surface.

To probe the local environment of Fe^{3+} dopants in as-prepared $\text{Fe}:\text{SrTiO}_3$ NCs, we employed conventional X-band EPR spectroscopy, which is well-reported to replace octahedral Ti^{4+} in SrTiO_3 lattice.^{36,37} For Fe^{3+} ($3d^5$) in an environment of cubic symmetry (ground state term ${}^6\text{A}_{1g}$), the spin degeneracy can only be lifted by a magnetic field, thus making Fe^{3+} a suitable dopant for investigation by EPR spectroscopy. Figure 3.1C shows the room temperature EPR spectrum of $\text{Fe}:\text{SrTiO}_3$ NCs. The broad isotropic resonance at ~ 350 mT ($g = 2.004$) corresponds to high-spin Fe^{3+} electronic states at the octahedral site in correspondence with literature.³⁸⁻⁴⁰ No evidence of pseudo-octahedral Fe^{3+} at the SrTiO_3 NC surface was detected in the low-field region.

We previously showed that both the spin-lattice (T_1) and spin-spin (T_2) relaxation times of Cr^{3+} ions in Cr-doped SrTiO_3 colloidal NCs could be effectively accelerated by the presence of fast-relaxing paramagnetic Ti^{3+} defects.²⁵ To probe the interaction between excess electrons and substitutional Fe^{3+} , we performed photodoping experiments on $\text{Fe}:\text{SrTiO}_3$ NCs. A detailed procedure for introducing excess electrons via photodoping is described in the experimental section. Briefly, UV irradiation of an anaerobic solution of as-prepared NCs creates electron/hole pairs

(excitons) followed by the rapid quenching of holes with ethanol which leaves excess electrons in NCs.

Figure 3.2 shows the electronic absorption spectra of Fe:SrTiO₃ NCs as a function of photodoping time. The absorption spectrum of as-prepared NCs is dominated by SrTiO₃ bandgap transition above ~3.25 eV. With increased photodoping time, a new absorption feature extending throughout the entire visible region and centered at ~1.5 eV (900 nm) appears, and the physical appearance of the sample changes to dark blue (see Figure 2 insets). This broad near-IR transition is assigned to the metal-to-metal charge transfer (MMCT) transition from localized Ti³⁺ sites to the conduction band reported previously in photodoped Cr-doped SrTiO₃ and TiO₂.^{25,41} The only change in the electronic absorption spectrum with increasing photodoping time is the growth of this MMCT transition. We have no evidence that photodoping creates delocalized electrons in the conduction band (e_{CB}^-). Recent reports showed that n-type doping of small anatase TiO₂ NCs can lead to accumulation of e_{CB}^- as observed spectroscopically by significant blue-shifts in the band-edge energy from the Moss-Burstein effect and formation of a localized surface plasmon resonance (LSPR) in the mid-IR region.⁵ Milliron reported similar spectroscopic changes caused by e_{CB}^- accumulation in Nb⁵⁺-doped TiO₂ NCs.^{6,42} The lack of absorption in the mid-IR region further confirms that e_{CB}^- are absent in our photodoped SrTiO₃ NCs (see Figure 3.10).

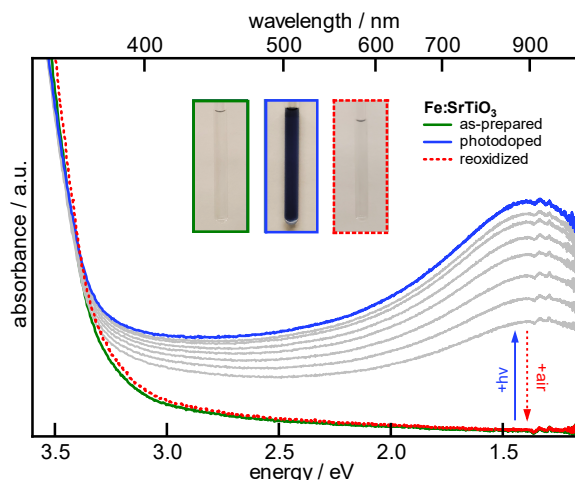


Figure 3.2 Electronic absorption spectra of an air-free solution of Fe:SrTiO₃ NCs in hexanes as a function of photodoping time. The green spectrum corresponds to as-prepared NCs. The intermediate grey lines show the data collected at various stages of photo-irradiation in the presence of EtOH. The blue spectrum belongs to maximally photodoped NCs, and the dotted red line overlapping the as-prepared data represents the reoxidized NCs upon opening the sample to air. The color photographs for as-prepared, photodoped, and reoxidized NCs are shown in insets.

EPR spectroscopy was also utilized to monitor any associated changes in substitutional Fe³⁺ with the concomitant occurrence of Ti³⁺ (3d¹, ²T_{2g} ground term) throughout the photodoping duration. Figure 3.3A shows the room temperature EPR spectra of Fe:SrTiO₃ NCs collected before and after photodoping. As shown in Figure 3.1C, the EPR spectrum of as-prepared Fe:SrTiO₃ NCs (green) exhibits only a single, broad resonance at $g = 2.004$ assigned to substitutional Fe³⁺ ions. However, after photodoping for 20 min, the Fe³⁺ signal completely disappears, and a new signal at ~355 mT ($g = 1.94$) is detected that is attributed to the formation of paramagnetic Ti³⁺ consistent with absorption measurements shown in Figure 3.2. The Ti³⁺ EPR signal in photodoped SrTiO₃ NCs is nearly isotropic. The disappearance of the Fe³⁺ signal and the simultaneous emergence of Ti³⁺ in the photodoped Fe:SrTiO₃ samples are entirely reversible processes. The above-mentioned spectroscopic changes were quantitatively recovered by exposing the photodoped sample to air, as shown by the red dotted spectra in Figures 3.2 and 3.3

These observations are similar to our recent report on Cr³⁺-doped SrTiO₃ NCs, where photodoping results in the disappearance of the Cr³⁺ EPR signal at room temperature.²⁵ The Cr³⁺ signal in that study reappeared after lowering the sample temperature to slow down the accelerated spin-relaxation. However, the EPR spectrum of photodoped Fe:SrTiO₃ NCs collected at 100 K (Figure 3.3A) indicates a permanent elimination of Fe³⁺ signal even at the cryogenic temperature. This observation suggests that Fe³⁺ may not have any spin-interactions with Ti³⁺ defects, unlike Cr³⁺ in photodoped SrTiO₃ NCs. The change in spin relaxation is caused by the near-resonant cross-relaxation phenomenon between Ti³⁺ and Cr³⁺ with similar *g*-values (Ti³⁺:1.945, Cr³⁺:1.978).^{43,44} We speculate that the Fe³⁺ signal at a much higher *g*-value (2.004) can make the spin cross-relaxation process with Ti³⁺ defects non-resonant and, therefore, less effective.

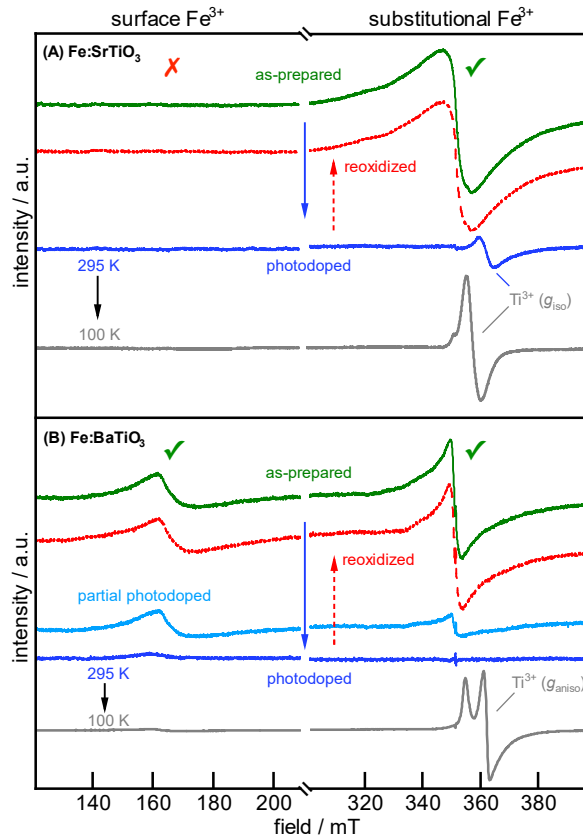


Figure 3.3 Room-temperature and cryogenic EPR spectra of an air-free solution of (A) Fe:SrTiO₃ and (B) Fe:BaTiO₃ NCs in hexanes. The green spectra in both samples correspond to as-prepared NCs, and the blue spectra denote the photodoped NCs. The dotted red lines overlapping the as-prepared data represent the reoxidized NCs upon opening the samples to air.

The observation of total disappearance of the Fe³⁺ signal in the SrTiO₃ NCs is consistent with at least some, if not all, of the substitutional Fe³⁺ dopants being reduced to EPR-silent Fe²⁺ dopants. This observation is also similar to our recent results on n-type Fe-doped ZnO NCs.^{9,45,46} To probe the relative positions of the Fe^{3+/2+} and Ti³⁺ energy levels in SrTiO₃ bandgap, we performed *in-situ* photodoping experiments (see SI Figure 3.11). The sequential disappearance of the Fe³⁺ EPR signal within a few seconds of photodoping followed by the appearance of the Ti³⁺ signal suggests (a) there is a gradual rise in Fermi level and (b) that the Fe^{3+/2+} redox level is situated deeper than the Ti³⁺ level. Figure 3.4 illustrates the proposed mechanism of electron trapping in Fe:SrTiO₃ NCs. The photoexcitation of charge carriers followed

by rapid hole quenching with EtOH raises the Fermi level below the conduction band, which results in the reduction of substitutional Fe^{3+} to Fe^{2+} prior to self-trapping at the Ti^{3+} defects. Upon exposing the photodoped sample to air, these changes are then reversed, and Fe^{3+} is quantitatively recovered.

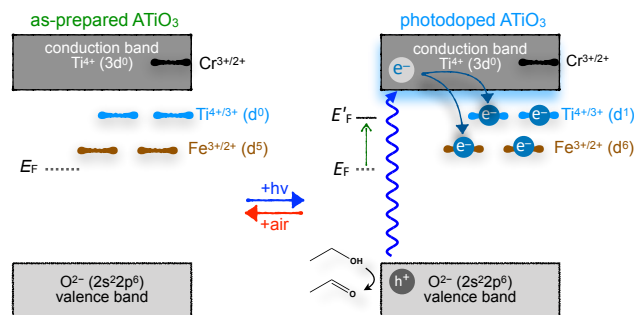


Figure 3.4 A schematic representation of reversible electron trapping at $\text{Fe}^{3+/2+}$ redox level in photodoped Fe:ATiO_3 NCs.

To validate the generality of this reversible electron trapping, we extended the study to Fe^{3+} -doped BaTiO_3 colloidal NCs prepared from a similar hydrothermal method (see SI Figure 3.12 and Table 3.3). Electronic absorption spectra of photodoped NCs of BaTiO_3 and Fe:BaTiO_3 show a similar but slightly blue-shifted broad absorption feature observed in photodoped SrTiO_3 NCs (see SI Figure 3.13). The EPR spectrum of as-prepared Fe:BaTiO_3 NCs is shown in Figure 3.3B is similar to the spectrum of Fe:SrTiO_3 but also displays an additional intense broad signal at ~ 160 mT ($g = 4.23$). This feature increases in relative intensity compared to the substitutional Fe^{3+} signal and is observed in nominally 1% Fe:SrTiO_3 NCs (see SI Figure 3.14). This feature at $g = 4.23$ is attributed to a fraction of Fe^{3+} dopants located near the surface of NCs with pseudo-octahedral symmetry. We attribute the higher relative intensity of this surface Fe^{3+} EPR signal to the smaller size of BaTiO_3 (~ 5 nm) that results in a larger fraction of Fe^{3+} dopants either on the surface or in near-surface sites compared to SrTiO_3 NCs (~ 8 nm). These surface Fe^{3+} sites have variation in the

coordinating ligands on the SrTiO₃ and BaTiO₃ NC surfaces. The alkylamine surface ligands used during the synthesis of the Fe:BaTiO₃ NCs are also known to remove surface Fe³⁺ dopants as observed in a post-synthetic ligand exchange process in Fe³⁺:ZnO NCs.¹⁰

Further support for the distinction between surface and substitutional Fe³⁺ in the BaTiO₃ NCs is based on a comparison of the spectral changes of this low-symmetry surface Fe³⁺ and substitutional Fe³⁺ at $g = 2.004$ as a function of photodoping time (see Figure 3.3B). Initially, the EPR signal for substitutional Fe³⁺ is more intense than the surface Fe³⁺; however, after photodoping for a short time (~5 min), the relative EPR signal intensities flip. Further photodoping leads to the complete disappearance of the substitutional Fe³⁺ with some residual surface Fe³⁺ signal still observed albeit weaker. The EPR spectrum of the photodoped Fe:BaTiO₃ NCs collected at 100 K also shows an anisotropic Ti³⁺ defect signal centered around ~355 mT as shown in Figure 3.3B. The anisotropy of the g -values observed for the Ti³⁺-related defect in photodoped colloidal Fe:BaTiO₃ NCs presented here ($g_{x,y} = 1.94$, $g_z = 1.97$) is exactly opposite to that reported in bulk BaTiO₃ powders and single crystals ($g_{x,y} = 1.97$, $g_z = 1.94$).⁴⁷ Both SrTiO₃ and BaTiO₃ NCs possess cubic structures by powder XRD. Therefore, the crystal field surrounding a self-trapped electron at the B-site (Ti³⁺) should be isotropic; however, Ti³⁺ is Jahn-Teller active and will distort to axial symmetry in both photodoped SrTiO₃ and BaTiO₃ albeit to varying extents.⁴⁸ The likelihood of Ti³⁺ defects being near or at the NC surface in the smaller BaTiO₃ NCs would also lead to distortion of the Ti³⁺ EPR signal due to mixed coordination with the lattice and surface ligands that adds a rhombic component to the g -anisotropy.

The anisotropic lineshape is less resolved in the photodoped SrTiO₃ and Fe:SrTiO₃ NCs and may be the result of the larger average particle size.

The existence of both axial and rhombic EPR signals associated with two distinct Ti³⁺ centers in thermal equilibrium was recently reported in photodoped TiO₂ nanoparticles.⁴⁹ This observation is unique to photodoped TiO₂ nanoparticles and has not been observed in the bulk. We postulate that the variation of the Ti³⁺ EPR signals in photodoped Fe:SrTiO₃ and Fe:SrTiO₃ NCs compared to the bulk n-type BaTiO₃ and Fe:SrTiO₃ is caused by differences in (1) the microenvironment surrounding the Ti³⁺ center and (2) its proximity to the NC surface. We are currently investigating the microscopic origins of the EPR signal in photodoped d⁰ semiconductor NCs. Nevertheless, there is no evidence of Fe³⁺ in photodoped Fe:BaTiO₃ NCs, which is consistent with our proposed model of electron trapping on Fe³⁺ in SrTiO₃ NCs. The drastic decrease in substitutional Fe³⁺ before surface Fe³⁺ suggests that the former is situated deeper than surface Fe^{3+/2+} redox level in BaTiO₃ NCs or that electronic coupling between surface Fe³⁺ and the self-trapped Ti³⁺ defects is weaker. Similar step-wise reductions were also observed in the recent report on n-type Fe³⁺:ZnO NCs containing both substitutional and surface Fe³⁺ species.^{9,46} Suzuki *et al.* have reported that the Fe^{3+/2+} redox level is located ~0.7-0.8 eV below the conduction band minimum in BaTiO₃ single crystals, while in SrTiO₃ the Fe^{3+/2+} level is nearly degenerate with the conduction band.⁵⁰ Our results are consistent with the Fe^{3+/2+} energy level lying at a more positive potential than the conduction band and Ti³⁺ level in both SrTiO₃ and BaTiO₃ NCs.

Several other reports show that at least some of the Fe^{3+} can be reduced to Fe^{2+} in SrTiO_3 and BaTiO_3 bulk powder and thin films under reducing conditions.⁵¹⁻⁵³ However, most of these reductions are limited to the surface and sub-surface Fe dopants and typically require elevated temperatures. Stabilization of Fe^{2+} in reduced SrTiO_3 and BaTiO_3 is also challenging as it requires additional charge compensation that could include the formation of oxygen vacancies.⁵⁴ The requirement of oxygen vacancies and effect on the lattice may be responsible for the reduction of only some of the Fe^{3+} dopants in bulk ATiO_3 . In contrast, the charge compensation during photodoping is accomplished by proton (H^+) adsorption at the surfaces, as reported in TiO_2 and other oxide NCs.^{14,55} Furthermore, photodoping can be achieved at room temperature with just UV photons and a sacrificial hole scavenger. This facile method offers a controlled and reversible reduction of every Fe^{3+} dopant throughout the entire ATiO_3 NC.

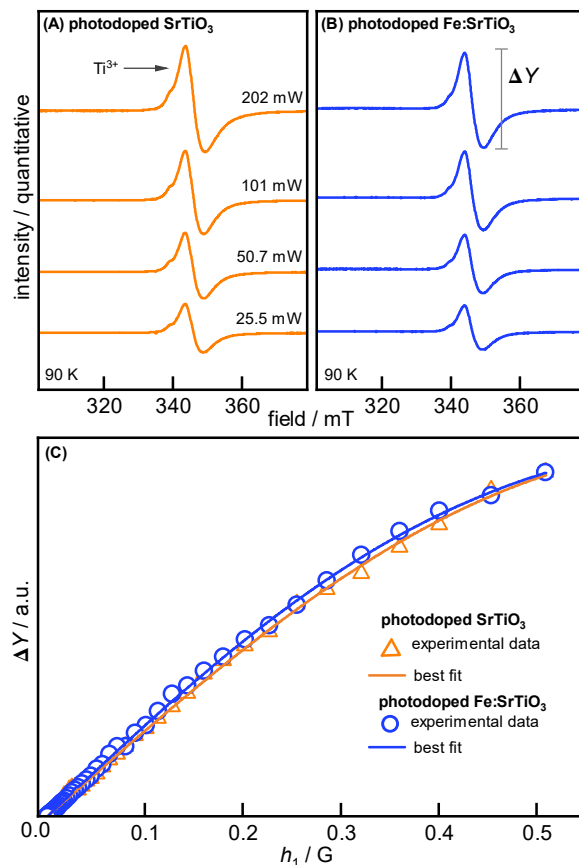


Figure 3.5 Representative CW-EPR spectra of maximally photodoped (A) SrTiO₃ and (B) Fe:SrTiO₃ NCs at 90 K at selected microwave powers. The signal intensity defined as ΔY and best fits to equation 3.3 (see Figure 3.16 and Table 3.4) for each measurement plotted as a function of microwave power (h_1) is shown in the bottom panel C.

These results provide sufficient evidence that photochemically introduced n-type carriers reduce the Fe³⁺ dopants to EPR-silent Fe²⁺ ions, in contrast to the acceleration of spin-relaxation time of Cr³⁺ dopants.²⁵ However, they do not rule out the possible cross-relaxation between Fe³⁺ and Fe²⁺ or Ti³⁺ and Fe²⁺ centers in photodoped ATiO₃ NCs. High-field EPR measurements to directly observe Fe²⁺ dopants are planned. However, with X-band CW-EPR, we were able to study the effect of Fe²⁺ on Ti³⁺ spin-dynamics. To probe this, we performed power saturation rollover experiments and linewidth analysis of Ti³⁺ signal in the absence and presence of Fe²⁺. These experiments were performed on photodoped SrTiO₃ (no Fe²⁺) and Fe:SrTiO₃ (Fe³⁺

fully reduced to Fe^{2+}) at 90 K (See Figures 3.5A-C, 3.15-3.16 and Table 3.4). No significant difference in saturation rollover behavior or T_2 time of the Ti^{3+} defects in both photodoped SrTiO_3 (2.65 ± 0.06 ns) and photodoped $\text{Fe}:\text{SrTiO}_3$ (2.50 ± 0.04 ns) was observed. The slight difference in T_2 can be attributed to variation in average Ti^{3+} concentrations in the photodoped NCs. These results suggest that the spin properties of Ti^{3+} defects are not affected by the presence of Fe^{2+} dopants in photodoped SrTiO_3 NCs.

3.4 Supporting Information

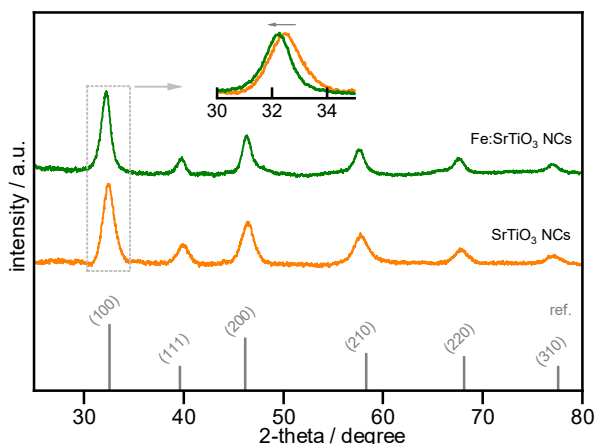


Figure 3.6 Powder X-ray diffraction pattern of as-prepared SrTiO_3 (bottom/black) and $\text{Fe}:\text{SrTiO}_3$ NCs (top/green) indexed to cubic phase of bulk SrTiO_3 with space group of $\text{Pm}\bar{3}\text{m}$.⁵⁶ The baseline was carefully subtracted from both patterns. A slight shift to lower 2θ value in $\text{Fe}:\text{SrTiO}_3$ NCs displayed in the inset is attributed to a typical lattice expansion due to the incorporation of relatively larger size Fe^{3+} (0.79 \AA) cation at cubic Ti^{4+} (0.75 \AA) site.

Table 3.1 The average lattice parameter and crystallite size calculated using Rigaku SmartLab Studio II for both samples.

Sample NCs	Lattice parameter (\AA)	Crystallite size (nm) [†]
SrTiO_3	3.9197	5.92 ± 0.42
$\text{Fe}:\text{SrTiO}_3$	3.9256	8.02 ± 0.54

[†]Crystallite size is the average size calculated from the three most intense reflections: (110), (200) and (211).

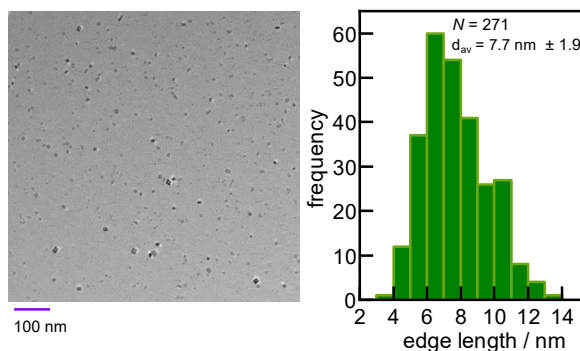


Figure 3.7 Representative transmission electron microscopy (TEM) image of as-prepared Fe:SrTiO₃ NCs deposited on 3 mm copper grid exhibiting a cubic morphology. The size distribution plot from analyzing a few hundred different crystals using ImageJ is shown in the right panel.

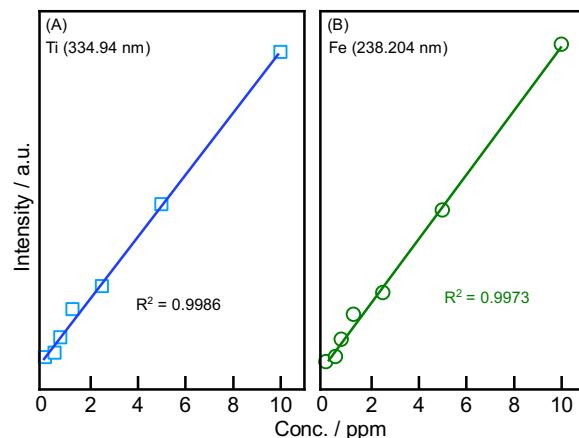


Figure 3.8 Calibration plots for Ti (left/blue) and Fe (right/green) for ICP-OES measurements using multielement QC-21 Perkin Elmer standard solution to determine percentage in Fe:SrTiO₃ NCs. In a typical sample preparation for ICP-OES analysis, solvent-evaporated NCs were digested in aqua regia (~1-2 mL) followed by the dilution in 5% HNO₃ solution.

Table 3.2 ICP-OES analysis of Fe content in Fe:SrTiO₃ NCs

Sample NCs	Nominal Fe content (%)	Fe from ICP-OES (%)
Fe:SrTiO ₃	1.0	1.70
Fe:SrTiO ₃	0.5	1.03

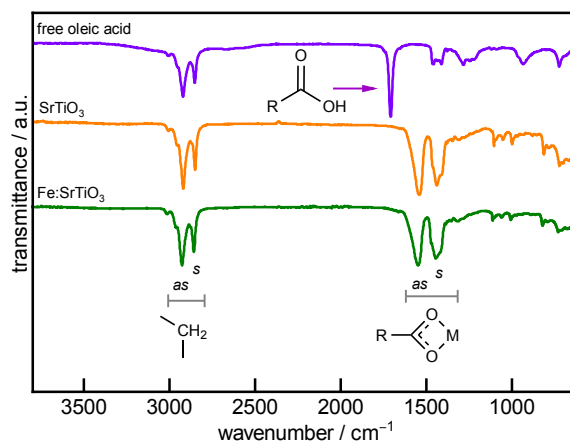


Figure 3.9 FTIR spectra of solvent evaporated SrTiO₃ (orange) and Fe:SrTiO₃ (green) NCs. The top purple spectrum belongs to free oleic acid. The pair of peaks at 1450 cm⁻¹ and 1550 cm⁻¹ corresponding to the symmetric (*s*) and asymmetric (*as*) stretches of bound carboxylate head group of oleate ligand is observed in both NCs. Aliphatic bands of oleate ligands are also present *ca.* 2900 cm⁻¹, which indicate the NCs surfaces are passivated by oleate ligands.

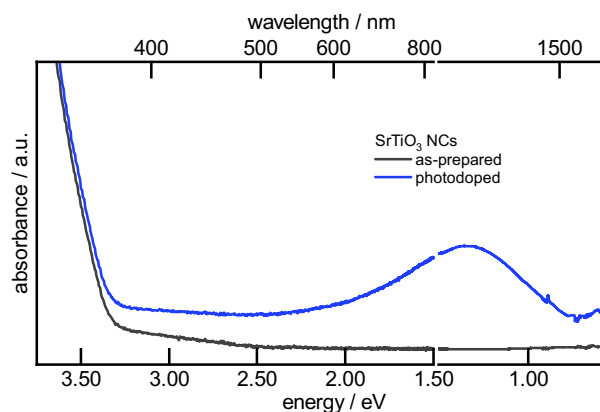


Figure 3.10 Electronic absorption spectrum of as-prepared and photodoped SrTiO₃ NCs. The full spectrum was made by splicing the spectra collected on a Cary 50 for $\lambda < 900$ nm and a Varian 670 for $\lambda > 900$ nm.

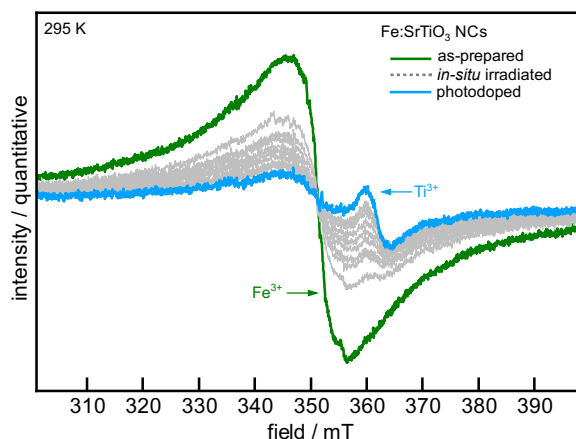


Figure 3.11 Room temperature in situ EPR spectra of Fe:SrTiO₃ NCs as a function of photodoping time. The green spectrum belongs to as-prepared NCs before turning on the Xe lamp. After turning on the lamp, the broad Fe³⁺ signal decreases in intensity and linewidth while a new feature at ~360 mT appears (grey lines). This new signal is assigned to Ti³⁺ defects. The blue spectrum shows the photodoped NCs with some residual Fe³⁺ signal which can be completely eliminated upon prolonged photodoping as discussed in the main paper.

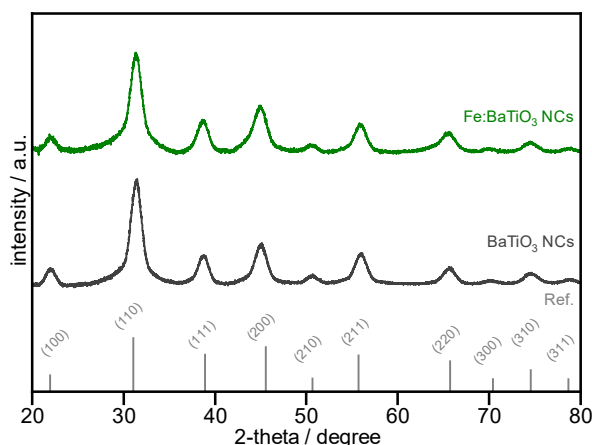


Figure 3.12 Powder X-ray diffraction pattern of as-prepared BaTiO₃ (bottom/black) and Fe:BaTiO₃ NCs (top/green) indexed to cubic phase of bulk BaTiO₃ with space group of Pm $\bar{3}$ m.⁵⁷ The average particle size calculated from peak analysis is shown below.

Table 3.3 The average lattice parameter and crystallite size calculated using Rigaku SmartLab Studio II for both samples.

Sample NCs	Lattice parameter (Å)	Crystallite size (nm) [†]
BaTiO ₃	4.0230	6.10 ± 0.4
Fe:BaTiO ₃	4.0523	5.04 ± 0.9

[†]Crystallite size is the average size calculated from the three most intense reflections: (110), (200) and (211).

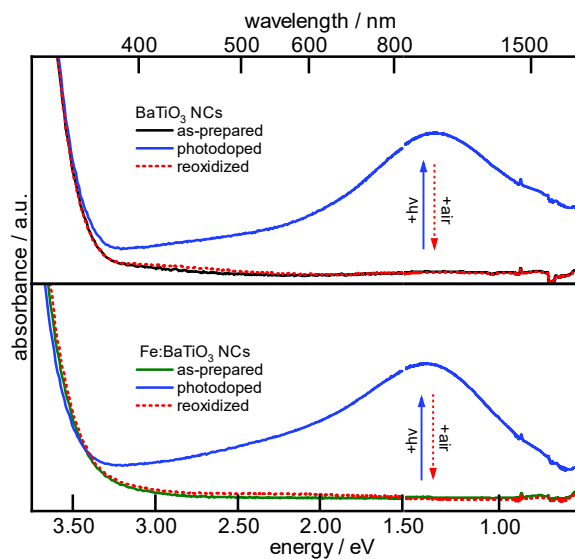


Figure 3.13 Electronic absorption spectra of an air-free solution of BaTiO_3 and Fe:BaTiO_3 NCs in hexanes before and after photodoping.

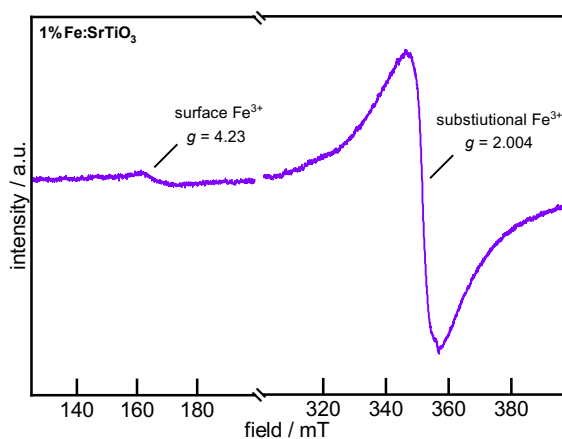


Figure 3.14 Room temperature EPR spectrum of nominally 1% Fe-doped SrTiO_3 NCs.

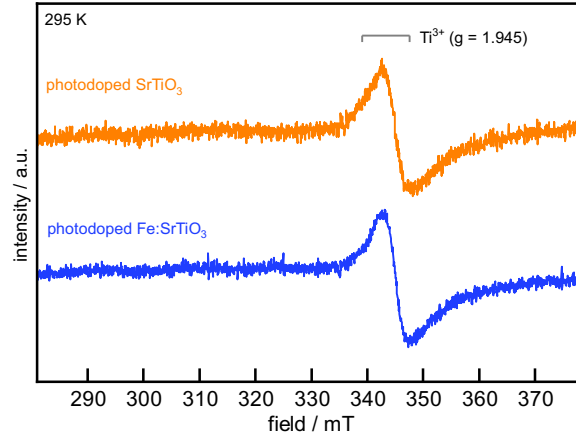


Figure 3.15 The CW-EPR spectra of collected at room temperature of maximally photodoped SrTiO₃ (top/orange) and Fe:SrTiO₃ NCs (bottom/blue) displaying the characteristics resonance at $g = 1.945$ attributed to Ti³⁺ defects. The concentration of NCs in each sample was kept the same, approximated from the combination of ICP-EOS and particle size analysis.

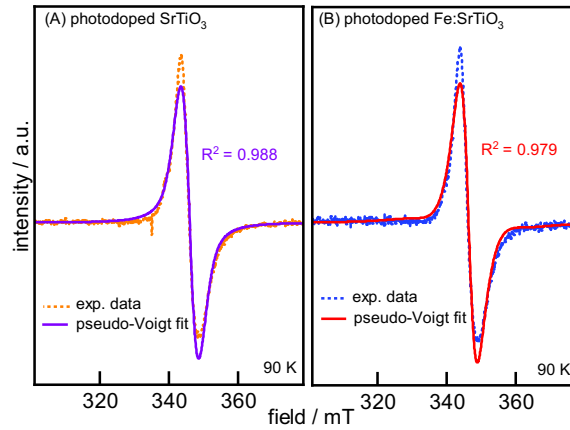


Figure 3.16 CW-EPR spectra collected well below the saturation power at 90 K of (A) photodoped SrTiO₃ and (B) photodoped Fe:SrTiO₃ NCs. The purple and red lines are best fits to pseudo-Voigt line profile in origin (Equation 3.1) to obtain the Lorentzian width (Γ_L) to calculate the spin-spin relaxation time (T_2) using equation 3.2.

$$y_0 + \left(-A_g \left(\frac{(x-x_0)}{W_g} \right)^2 \right) \exp \left(-0.5 \left(\frac{(x-x_0)}{W_g} \right)^2 \right) + \left(-2A_l \left(\frac{(x-x_0)}{W_l} \right)^2 \right) / \left((W_l^2) \left(1 + \left(\frac{(x-x_0)}{W_l} \right)^2 \right)^2 \right) \quad (\text{eq. 3.1})$$

$$T_2 = 2/(\gamma_e \Gamma_L) \quad (\text{eq. 3.2})$$

Where γ is the electron gyromagnetic ratio ($1.760859708 \times 10^{11} \text{ s}^{-1}\text{T}^{-1}$).

Equation 3.3 shown below was used to fit the power saturation data shown in Figures 4A-C in the main text to determine the term P_2 which is the product of the spin-lattice ($1/T_1$) and spin-spin relaxation rate ($1/T_2$) as shown in equation 3.4. Where c is a scalar and ϵ is a measure of line homogeneity.

$$\Delta Y = c \cdot h_1 \left(1 + \frac{(h_1)^2}{P_2} \right)^{-\epsilon} \quad (\text{eq. 3.3})$$

$$P_2 = (\gamma_e^2 T_1 T_2)^{-1} \quad (\text{eq. 3.4})$$

Table 3.4 Spin-spin relaxation time (T_2) calculated from Lorentzian width (Γ_L) from fitting EPR spectra of both samples to first derivative pseudo-Voigt line profile as displayed in figure 3.16.

Sample NCs	Γ_L (G)	T_2 (ns)
SrTiO ₃	42.88 ± 1.11	2.65 ± 0.06
Fe:SrTiO ₃	45.75 ± 0.62	2.50 ± 0.04

3.5 References

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CHAPTER 4

CONTROLLING THE VALENCE STATE OF IRON DOPANTS IN COLLOIDAL TITANIUM DIOXIDE NANOCRYSTALS

4.1 Introduction

The intentional incorporation of impurities or dopants in semiconductors is fundamental to manipulate the properties that render them useful for various applications.¹ The example of such impurities in semiconductors includes the magnetic dopants such as Mn^{2+} and Ni^{2+} in ZnO to activate the room temperature ferromagnetism for their potential use in spin-based applications^{2,3} and the Cr dopants in TiO_2 to improve its optical properties.⁴ The changes in the electronic structure of semiconductors are very sensitive to the individual dopants properties such as the energy level alignment, oxidation states, and their speciation within the host lattices. As an example, the substitution of Al^{3+} in ZnO acts as an n-type dopant while the same dopant produces p-type carriers in SnO_2 .^{5,6}

Our understanding of the dopant-induced properties particularly in free-standing nanoscale semiconductors is incomplete because of the unavailability of robust synthetic methods to prepare doped nanocrystals and their complicated surface and defect chemistries.⁷ Moreover, many of the transition metal dopants exist in multiple oxidation states within the host lattices.^{8,9} The origin of the dopant-induced properties oftentimes remains elusive as each oxidation state imparts unique changes to the electronic structure. For example, the correlation between the mixed-oxidation states of redox-active dopants such as $\text{Mn}^{3+/2+}$ and $\text{Fe}^{3+/2+}$ and the room temperature ferromagnetic ordering in oxide lattices has been a long-standing

debate.¹⁰⁻¹² Controlling the properties of individual dopants within the host lattice is, therefore, a key to understand the underlying mechanisms for the dopant-induced changes in semiconductors.

In our recent report, we demonstrated that the oxidation state of Fe dopants in SrTiO₃ and BaTiO₃ nanocrystals (NCs) can be reversibly modulated between Fe³⁺ and Fe²⁺ using a photodoping method.¹³ However, these changes in the dopant oxidation state required EtOH as sacrificial reductant and UV photons under extremely inert conditions. The traces of air can make the photodoping process completely ineffective. Here, we present an alternative method to control the oxidation state of Fe dopants in TiO₂ NCs under ambient conditions without requiring any sacrificial reductant or UV photons. We prepared and characterized Fe³⁺-doped TiO₂ colloidal NCs using a solvothermal method. The spectroscopic analysis shows that Fe³⁺ dopant is substituted at Ti⁴⁺ site in as-prepared NCs. The introduction of excess electrons using the photodoping method in as-prepared NCs shows a reversible electron trapping at Fe³⁺ dopants converting it into Fe²⁺. This observation consistent with our previous report on Ti-based oxide NCs¹³ helped to build the electronic structure of Fe-doped TiO₂ NCs. We further extended this work to incorporate the air-table excess electrons by co-doping the NCs with aliovalent Nb⁵⁺ dopants. The substitutional Nb⁵⁺ dopants add n-type carriers in the TiO₂ NCs which are stable under the ambient condition as confirmed through electronic absorption and electron paramagnetic resonance (EPR) spectroscopies. The presence of n-type carriers is correlated with the complete disappearance of Fe³⁺-related EPR signal that we attribute to the reduction of Fe³⁺ to EPR-silent Fe²⁺.

4.2 Experimental

4.2.1 Materials

Titanium(iv) *n*-butoxide ($\text{Ti}(\text{OBu})_4$, 99%, Acros Organics), oleylamine (>50%, TCI America), oleic acid (>95%, Fisher Chemical), iron nitrate nonahydrate ($\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$, 99.99% Sigma Aldrich), niobium(v) chloride (NbCl_5 , 99.95%, Alfa Aesar), ethanol (200 proof, PHARMCO-AAPER), QC standard, 21 elements (5% HNO_3 , PerkinElmer), nitric acid (Certified ACS Plus – Safe Cote, Fisher Chemical), hydrochloric acid (Certified ACS Plus, Fisher Chemical), toluene (optima, Fisher Chemicals) and hexanes (optima, Fisher Chemicals) were all used as received.

4.2.2 Synthesis of Nanocrystals

The synthesis of colloidal TiO_2 NCs was carried out by a solvothermal method. In a typical preparation, 1.5 mmol of $\text{Ti}(\text{OBu})_4$ was mixed with 7.5 mmol of oleic acid and 7.5 mmol of oleylamine in 6 mL of 99% EtOH in water (v/v). The obtained mixture was then transferred to a 45 mL Teflon-lined autoclave and heated to 200 °C for 18 h. After the synthesis, the autoclave was cooled down to room temperature and the solid product was collected, washed with ethanol several times, and suspended in non-polar solvents such as toluene to produce a turbid solution. Finally, a transparent layer containing NCs was extracted by centrifuging (5 min at 4000 rpm) the cloudy solution. The synthesis of nominally 0.2% Fe-doped TiO_2 NCs was carried out by adding 30 μL of 0.1 M aqueous solution of the $\text{Fe}(\text{NO}_3)_3$. Similarly, for the preparation of 8% Nb-doped TiO_2 NCs, 32 mg of NbCl_5 was added before closing the autoclave.

4.2.3 Physical Characterization

All measurements presented below were collected on either undoped or 0.2% Fe-doped TiO_2 NCs at room temperature unless specified otherwise. Electronic

absorption and emission spectra were collected at room temperature on colloidal NCs in an air-tight quartz cuvette on a Cary 50 Bio and Cary Eclipse, respectively. Near-IR absorption measurements were collected on a Varian 670 FT spectrometer equipped with an InGaAs detector and near-IR quartz beam splitter using a 1-mm pathlength quartz cuvette. The functional groups and surface chemistry analysis of solvent-free NCs were carried out on a Bruker Alpha-P Fourier transform infrared (FTIR) spectrophotometer equipped with an attenuated total reflectance (ATR) platinum diamond optic. EPR measurements were carried out on a continuous wave (CW) Bruker Eleksys-500 fitted with Super High QE X-band cavity (ER4123SHQE) and cryostat (ESR-900, Oxford) for low-temperature measurements. For the photodoping experiments, a 75 W xenon lamp was used (60000 Q series, Oriel Corporation). Transmission electron microscopy (TEM, JEOL 2000FX) images were measured from freshly prepared NCs drop-casted onto carbon-coated (3nm) copper grids (CF400-CU-50, Electron Microscopy Sciences). The average size distribution and standard deviation were generated by analyzing a few hundred particles from TEM images using ImageJ software. Powder X-ray diffraction patterns were collected in a Bragg-Brentano configuration and Cu K- source (Rigaku Smart Lab SE). Metal ion stoichiometry was measured on dissolved ensembles of NCs and measured by inductively coupled plasma-optical emission spectrometry (ICP-OES, Perkin Elmer Optima 4300 DV).

4.3 Results and Discussion

The synthesis of nominally 0.2% Fe-doped TiO₂ NCs (to referred as Fe:TiO₂ hereafter) was carried out using a solvothermal method by modifying a literature

procedure¹⁴ as described in experimental section. This synthesis employs water vapor to hydrolyze the $\text{Ti}(\text{OBu})_4$ precursor over 24 h that allows the successful dopant incorporation. The as-prepared NCs are soluble in non-polar solvents such as toluene to produce a transparent colloidal solution. The powder X-ray diffraction (PXRD) pattern of the as-prepared product is indexed to anatase crystal structure of TiO_2 (I4₁/amd) with no detectable secondary phases (see SI Figure 4.5).¹⁵ The small increase in the lattice parameters of $\text{Fe}:\text{TiO}_2$ compared to undoped TiO_2 NCs (see SI Table 4.1) is attributed to the substitution of a larger size Fe^{3+} (0.79 Å) dopants at Ti^{4+} (0.74 Å) sites consistent with other spectroscopic evidence discussed below.¹⁶ The TEM images of TiO_2 and $\text{Fe}:\text{TiO}_2$ NCs (see SI Figures 4.6 and 4.7) show an average size of 10.4 ± 2.2 and 7.45 ± 1.5 nm, respectively. The size of NCs measured from TEM images agrees well with the crystallite size calculated from PXRD analysis shown in Table 4.1.

The electronic absorption spectra of as-prepared TiO_2 (black) and $\text{Fe}:\text{TiO}_2$ NCs (green) are shown in Figure 4.1A-B. Both samples are dominated with the bandgap absorption above ~ 3.25 eV consistent with bandgap energy of single crystal of TiO_2 .¹⁷ However, $\text{Fe}:\text{TiO}_2$ NCs show an absorption smearing out of the band edge. This absorption feature has been previously observed in Fe-doped TiO_2 nanoparticles and is assigned to the charge transfer transitions from the valence band to the Fe^{3+} impurity levels located below the conduction band.¹⁸ To probe the local environment of Fe^{3+} dopants in TiO_2 , we employed the X-band EPR spectroscopy. The room temperature EPR spectrum of as-prepared $\text{Fe}:\text{TiO}_2$ NCs shown in Figure 4.2 (green) displays a broad anisotropic signal centered around 350 mT. Although the shape and

the g -value for this EPR signal agree with the substitutional Fe^{3+} dopants in anatase,¹⁹⁻
²² we observe much resolved features spread around the central peak at g -value = 1.992. To confirm the identity of all these features, we simulated the EPR spectrum of Fe^{3+} ($S = 5/2$) in TiO_2 with both the axial (D) and rhombic (E) components to the zero-field splitting (ZFS) using the spin-Hamiltonian (eq. 4.1) shown below in EasySpin software.

$$H = g\beta BS + D \left(\hat{S}_z^2 - \frac{1}{3}(S^2 + S) \right) + E(\hat{S}_x^2 - \hat{S}_y^2) + \sum_k \sum_{q=-k}^k B_k^q O_k^q \quad (\text{eq. 4.1})$$

where g is the Landé g -factor, β is the Bohr magneton and B represents the applied magnetic field. The last term in the equation represents the fine structure parameters (a and F). With the initial parameters for Fe-doped TiO_2 taken from the literature,²³ the simulated spectrum is displayed as a pink dotted line in Figure 4.2. A good agreement between experimental and simulated EPR spectrum identifies this broad signal including the satellite features spanning through 300-400 mT as a substitutional Fe^{3+} dopant with prominent axial ZFS ($D = 151.758 \times 10^{-4} \text{ cm}^{-1}$) and a negligible rhombic anisotropy ($E = 7.35 \times 10^{-10} \text{ cm}^{-1}$). The fine structures around the central peak are originating from cubic parameters a (-63.5026 MHz) and F (-529.5953 MHz). No evidence of surface Fe^{3+} species was detected in the low-field region of EPR spectrum collected on a few different samples (see SI Figure 4.8).

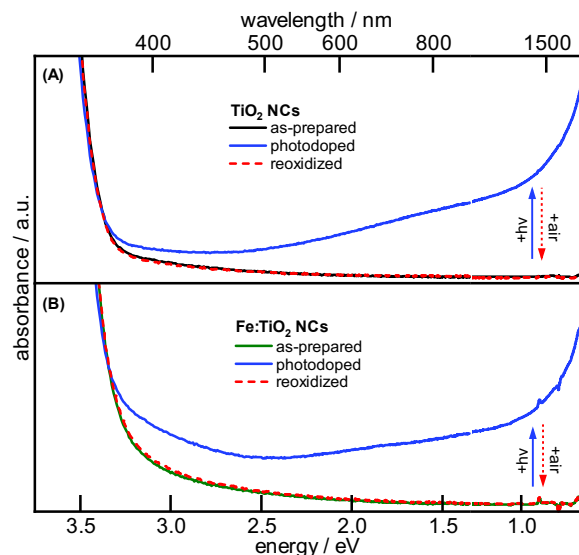


Figure 4.1 Electronic absorption spectra of (A) TiO_2 and (B) $\text{Fe}:\text{TiO}_2$ NCs before and after the photodoping. The full absorption spectra were generated by splicing the data collected on a Cary 50 for energy > 1.2 eV and a Varian 670 for energy < 1.2 eV.

To probe the effect of excess electrons on the substitutional Fe^{3+} dopants, photodoping of as-prepared samples was performed. The detailed procedure for the photodoping experiments is adopted from our previous report.¹³ Briefly, UV irradiation of air-free solution of NCs creates electron-hole pairs. The rapid hole quenching with EtOH as a sacrificial reductant leaves excess electrons in the NCs. The electronic absorption spectra displayed in Figure 4.1A-B show the appearance of a broad absorption feature tailing through the band-edge with photodoping both TiO_2 and $\text{Fe}:\text{TiO}_2$ NCs. This characteristic broad absorption that maximizes in the mid-IR along with the deep blue coloration of photodoped samples (see SI Figure 4.9) has been previously observed in n-type TiO_2 nanoparticles and is attributed to the localized surface plasmon resonance (LSPR).^{24,25} A clear blue-shift in the band-edge absorption with photodoping due to the Moss-Burstein effect is also indicative of the free electrons in the conduction band (e_{CB}^-). It is worth noticing that in addition to the LSPR-related broad absorption, a weak feature around ~ 1.75 eV (800 nm) also

appears with the photodoping. The energy of this feature is consistent with the trapped electrons at Ti^{3+} sites as reported by Mayer and co-workers in photodoped TiO_2 amorphous nanoparticles.²⁶ The amorphous TiO_2 does not have any well-defined bands. Therefore, added electrons did not produce any free electrons or surface plasmon but Ti^{3+} trap sites in photodoped TiO_2 nanoparticles in that study. However, TiO_2 employed in current work is phase pure crystalline anatase and possess well-defined valence and conduction band as depicted from band edge absorption. Valentin et al., found that quasi-localized or localized Ti^{3+} states in crystalline n-type TiO_2 thin films are very close in energy (~ 0.1 eV) to the delocalized electrons that can be thermally excited to the conduction band.²⁷ We speculate that a fraction of electrons in photodoped NCs creates thermally-detrapped Ti^{3+} defects, and the charge transfer transition from these Ti^{3+} trap states to the conduction band gives rise to this small absorption feature at ~ 1.75 eV. The electronic absorption spectra taken periodically during the whole photodoping process on both the TiO_2 and Fe:TiO_2 NCs show the immediate appearance of the broad plasmonic absorption followed by the weak feature at ~ 1.75 eV (see SI Figures 4.10 and 4.11). This observation suggests that the majority of photoexcited electrons are delocalized in the conduction band while that a fraction of them creates quasi-localized Ti^{3+} defects with increased photodoping time.

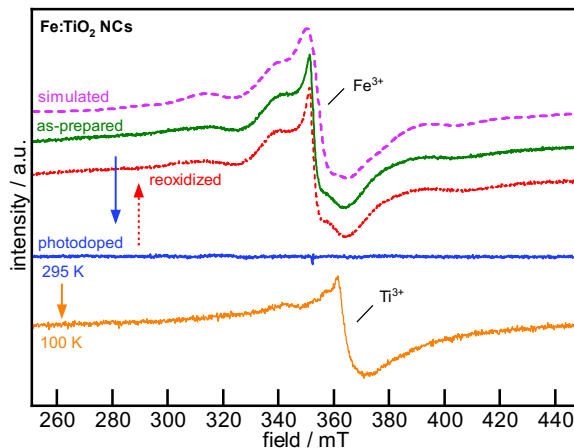


Figure 4.2 The EPR spectrum of as-prepared NCs (green) compared to the simulated (pink/dotted), the photodoped NCs at room temperature (blue) and the 100 K (orange), and the EPR spectrum of reoxidized NCs represented with the red dotted line.

We further employed EPR spectroscopy to explore the effect of added electrons on Fe^{3+} dopants. The room temperature EPR spectra of as-prepared and photodoped $\text{Fe}:\text{TiO}_2$ NCs are shown in Figure 4.2. As discussed in the earlier section, the as-prepared sample exhibits only a broad signal associated with the substitutional Fe^{3+} dopants. With a short photodoping time (~ 20 min), the Fe^{3+} -related EPR signal is completely disappeared. This observation persists down to 100 K where a new axial signal at ~ 365 mT is detected. This broad and poorly resolved anisotropic signal is associated with the formation of Ti^{3+} (d^1) defects consistent with absorption measurements and the literature.^{24,26} The Ti^{3+} signal is observed only at cryogenic temperature owing to its short spin relaxation times. The increased spin-spin relaxation pathways in the photodoped NCs with excess electrons cause the broadening of this Ti^{3+} signal. The disappearance of the Fe^{3+} signal with the simultaneous appearance of Ti^{3+} defects is consistent with our recent report on Fe dopants in SrTiO_3 and BaTiO_3 NCs.¹³ The Fe^{3+} dopants in TiO_2 have been reported to create electron trap states below the conduction band.²⁸⁻³⁰ We propose that photochemically added electrons reduce these trap states to EPR-silent Fe^{2+} as

illustrated in Figure 4.3. The photoexcitation of charge carriers followed by the rapid hole quenching with EtOH raises the Fermi level into the conduction band. The substitutional Fe^{3+} dopants act as trap states below the conduction band and get reduced to Fe^{2+} during photodoping. Upon reopening the sample to air, these changes are reversed and the Fe^{3+} signal is quantitatively recovered.

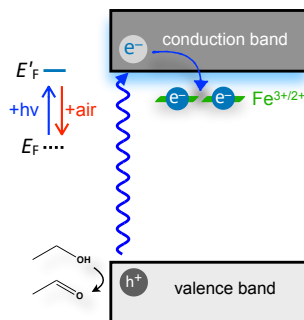


Figure 4.3 Reversible electron trapping at the $\text{Fe}^{3+/2+}$ in $\text{Fe}:\text{TiO}_2$ colloidal NCs with photodoping.

Although the photodoping method offers a tunable and reversible control of the valance state of Fe^{3+} dopants, the whole process is very sensitive to the presence of air. The traces of oxygen can make the photodoping process completely ineffective. The photochemically added electrons are typically stabilized by the adsorption of protons at the surface of NCs.^{26,31} Upon exposing the sample to air, these electrons can easily react with molecular oxygen near the surfaces to produce superoxide radicals which we also observed in our previous report on Cr-doped SrTiO_3 NCs.³² This instability of the added electrons limits the exploration of the photodoped NCs, and thus requires alternative methods. The incorporation of stable carriers in TiO_2 has been studied using anaerobic annealing methods and the aliovalent doping such as Nb^{5+} and Mo^{5+} etc., at substitutional Ti^{4+} sites.^{33,34} Although both approaches produce free electrons, we find post-synthetic annealing a least suitable method as it requires extreme temperatures and may affect the morphology and crystalline phase

of the NCs. Preparation of n-type TiO₂ NCs on the other hand using aliovalent doping can be achieved by simply adding appropriate metal salt during the synthesis without requiring any additional steps.

The synthesis of nominally 0.2% Fe 8% Nb (high enough to produce excess electrons) co-doped TiO₂ (to referred as (Fe,Nb):TiO₂ hereafter) NCs was carried out following the solvothermal procedure described in experimental section. The as-prepared NCs exhibit the crystal structure of anatase evident from the PXRD pattern (see SI Figure 4.5). Although no detectable secondary phases are observed, there is a significant elongation along the z-axis as summarized in Table 4.1. Such elongation has been previously observed in Nb-doped TiO₂ NCs and is related to changes in the morphology of NCs from cubic to “peanut” shape.²⁵ However, the TEM images in our work (see SI Figure 4.12) do not show any significant changes in the morphology of (Fe,Nb):TiO₂ NCs compared to undoped TiO₂ or Fe:TiO₂ NCs. Alternatively, we propose that a fraction of the added electrons with Nb⁵⁺ doping can create localized Ti³⁺ defects like what we observed in photodoped TiO₂ NCs. The Ti³⁺ with electrons primarily occupying the ground state d_{xz} and d_{yz} orbitals can contribute to the elongation along the z-axis in the crystal structure.

The electronic absorption spectrum of colloidal solution of as-prepared (Fe,Nb):TiO₂ NCs exhibits a broad absorption feature tailing from the band-edge into the mid-IR region (Figure 4.4A) similar to photodoped NCs. We assign this characteristic broad absorption to the LSPR consistent with the other reports on Nb-doped TiO₂ NCs.^{25,34} There is also clear evidence of a blue shift in the band-edge energy due to the Moss-Burstein effect which further confirms the incorporation of

e_{CB}^- with Nb-doping (Figure 4.4A inset). In contrast to the photodoping method, we do not observe any additional absorption feature related to the Ti^{3+} defects at ~ 1.75 eV. Despite the absence of this feature, we cannot rule out the existence of Ti^{3+} defects because the plasmonic absorption typically have a large oscillator strength that can occlude the absorption feature from trapped electrons. Nevertheless, these results provide convincing evidence that the substitution of Nb^{5+} adds air-stable electrons and majority of which are free carriers in the conduction band of TiO_2 NCs.

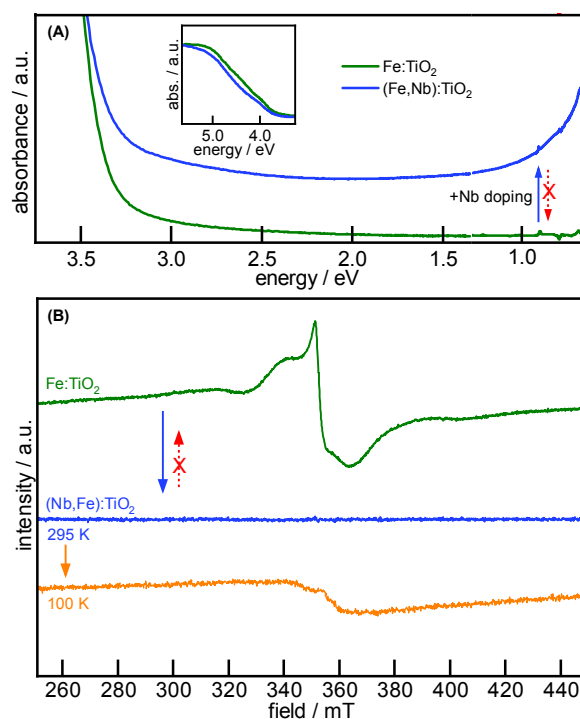


Figure 4.4 The EPR spectra of as-prepared Fe:TiO₂ (green/top) and 0.2% Fe 8% Nb co-doped TiO₂ NCs collected at room temperature (blue/middle) and 100 K (orange/bottom).

To exploit the effect of the added electrons on Fe^{3+} dopants, EPR measurements were carried out. The room temperature EPR spectrum of (Fe,Nb):TiO₂ compared to Fe:TiO₂ NCs shows the complete disappearance of Fe^{3+} related EPR signal (see Figure 4.4B). This observation persists down to 100 K where a weak broad absorption feature corresponding to Ti^{3+} defects is observed. We attribute the broadening of the

Ti³⁺ signal to the increased spin-spin interactions in these heavily Nb⁵⁺ doped TiO₂ NCs. The controlled experiments on nominally 2% and 5% Nb-doped TiO₂ NC further support the broadening of Ti³⁺ signal with increased dopant concentration as shown in Figure 4.13. The disappearance of Fe³⁺ signal both at room temperature and 100 K consistent with the absorption measurement in Figure 4.4A suggests that excess carriers introduced through aliovalent Nb⁵⁺ doping reduces the Fe³⁺ dopant to Fe²⁺. This reduction process is completely *irreversible and stable* under ambient conditions compared to photodoping method. This observation opens the avenue for further exploration of mixed-valence states of Fe³⁺ and other redox-active dopants in oxide lattices.

4.5 Supporting Information

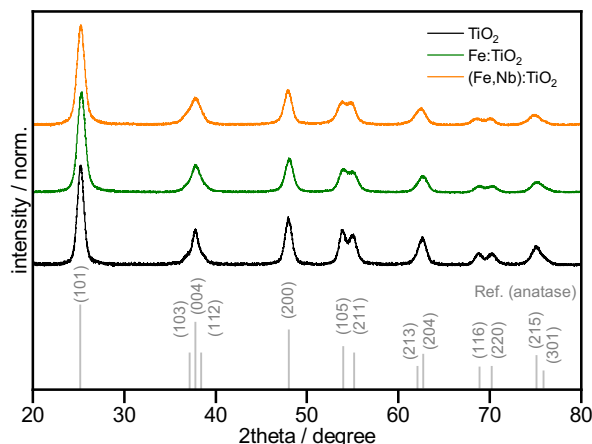


Figure 4.5 Powder x-ray diffraction patterns of TiO₂ (black/bottom), Fe:TiO₂ (green/middle) and (Fe,Nb):TiO₂ (orange/top) NCs indexed to anatase.

Table 4.1 The average lattice parameter and crystallite size calculated using Rigaku SmartLab Studio II on diffraction data shown in Figure 4.5.

Sample NCs	Lattice parameters (Å)			Crystallite size (nm) [†]
	a	b	c	
TiO ₂	3.7865	3.7865	9.5067	10.03 ± 0.024
Fe:TiO ₂	3.7902	3.7902	9.5158	8.63 ± 0.113
(Fe,Nb):TiO ₂	3.8009	3.8009	9.5430	8.93 ± 0.431

†Crystallite size is the average size calculated from (101) and (200) reflections.

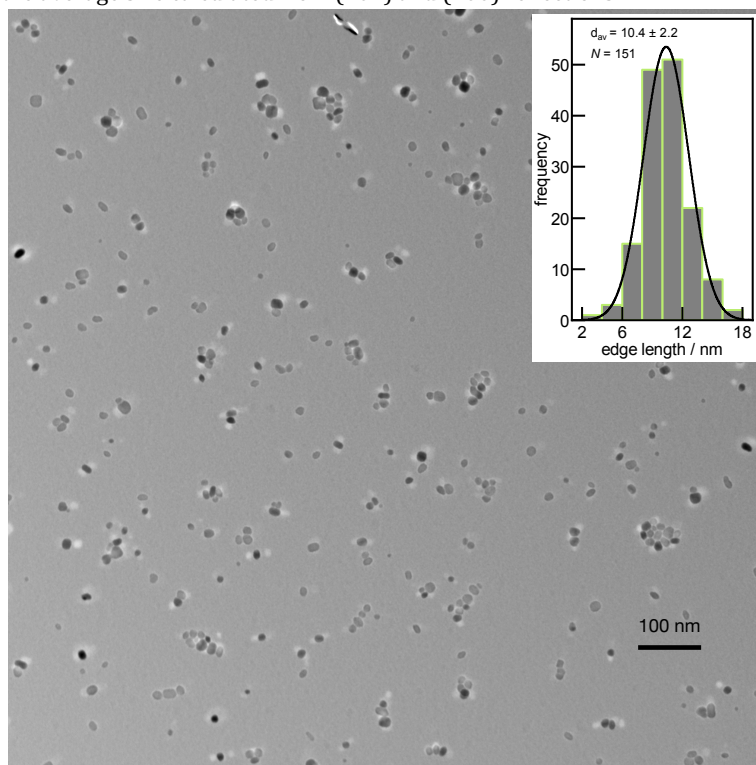


Figure 4.6 Transmission electron microscopy (TEM) image of as-prepared TiO₂ NCs and the size distribution plot from analyzing crystals using ImageJ software.

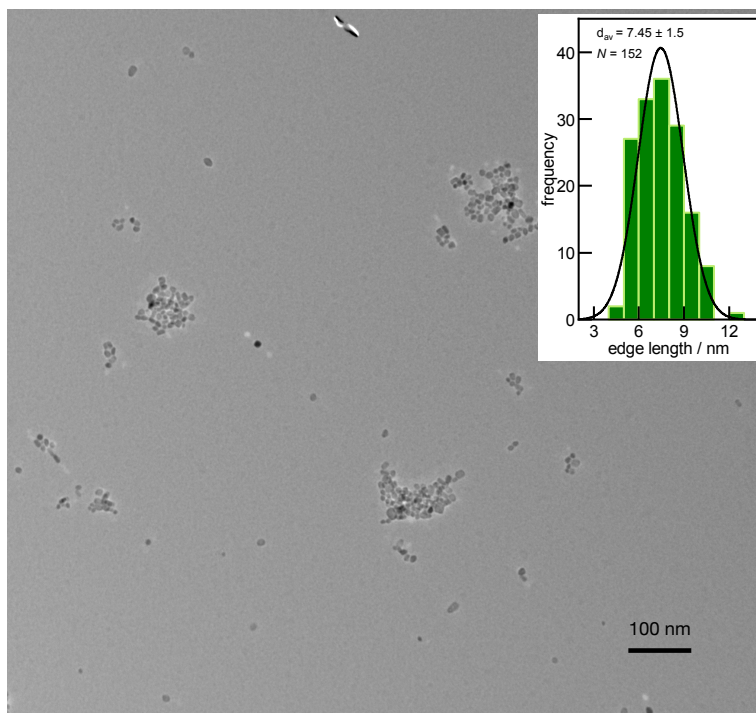


Figure 4.7 The TEM image and the size distribution plot of as-prepared Fe:TiO₂ NCs.

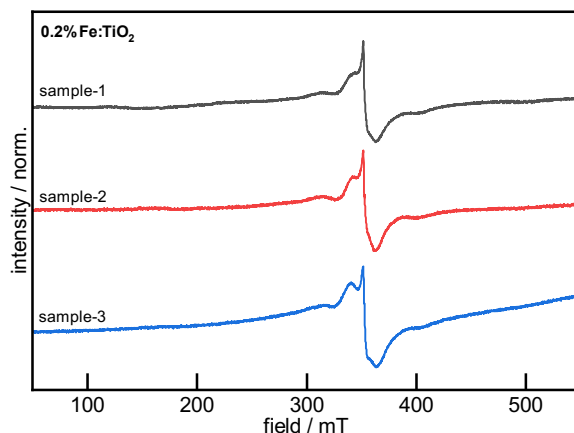


Figure 4.8 EPR spectrum of as-prepared Fe:TiO₂ NCs collected with a wide field range. All the samples were prepared under identical conditions to validate the reproducibility of EPR spectrum. No baseline correction was performed on the data.

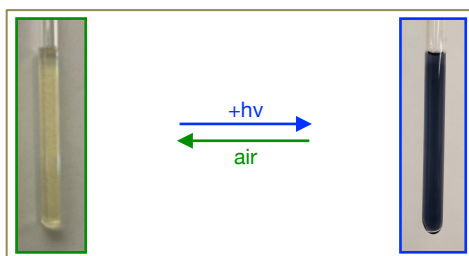


Figure 4.9 Photographs of air-free solution of Fe:TiO₂ NCs before (green/left) and after photodoping (blue/right) showing a clear change in color.

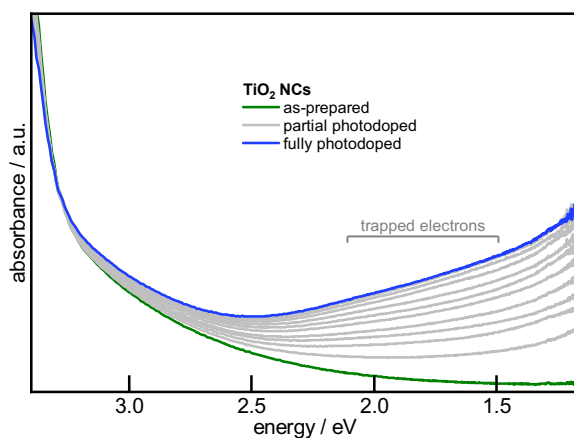


Figure 4.10 UV-Vis spectra of as-prepared (green), partially photodoped (grey line collected at various stages) and fully photodoped (blue) TiO₂ NCs.

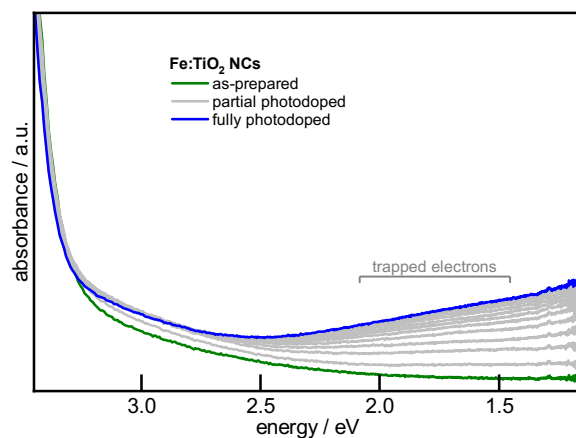


Figure 4.11 UV-Vis spectra of as-prepared (green), partially photodoped (grey line collected at various stages) and fully photodoped (blue) Fe:TiO₂ NCs.

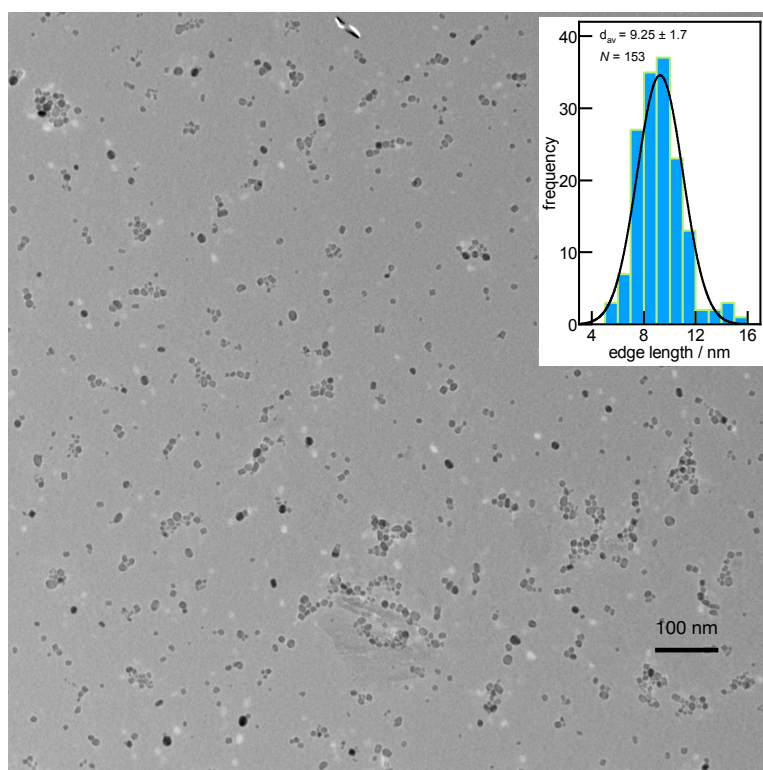


Figure 4.12 The TEM image and the size distribution plot of as-prepared (Nb,Fe):TiO₂ NCs.

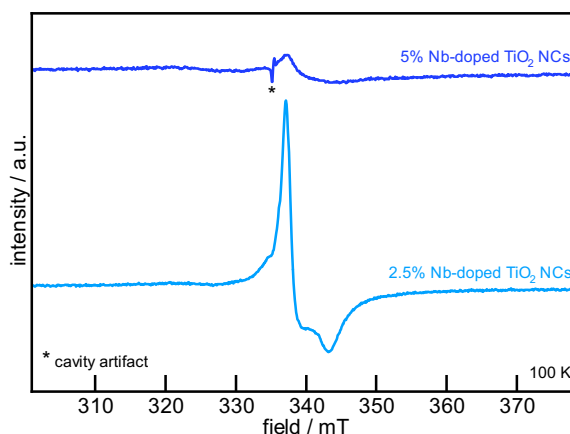


Figure 4.13 EPR spectra of nominally 2.5% Nb (light blue/bottom) and 5% Nb-doped TiO₂ NCs (dark blue/top) collected at 100 K.

4.6 References

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CHAPTER 5

EFFECT OF N-TYPE CARRIERS ON SUBSTITUTIONAL CHROMIUM DOPANTS IN TI(IV)-BASED OXIDE NANOCRYSTALS

5.1 Introduction

The introduction of paramagnetic defects into semiconducting materials has recently attracted a great deal of attention for developing the field of quantum computing – where a spin degree of freedom is exploited in addition to electronic charge to create new functionalities.^{1,2} The advancement of magnetic semiconductors has succeeded in creating intrinsic defects such as vacancies^{3,4} and doping the paramagnetic impurity ions into semiconducting host-lattices.⁵⁻⁷ The coexistence of magnetism and semiconducting properties leads to exchange interactions between localized magnetic moments and the charge carriers resulting in new magneto-optical properties.^{8,9} Similarly, multiferroic materials that exhibit ferroelectricity and magnetism are also widely investigated for various spin-based applications.¹⁰ Despite the substantial progress made over the past few decades, the quest for new materials with unique magnetic properties, particularly at room temperature, continues.

Barium titanate (BaTiO_3), a ternary oxide perovskite, has emerged as one of the prime candidates for spintronics owing to its multifunctional properties.¹¹ In addition to its intrinsic high dielectric constant and ferroelectricity, magnetism by substituting paramagnetic ions at B-site or n-type conductivity by partial reduction of Ti^{4+} to Ti^{3+} has also been well studied.^{12,13} Though significant progress has been made to achieve various properties, the interaction between paramagnetic dopants

and n-type charge carriers in a single host-lattice, particularly at nanoscale, remains elusive.

We recently reported that the introduction of n-type carriers could control the spin-relaxation time of Cr^{3+} dopants in SrTiO_3 .¹⁴ The discovery of this tunability over spin properties has led to our current work to explore such spin interactions in other Ti-based lattices to probe the role of host crystal structure and the site symmetry. Herein we report the synthesis and characterization of colloidal Cr^{3+} -doped BaTiO_3 NCs using a hydrothermal method. The dopant-specific spectroscopic studies show that dopant exists in the Cr^{3+} valence state occupying the octahedral Ti^{4+} site. The post-synthetic introduction of n-type charge carriers using the photodoping method shows that these carriers are localized at Ti^{3+} defects creating a small-polaron. Furthermore, the presence of two paramagnetic centers (Cr^{3+} and Ti^{3+}) in the same lattice structure shows a spin cross-relaxation phenomenon where fast-relaxing Ti^{3+} accelerates the Cr^{3+} spin-relaxation time. We further validate the existence of such spin processes in colloidal Cr-doped TiO_2 NCs. The added electrons in TiO_2 are delocalized in the conduction band yet show spin cross-relaxation with substitutional Cr^{3+} dopants.

5.2 Experimental

5.2.1 Materials

Barium hydroxide octahydrate ($\text{Ba}(\text{OH})_2 \cdot 8\text{H}_2\text{O}$, 99%, Alfa Aesar), titanium(IV) bis(ammonium lactate)dihydroxide (TALH, 50% in water, Alfa Aesar), sodium hydroxide (NaOH, ACS Certified, Fisher Chemical), Titanium(IV) *n*-butoxide ($\text{Ti}(\text{OBu})_4$, 99%, Acros Organics), oleylamine (>50%, TCI America), oleic acid (90%,

Fisher Chemicals), chromium nitrate nonahydrate ($\text{Cr}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$, 99.99% Sigma Aldrich), ethanol (200 proof, PHARMCO-AAPER), 21 elements standard (5% HNO_3 , PerkinElmer), hydrazine hydrate (100% Acros Organics), nitric acid (HNO_3 , Certified ACS Plus – Safe Cote, Fisher Chemicals), hydrochloric acid (HCl , Certified ACS Plus, Fisher Chemical), toluene (optima, Fisher Chemicals) and hexanes (optima, Fisher Chemicals) were used as received.

5.2.2 Synthesis of Nanocrystals

Pure and Cr-doped BaTiO_3 nanocrystals. The Synthesis of colloidal BaTiO_3 NCs was also carried out by a hydrothermal method reported in our recent study. In a typical preparation, 1.5 mmol of each TALH and $\text{Ba}(\text{OH})_2$ were dissolved in 24 mL distilled water, followed by 6 mL of NaOH (5 M). The reaction solution was then transferred to a 45-mL Teflon-lined autoclave, and then hydrazine hydrate (5 mmol), oleylamine (6 mmol), and oleic acid (6 mmol) were added. The pH of the reaction solution was 12. The sealed autoclave was placed in custom-made aluminum block housing that was heated to 215 °C and constantly stirred for 24 h using a stirring hotplate. After the synthesis, the autoclave was cooled to room temperature. The solid product was collected, washed with EtOH, and dissolved in nonpolar solvents such as hexanes to produce a transparent solution. To synthesize nominally 0.1% Cr-doped BaTiO_3 colloidal NCs, 1.5 mL of 0.001 M aqueous solution of $\text{Cr}(\text{NO}_3)_3$ was added before adding the NaOH solution.

Pure and Cr-doped TiO_2 NCs. Synthesis of colloidal TiO_2 NCs was carried out by a modified solvothermal method. In a typical preparation, 1.5 mmol of $\text{Ti}(\text{O}i\text{Bu})_4$ were mixed with 7.5 mmol of oleic acid and 7.5 mmol of oleylamine in 1.1 mL of ethanol.

The obtained mixture was then transferred to a 50 mL Teflon-lined autoclave containing 6 mL of 98% ethanol in water (v/v) and heated to 200 °C for 18 hours. After the synthesis, the autoclave was cooled down to room temperature, and the solid product was collected, washed with ethanol several times, and suspended in nonpolar solvents such as toluene to produce colloidal solutions. For the synthesis of 0.1% Cr-doped TiO₂, 1.5 mL of 0.001 M aqueous solution of Cr(NO₃)₃ was added before closing the autoclave.

5.2.3 Physical Characterization

All measurements were collected at room temperature and either on undoped or 0.1% Cr-doped BaTiO₃ and undoped or 0.1% Cr-doped TiO₂ NCs unless specified otherwise. Electronic absorption and emission spectra were obtained on colloidal suspensions in an air-tight quartz cuvette with 1-cm pathlength on Cary 50 Bio and Cary Eclipse, respectively. Near-IR absorption measurements were collected on colloidal solutions in a custom-made 1-mm pathlength quartz cuvette on Varian 670 Fourier transform infrared (FTIR) spectrometer equipped with InGaAs detector and a near-IR quartz beam splitter. The functional groups and surface chemistry analysis of solvent-free NCs were carried out on Bruker Alpha-P spectrophotometer equipped with a diamond attenuated total reflectance (ATR) crystal. Electron paramagnetic resonance (EPR) spectra were measured at X-band frequency (9.6 GHz) with a Bruker Elexsys-500 equipped with a Super High QE (ER4123SHQE) cavity. Transmission electron microscopy (TEM) images of freshly prepared NCs drop-casted onto carbon-coated (3nm) copper grids (CF400-CU-50, Electron Microscopy Sciences) were collected on JEOL 2000FX electron microscopy. The average size distribution and

standard deviation were generated by analyzing particles using ImageJ software. Powder X-ray diffraction patterns were collected on Rigaku Smart Lab SE in the Bragg-Brentano configuration with a Cu K-source. An inductively coupled plasma optical emission spectrometer (ICP-OES, Optima 4300 DV - PerkinElmer) was employed to analyze the stoichiometry of metal ions in NCs.

5.2.4 Photochemical Reduction

Photochemical reduction, also known as photodoping, was performed to introduce extra electrons in as-prepared colloidal NCs. Sample preparation was completed in the Argon-filled glove box where NCs suspended in hexanes with ~5% EtOH as sacrificial reductant were transferred to a 1 mm air-tight cuvette before removal. Samples were photodoped by prolonged exposure to irradiation from 75 W xenon lamp. The electronic absorption measurements were taken periodically during the entire photodoping process. When absorption did not change over an exposure of a few hours, samples were considered to have reached their maximum photodoping level. For EPR measurements, NCs were transferred to a 4 mm air-tight customized quartz tube, and photodoping inside the EPR resonator was carried out by UV irradiation using fiber optics. To maintain a similar irradiation intensity, the lamp was warmed up for about 30 minutes before each photodoping experiment.

5.3 Results and Discussion

The synthesis of BaTiO₃ and 0.1% Cr-doped BaTiO₃ NCs (referred to as Cr:BaTiO₃ hereafter) was carried out by a facile hydrothermal method developed previously by our group.^{15,16} Various optimization parameters such as pH, temperature, and time are summarized in supporting information (see SI Figure 5.6-

5.10 and Table 5.1-5.4). The crystallinity and phase purity of as-prepared NCs were confirmed by powder X-ray diffraction measurements. The diffraction patterns of both samples are indexed to the cubic phase of BaTiO₃ with space group Pm $\bar{3}$ m and no considerable secondary phases were observed (see SI Figure 5.11).^{17,18} The TEM images shown in the supporting information (see SI Figure 5.12) and its analysis confirmed that NCs exhibit cuboidal morphology with an average diameter of 12.8 ± 1.8 nm. FTIR measurements showed that the surface of NCs is passivated with oleic acid and oleylamine to provide colloidal stability (see SI Figure 5.13).

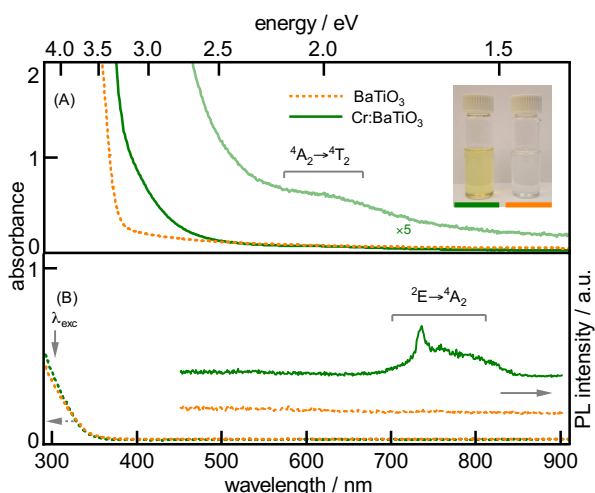


Figure 5.1 (A) Electronic absorption spectra of a concentrated solution of as-prepared BaTiO₃ and Cr:BaTiO₃ NCs in hexanes. The color photographs of both samples are shown in the inset. (B) Photoluminescence spectra of dilute samples excited at 300 nm. The absorption spectra of both samples with absorbance less than 0.5 at the excitation wavelength, as shown on the left-axis.

Figure 5.1A shows the room temperature UV-Vis spectra and color photographs of concentrated suspensions of BaTiO₃ and Cr:BaTiO₃ NCs in hexanes. The electronic absorption spectra of both samples are dominated by the band-edge transition *ca.* 3.25 eV in good agreement with single-crystal BaTiO₃.¹⁴ However, Cr:BaTiO₃ NCs also exhibit an absorption tail into the visible region and appear crispy yellow as depicted from color photographs shown in Figure 5.1A inset. This absorption smearing out of the band-edge is attributed to a combination of charge

transfer transition involving Cr^{3+} sub-bandgap impurity levels and ligand-field $^4\text{A}_2 \rightarrow ^4\text{T}_1$ transition of Cr^{3+} in cubic symmetry.¹⁵ The low energy $^4\text{A}_2 \rightarrow ^4\text{T}_2$ transition of Cr^{3+} is also weakly observed, as shown in enlarged magnification *ca.* 2.0 eV. Figure 5.1B shows the room temperature emission spectra of dilute solutions of both NCs excited at 300 nm. The characteristic emission at ~ 1.68 eV originates from spin-forbidden $^2\text{E} \rightarrow ^4\text{A}_2$ transition of octahedral Cr^{3+} in $\text{Cr}:\text{BaTiO}_3$ NCs.^{16,17}

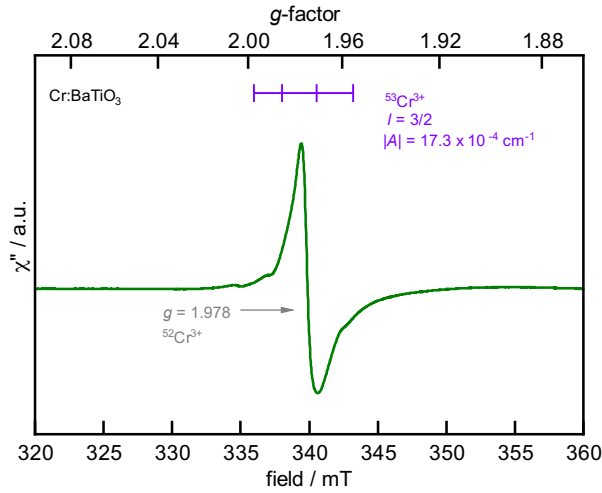


Figure 5.2 Room temperature EPR spectrum of as-prepared $\text{Cr}:\text{BaTiO}_3$ NCs. A weak hyperfine from $^{53}\text{Cr}^{3+}$ nuclear spin ($I = 3/4$) is represented with a four-line bar around broad resonance.

To probe the Cr^{3+} speciation, we employed EPR spectroscopy at X-band frequency (9.5 GHz). Cr^{3+} dopants in the BaTiO_3 lattices preferentially substitute Ti^{3+} ions because of the similar ionic radii in octahedral cubic symmetry [$(\text{Cr}^{3+}) = 0.62 \text{ \AA}$, $(\text{Ti}^{4+}) = 0.61 \text{ \AA}$].¹⁸ Figure 5.2 shows the room temperature EPR spectrum of as-prepared $\text{Cr}:\text{BaTiO}_3$ NCs. The broad isotropic resonance at $g = 1.978$ corresponds to spin-allowed ($\Delta m_s = \pm 1$) transition of $^{52}\text{Cr}^{3+}$ ($S = 3/2$) in octahedral symmetry consistent with Cr -doped SrTiO_3 NCs.^{19,20} A weak 4-line hyperfine ($|A| = 17.3 \times 10^{-1} \text{ cm}^{-1}$) originating from $^{53}\text{Cr}^{3+}$ ($I = 3/2$) nuclei with natural abundance of 9.5% is displayed with a top bar above broad resonance.

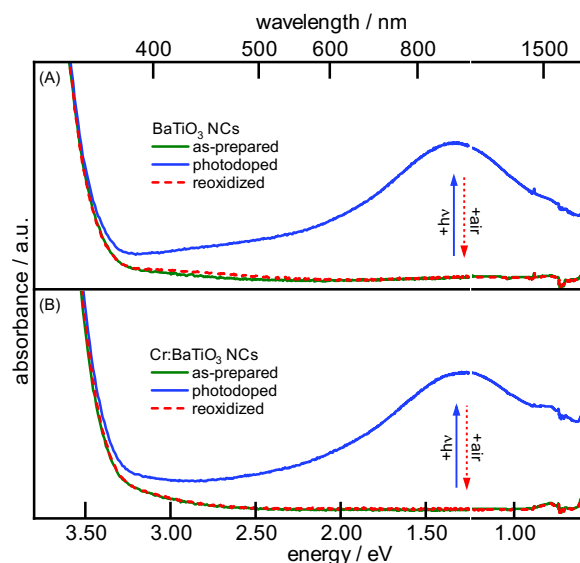


Figure 5.3 Electronic absorption spectra of (A) BaTiO₃ and (B) Cr:BaTiO₃ NCs solution in hexanes as a function of photodoping time. The green spectrum belongs to as-prepared NCs, and the blue spectrum corresponds to maximally photodoped NCs. The red dotted line overlapping the as-prepared NCs belongs to the reoxidized sample.

To introduce n-type charge carriers, photodoping experiments were performed (see the experimental section for detailed procedure). Briefly, UV irradiation of anaerobic solution of as-prepared NCs creates electron/hole pairs (excitons), and the rapid quenching of holes with ethanol leaves extra electrons in the NCs. Figure 5.3 shows the UV-visible spectra of NCs as a function of photodoping time. The electronic absorption spectrum of as-prepared NCs is dominated by the band-edge transition above 3.25 eV. With increased photodoping time, the physical appearance of the samples gradually changes to blue with a broad absorption feature that extends throughout the entire visible region into near-IR. The continuous growth of this new absorption feature without changing the band-edge energy suggests that added electrons are not free or delocalized in the conduction band (e_{CB}^-), which would have otherwise displayed a significant blue-shift in band-edge absorption due to the Moss-Burstein effect.^{20,21} Instead, added electrons create localized Ti³⁺ trap states. In our recent study on the photodoped SrTiO₃ NCs,¹⁴ similar blue colorations and an

absorption tail through the visible region was observed. We attributed this broad absorption centered at *ca.* 1.2 eV to charge transfer transition of trapped electrons at polaronic Ti^{3+} sites to the conduction band of BaTiO_3 NCs. There is no significant difference in absorption features between BaTiO_3 and $\text{Cr}:\text{BaTiO}_3$ NCs.

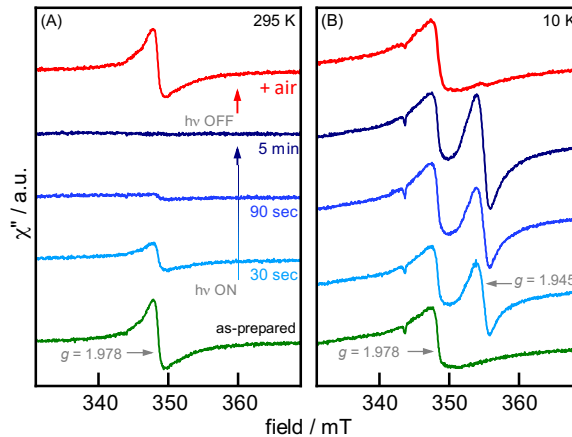


Figure 5.4 EPR spectra of $\text{Cr}:\text{BaTiO}_3$ colloidal solutions in hexanes with increased photodoping time at room temperature (left panel-A) and at cryogenic temperature 10 K (right panel-B) with increased photodoping time. The red spectra in both panels correspond to reoxidized samples with air.

To probe the interplay between these polaronic electrons and the substitutional Cr^{3+} dopants, we employed EPR spectroscopy. Figure 5.4A shows the room temperature EPR spectra of $\text{Cr}:\text{BaTiO}_3$ NCs collected at various photodoping times. With very short photodoping time, the EPR signal at $g = 1.978$ corresponding to Cr^{3+} sequentially decreases in intensity and completely disappears in 5 min of photodoping time. This EPR signal is quantitatively recovered upon reopening the photodoped samples to air. However, the acquisition of EPR measurements at 10 K as displayed in Figure 5.4B shows that the Cr^{3+} EPR signal remains *intact* with the photodoping, while a new $g = 1.945$ signal starts appearing. This new signal corresponds to the paramagnetic Ti^{3+} ($d^1, S = \frac{1}{2}$) defects that appear only at cryogenic temperature owing to its fast spin relaxation time. This observation suggests that Cr^{3+} dopants are mimicking the spin dynamics of Ti^{3+} in photodoped NCs - acceleration of

spin relaxation times through a near-resonant cross-relaxation phenomenon consistent with our previous report.²¹

To further confirm the existence of spin acceleration, we measured the power saturation rollover behavior of Cr^{3+} dopants in the absence and presence of Ti^{3+} defects. In this particular experiment, a series of EPR spectra on as-prepared and photodoped NCs were collected at various applied microwave powers traversing from lowest to the maximum capacity of the equipment as discussed in an earlier report and also in chapter 3 above.²² Figure 5.5A-B plots the peak-to-peak signal intensity (ΔY) extracted from EPR spectra of Cr^{3+} in both the as-prepared and the photodoped NCs and Ti^{3+} in photodoped NCs against the microwave power (h_1) at 10 K. At low power end, ΔY for Cr^{3+} in as-prepared NCs (green data) increases approximately linearly with increasing h_1 , but at high power, it rolls over and begins to saturate. While at the same power range, Cr^{3+} and Ti^{3+} signal intensities in photodoped NCs (blue data) increase continuously without showing any signs of saturation. This typical behavior of signal intensities in photodoped NCs reflects the acceleration of spin relaxation times as observed previously in Cr^{3+} -doped SrTiO_3 NCs. Due to the absence of rollover behavior (saturation) in photodoped NCs, fitting this data is challenging to extract the spin relaxation rates. However, pulse EPR measurements could be carried out to quantify the T_1 and T_2 times directly. Nevertheless, this data suggests that spin interactions between Ti^{3+} and Cr^{3+} exist regardless of the A-site cation in the ABO_3 -type cubic crystal structure.

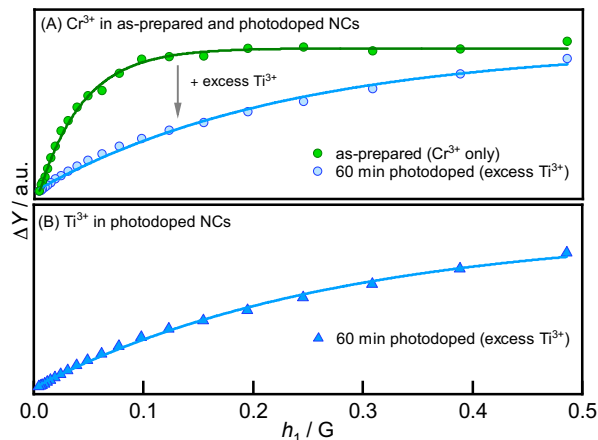


Figure 5.5 Normalized peak-to-peak intensity (ΔY) of 10 K CW-EPR spectra of (A) Cr^{3+} in as-prepared and photodoped NCs and (B) Ti^{3+} in photodoped NCs against the microwave power (h_1).

To validate the existence of such interactions and the effect of site symmetry, we extended this study to 0.1% Cr-doped TiO_2 NCs (referred to as Cr: TiO_2 hereafter) prepared from a solvothermal method (see the experimental section for the procedure and SI Figure 5.14 for PXRD measurements). Room temperature EPR measurements on as-prepared samples show anisotropic signal at g -value = 1.978 with a prominent zero-field splitting consistent with simulated EPR spectrum using EasySpin (see SI Figure 5.15). This signal is assigned to the substitution of Cr^{3+} at Ti^{4+} site in anatase TiO_2 NCs. The electronic absorption spectra displayed in Figure 5.16A-B show the appearance of a broad absorption feature tailing through the band-edge with photodoping. This characteristic broad absorption that maximizes into the mid-IR has been previously observed in n-type TiO_2 nanoparticles and is attributed to the localized surface plasmon resonance (LSPR).^{23,24} A clear blue-shift in the band-edge absorption with photodoping due to the Moss-Burstein effect indicates the free or delocalized electrons in the conduction band. It is important to notice that in addition to the LSPR-related broad absorption, a weak feature around $\sim 1.75 \text{ eV}$ (800 nm) also appears. The energy of this feature is consistent with the trapped electrons at Ti^{3+}

sites, as reported by Mayer and coworkers in photodoped TiO₂ nanoparticles.²⁵ Valentin *et al.* found that quasi-localized or localized Ti³⁺ states in the n-type TiO₂ lattices are very close in energy (~ 0.1 eV) to the delocalized electrons, and oftentimes they are thermally excited into the conduction band.²⁶ We speculate that a fraction of electrons in photodoped NCs creates thermally-detrapped Ti³⁺ defects, and the charge transfer transition from these Ti³⁺ trap states to the conduction band gives rise to this small absorption feature at ~ 1.75 eV. This observation suggests that the majority of photoexcited electrons are delocalized in the conduction band while that a fraction of them creates quasi-localized Ti³⁺ defects with increased photodoping time.

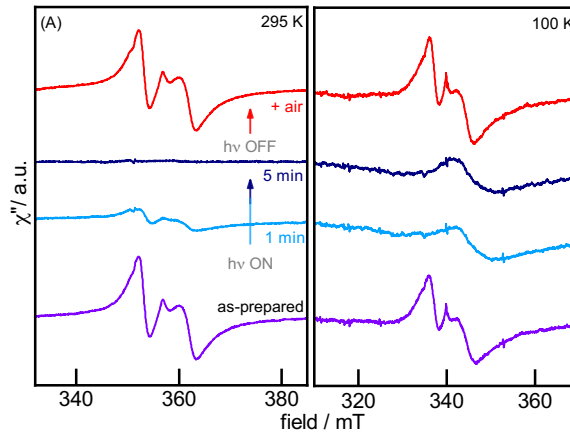


Figure 5.6 EPR spectra of Cr:TiO₂ NCs as a function of photodoping time collected at (A) 295 K and (B) 100 K.

The room temperature EPR spectra of Cr:TiO₂ NCs as a function of photodoping time are shown in Figure 5.6A. With a short photodoping time, the Cr³⁺-related EPR signal is completely disappeared and returns to its original states upon reopening the samples to air. This observation seems consistent with the previous discussion on Cr:BaTiO₃ NCs. However, the EPR measurements at 100 K (Figure 5.6B) show a broad and poorly resolved Ti³⁺ signal. The Ti³⁺ signal in general is observed only at

cryogenic temperature owing to its short spin relaxation times. The increased spin-spin relaxation pathways in the photodoped NCs with a lot of Ti^{3+} in the vicinity of each other cause the broadening of this Ti^{3+} signal. We further suspect the overlapping of Cr^{3+} and Ti^{3+} signals also led to this broadening. Nevertheless, this observation suggests spin interaction between excess electrons and substitutional Cr^{3+} dopants in colloidal TiO_2 . We suspect this phenomenon to exist in other Ti(IV) -based oxide NCs as well. High-frequency EPR measurements to resolve these overlapping signals and pulse EPR would be an alternative approach to further study the changes in spin properties of dopants in these oxide NCs.

5.4 Supporting Information

Optimization of synthesis condition. For the synthesis of high-quality colloidal NCs, several different synthesis parameters were tuned, and the resulting products were analyzed using powder X-ray diffraction, electronic absorption and EPR spectroscopy.

Effect of NaOH concentration. To investigate the effect of hydroxide on NCs growth, the concentration of NaOH was gradually increased from 0.5 to 1.5 M keeping all other conditions unchanged. The PXRD pattern displayed in Figure 5.7 shows that (1) BaTiO_3 phase formation requires more than 0.5 M NaOH concentration and (2) NCs size increases at higher NaOH concentration as summarized in Table 5.1. In addition to that, the solubility of NCs in hexanes also decreases with increased NaOH. Therefore, we employed 1 M NaOH for all follow-up experiments.

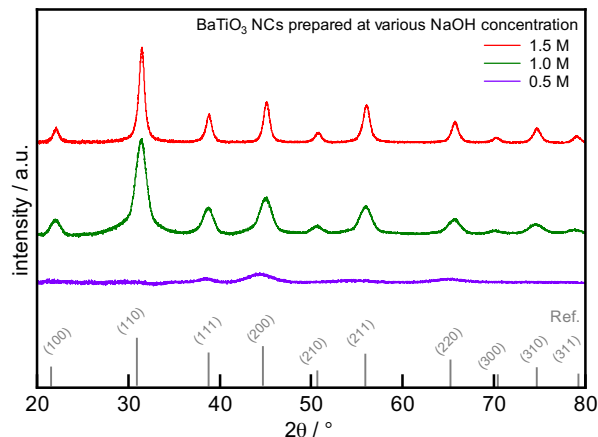


Figure 5.7 Powder X-ray diffraction patterns of BaTiO₃ NCs prepared with different NaOH concentrations at 215 °C for 24 hours.

Table 5.1 The average lattice parameter and crystallite size of BaTiO₃ NCs prepared with different NaOH concentrations at 215 °C for 24 hours.

[NaOH] (M)	Lattice parameter (Å)	Crystallite size (nm)†
1.5	4.0161	11.5 ± 0.41
1.0	4.0230	6.1 ± 0.37
0.5	--	--

†Crystallite size is the average size calculated from the three most intense reflections: (110), (200) and (211) using Rigaku SmartLab Studio II

Effect of temperature. To achieve the optimized growth temperature, hydrothermal synthesis was carried out at different temperatures. The PXRD analysis shows that a high temperature (> 200 °C) is required for the BaTiO₃ phase formation. The analysis summarized in Table 5.2 indicates no significant effect on crystallite size between these three temperatures.

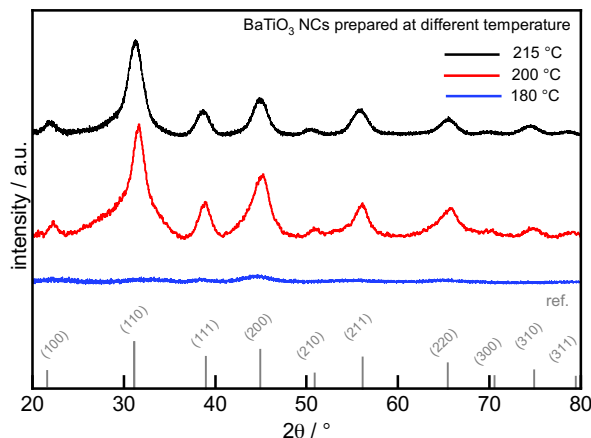


Figure 5.8 Powder X-ray diffraction patterns of BaTiO₃ NCs prepared at different temperatures.

Table 5.2 The average lattice parameter and crystallite size of BaTiO₃ NCs prepared at different temperatures and 1 M NaOH.

Prep. Temp. (°C)	Lattice parameter (Å)	Crystallite size (nm)†
180	--	--
200	4.044	4.29 ± 0.91
215	4.033	4.42 ± 0.54

†Crystallite size is the average size calculated from the three most intense reflections: (110), (200) and (211) using Rigaku SmartLab Studio II

Evolution of BaTiO₃ phase and dopant incorporation. The growth of BaTiO₃ phase was studied by synthesizing the NCs for various times. The PXRD data suggests that phase of BaTiO₃ starts growing at a very short time and then NCs size continuously grow in size with increased time, as summarized in Figure 5.9 and Table 5.3. To analyze the dopant incorporation with increased time, EPR measurements and ICP-OES analysis were performed (see Figure 5.10 and Table 5.4). The measurements show that dopant incorporation happens as the NCs grow in size.

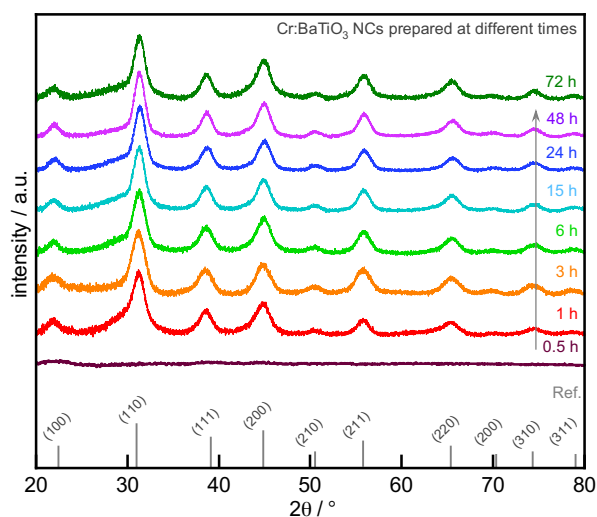


Figure 5.9 Powder X-ray diffraction pattern of NCs prepared at different times showing the evolution of BaTiO₃ phase.

Table 5.3 The average lattice parameter and crystallite size of BaTiO₃ NCs prepared at different times and 215 °C.

Time (h)	Lattice parameter (Å)	Crystallite size (nm)†
0.5	--	--
1	4.034	3.69 ± 1.4
3	4.033	4.31 ± 0.48
6	4.051	5.00 ± 0.39
15	4.055	5.40 ± 0.62
24	4.0497	5.36 ± 0.38
48	4.0404	5.61 ± 0.85
72	4.0348	5.67 ± 0.68

†Crystallite size is the average size calculated from the three most intense reflections: (110), (200) and (211) using Rigaku SmartLab Studio II

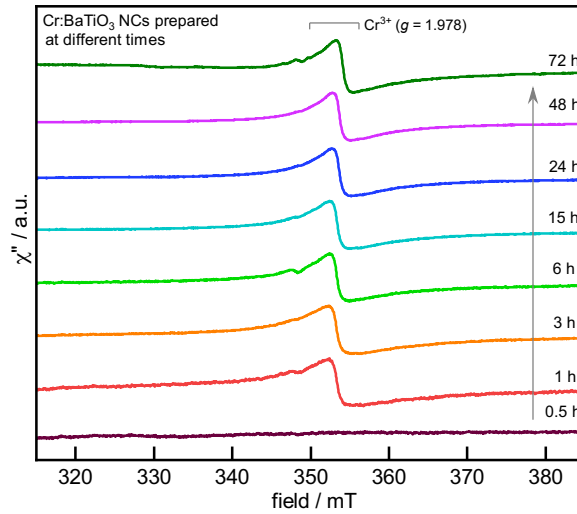


Figure 5.10 Room temperature EPR spectra of Cr-doped BaTiO₃ NCs prepared at various times. The resonance at $g = 1.978$ is attributed to Cr³⁺ substituted at Ti³⁺ sites.

Table 5.4 ICP-OES analysis of Cr-doped BaTiO₃ NCs prepared at different times.

Time (h)	Nominal Cr content (%)	ICP-OES Cr content (%) †
0.5	0.1	0.253
1	0.1	0.297
3	0.1	0.263
6	0.1	0.271
15	0.1	0.282
24	0.1	0.313
48	0.1	0.290
72	0.1	0.290

†The percentage of Cr is calculated based on total Ti content in NCs samples

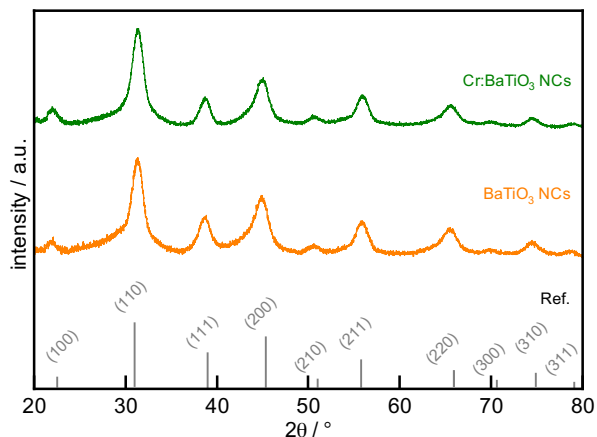


Figure 5.11 Powder x-ray diffraction patterns of BaTiO₃ and Cr:BaTiO₃ NCs indexed to cubic phase of BaTiO₃ with space group $\text{pm}\bar{3}\text{m}$.

Table 5.5 The lattice parameters and crystallite size of BaTiO₃ and Cr:BaTiO₃ NCs determined from powder X-ray diffraction patterns.

Sample NCs	Lattice parameter (Å)	Crystallite size (nm)†
BaTiO ₃	4.0400	5.25 ± 0.58
Cr:BaTiO ₃	4.0272	5.18 ± 0.90

†Crystallite size is the average size calculated from the following reflections: (110), (200) and (211) using Rigaku SmartLab Studio II.

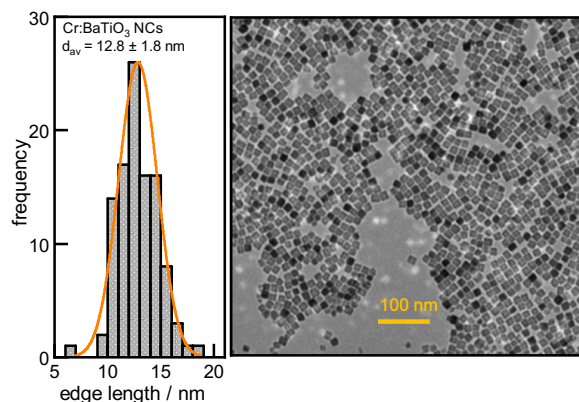


Figure 5.12 Transmission electron microscopy image of as-prepared Cr:BaTiO₃ NCs and size distribution plot showing the average edge length.

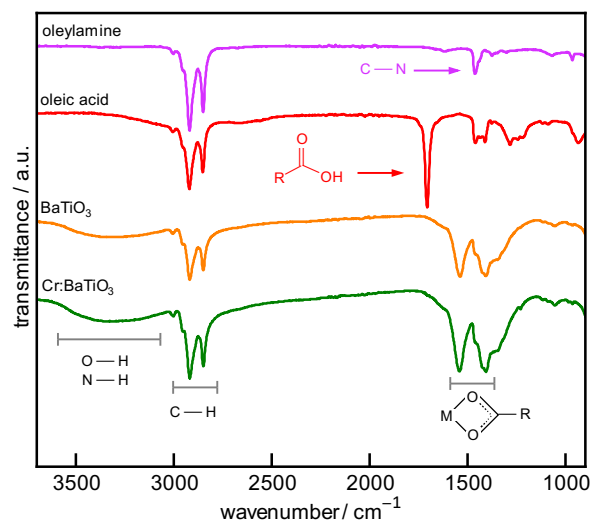


Figure 5.13 FTIR spectra of solvent-evaporated BaTiO₃ (orange) and Cr:BaTiO₃ (green) NCs compared to free oleic acid and oleylamine.

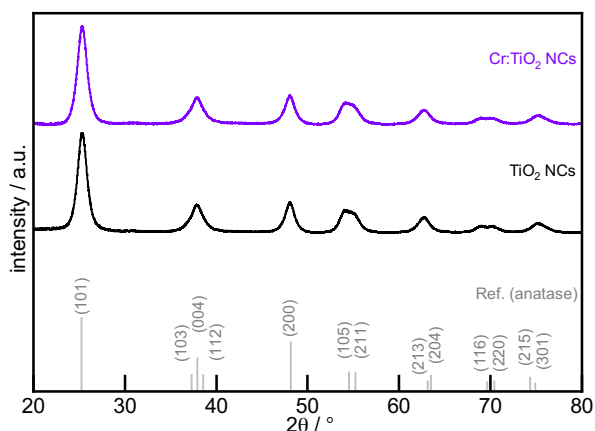


Figure 5.14 Powder X-ray diffraction pattern of TiO₂ and Cr-doped TiO₂ indexed to anatase crystal structure.

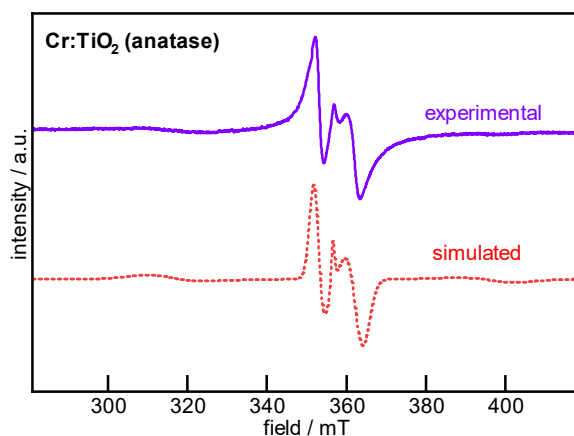


Figure 5.15 Room temperature EPR spectrum of 0.1% Cr-doped TiO₂ NCs compared to simulated EPR spectrum using EasySpin.

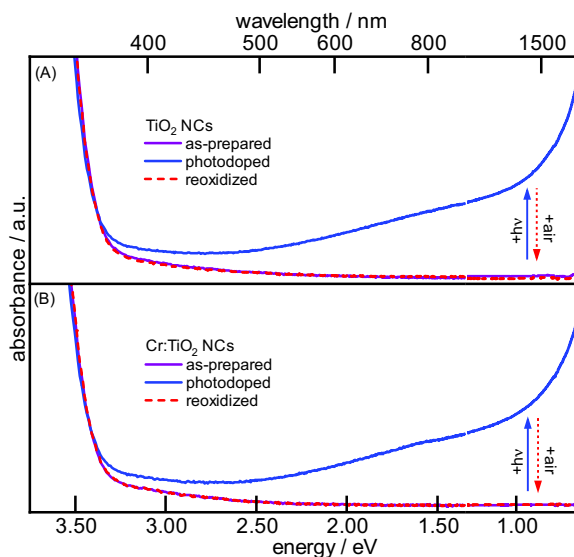


Figure 5.16 Electronic absorption spectra of as-prepared, photodoped and reoxidized samples of (A) TiO_2 and (B) Cr:TiO_2 NCs.

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CHAPTER 6

INVESTIGATING THE PHOTODOPING AND CARRIER DENSITIES OF COLLOIDAL STRONTIUM TITANATE NANOCRYSTALS

6.1 Introduction

The generation and manipulation of the charge carriers in semiconducting materials give rise to new exciting properties which are fundamental to their use in optoelectronic, solar-energy conversions and redox-based applications.¹⁻³ The examples include carrier-induced ferromagnetic ordering in ZnO, MgO and EuTiO₃ nanostructures⁴⁻⁶ and the reversible control of the dopants oxidation state in SrTiO₃ bulk powders.⁷ The interactions between n-type carriers and other defects are susceptible to the nature of host material, and they change significantly from bulk to nanoscale lattices. For instance, the addition of extra electrons in Cr-doped colloidal SrTiO₃ NCs modulates the spin-dynamic of Cr³⁺ dopants while the same carriers cause the reduction of the higher oxidation states of Cr dopant in bulk powder.^{8,9} Similarly, reversible electron accumulation in ZnO NCs has recently gained a lot of interest for energy storage applications.¹⁰ To study the interactions of these carriers in semiconductors, some examples of successful strategies to introduce extra electrons include photochemical, electrochemical and aliovalent doping.¹¹⁻¹³ Among these approaches, photochemical reduction is the most extensively investigated method for modulating carrier densities and is widely applicable to colloidal nanocrystals.^{1,14} It has been demonstrated that kinetics of electrons accumulation and their densities in ZnO NCs are sensitive to nature and concentration of molecular reductant, size of the nanocrystals and identity of counter-cation to stabilize the added electrons.^{3,15}

However, photochemical electrons doping of colloidal SrTiO₃ NCs, its dependence on sacrificial reductant and solvent, and their effect on carrier densities remain elusive.

In this study, a photochemical method for introducing extra charge carriers in colloidal SrTiO₃ NCs is discussed. Photoexcitation of SrTiO₃ NCs in the presence of a sacrificial hole quencher such as ethanol produces extra electrons that are stable indefinitely under anaerobic conditions. Electron paramagnetic resonance (EPR) and electronic absorption spectroscopies measurements provide evidence that these electrons are localized at Ti³⁺ trap sites. The effect of sacrificial hole quencher and irradiation time on the rate of electrons accumulation is thoroughly discussed. We also estimated the average carrier densities per NC by chemical titration against the known concentration of oxidizing agent and find it exceeds that found for other known semiconductors.

6.2 Results and Discussion

The synthesis and characterization of SrTiO₃ NCs were completed by following our recent work.¹⁶ The photodoping experiments were performed on as-prepared SrTiO₃ NCs following the previously reported method.⁹ Briefly, UV irradiation of an anaerobic solution of as-prepared NCs creates electron/hole pairs (excitons) followed by the rapid quenching of holes with sacrificial reductant which leaves excess electrons in the NCs. Figure 6.1 shows the electronic absorption spectra of SrTiO₃ NCs as a function of photodoping time. The absorption spectrum of as-prepared NCs is dominated by the bandgap energy above ~3.25 eV. With increased photodoping time, the physical appearance of the sample changes to dark blue and a new absorption feature extends throughout the entire visible region and is centered

at ~ 1.5 eV (900 nm). This broad near-IR transition is assigned to the metal-to-metal charge transfer (MMCT) transition from localized Ti^{3+} sites to the conduction band reported previously in photodoped SrTiO_3 and TiO_2 .^{17,18} We have no evidence that photodoping creates free or delocalized electrons in the conduction band (e_{CB}^-). Recent reports showed that n-type doping of small anatase TiO_2 NCs could lead to accumulation of e_{CB}^- as observed spectroscopically.^{19,20} The presence of free electrons causes significant blue-shifts in the band-edge energy from the Moss-Burstein effect with the formation of a localized surface plasmon resonance (LSPR) in the mid-IR region. The lack of both the absorption in the mid-IR region and the Moss-Burstein effect confirms that e_{CB}^- are absent in our photodoped SrTiO_3 NCs. Instead, photochemically added electrons create quasi-localized Ti^{3+} defects ($S = \frac{1}{2}$) which can be observed in the EPR spectrum of photodoped NCs (see SI Figure 6.5).

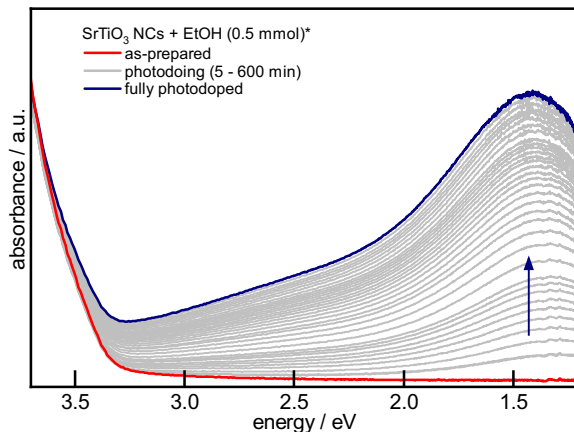


Figure 6.1 Electronic absorption spectra of SrTiO_3 NCs as a function of photodoping time in the presence of ethanol (0.5 mmol) as a sacrificial reductant. The blue arrow shows the progression of time with time.

Although the continuous growth of the absorption feature with increased photodoping time shows the successive addition of electrons, it does not represent the rate of electron accumulation. To gain further insight into it, the relative absorbance at 800 nm ($A_{\text{photodoped}} - A_{\text{as-prepared}}$) plotted against time is displayed in

Figure 6.2. The absorbance (light-blue circle for the data extracted from absorption measurements in Figure 6.1) shows a rapid growth at the initial stages of photodoping which eventually reaches the plateau. The photophysical origin of this rapid increase followed by a slow accumulation of electrons could be various competing processes during the photodoping, as discussed following. The successful electron addition is fundamentally governed by the reactivity of photogenerated holes towards (i) sacrificial reductant, i.e., EtOH and (ii) recombination processes, as demonstrated in Figure 6.3. Each hole being transferred to the sacrificial reductant would result in the addition of one electron. However, after the successful addition of the first electron, further electron accumulation gets complicated, and it must have to go through trion recombination involving already-added electrons and the sacrificial reductant as discussed previously by Schimpf and coworkers in photodoped ZnO NCs.¹⁵ If photogenerated holes are recombined with already-added electrons before they are transferred to sacrificial reductant, further electron addition will not happen. For the successive electron addition, the sacrificial reductant must beat the recombination processes. The relative absorbance in Figure 6.2 shows that one-half of the maximum photodoping is achieved within the first 100 min, and the remaining half is slowly attained in the next 500 min. We speculate that a higher concentration of EtOH as sacrificial reductant ($0.5 \text{ mmol (29 } \mu\text{L)} = 10^3 \times [\text{NCs}]$) and relatively a few added electrons at the early stages of photodoping make the hole quenching reaction much effective. However, every successive electron addition makes the recombination processes more competitive and EtOH less effective, slowing down the accumulation of electrons as the photodoping time progresses.

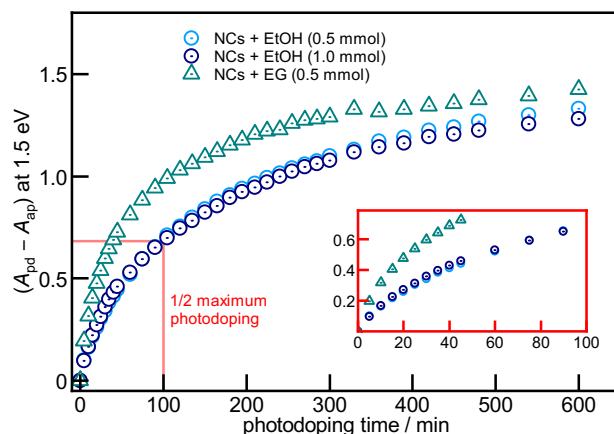


Figure 6.2 The increase in absorbance at 800 nm as a function of photodoping in the presence of two different concentrations of ethanol (0.5 and 1 mmol) and ethylene glycol as sacrificial reductant.

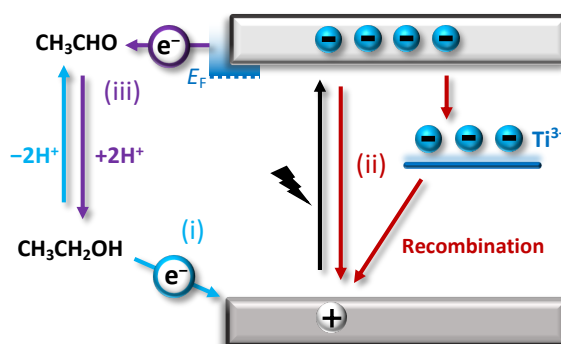


Figure 6.3 Schematic band structure of SrTiO₃ NCs demonstrating the various electron transfer processes during photodoping.

To test the dependence of this electron accumulation on the concentration of sacrificial reductant, the photodoping experiments were performed on the same batch of samples under identical conditions but double the EtOH concentration (1 mmol) and a different sacrificial reductant ethylene glycol (EG) (see SI Figure 6.6 and 6.7). The relative absorbance at 800 nm for both samples plotted against photodoping time is shown in Figure 6.2. The relative change in the absorbance indicates that the rate of electrons accumulation is negligibly faster when the concentration of EtOH is increased to 1 mmol. There are no significant changes in the time scale in achieving ½-maximum photodoping (see inset Figure 6.2). We suspect that the concentration of EtOH (0.5 mmol) is already in significant excess, effectively carrying out the hole scavenging reaction. Furthermore, there is a diffusion limit as well. Therefore, any

further increase in EtOH concentration is not significantly affecting the rate of electron accumulation. For the samples being photodoped in the presence of EG, there is a dramatic increase in the changes in absorbance over time. It quickly reaches the $\frac{1}{2}$ -maximum photodoping in about 50 min compared to 100 min with EtOH (see Figure 6.2 insets). This drastic difference in the rate of electron addition between EtOH and EG possibly arises from their relative reactivities towards holes and the number of electrons they can donate per molecule. Theoretically, each EtOH molecule as a sacrificial reductant donates two electrons upon its complete oxidation to acetaldehyde. The H-NMR data collected on the photodoped samples with EtOH shows a signature low-field peak for aldehydic proton at ~ 9.6 ppm (see SI Figure 6.8). However, methyl protons of acetaldehyde are not observed, and we suspect they might have been occluded by the broad signals from surface-bound oleate ligands. Nevertheless, the peak at ~ 9.6 ppm suggests that at least a fraction of EtOH is oxidized to acetaldehyde during photodoping. The H-NMR spectrum collected on NCs photodoped with EG also shows a peak at ~ 9.6 ppm (see SI Figure 6.9) which we attribute to the presence of acetaldehyde from the oxidation of residual EtOH from washing procedures. We do not observe any other low-field signal in the H-NMR spectrum.

EG has been previously reported as a highly effective sacrificial reductant to suppress the photoexcited charge recombination in TiO_2 nanoparticles.^{21,22} The oxidation of EG produces a mixture of various products, including but not limited to glycolaldehyde, acetaldehyde, formaldehyde, formic acid, and the multiple electrons. However, the mechanisms of such degradation reactions are still unknown. We

speculate that twice the number of active functional groups (-OH) in EG compared to EtOH and its multiple electron donation abilities can play a key role in the faster electron accumulation at the initial stages of photodoping. Although photodoping with EtOH and EG have different rates, they eventually reach about a similar absorbance where further UV irradiation yields no significant changes. The microscopic origin of this saturation has been previously identified as the reduction of acetaldehyde back to EtOH in photodoped ZnO NCs.³ This hydrogenation reaction is sensitive to changes in Fermi level which at low electron densities is negligible but dominates with increased photodoping time (see pathway iii in Figure 6.3). The saturation in absorbance with prolonged photodoping time suggests the exitance of limit for electron doping in these SrTiO₃ NCs. This further motivated us to determine the maximum number of electrons added in these photodoped colloidal NCs.

Although stable indefinitely under anaerobic conditions, these photodoped NCs can be reversed quantitatively by simply reoxidizing the samples with air.^{9,17} This reversibility makes the photodoped colloidal NCs amenable to direct redox titration to estimate the added electrons, something not practicable with chemically reduced bulk powders.⁷ The detailed method for determining the added electrons was adopted from a procedure developed by Mayer *et al.* on colloidal ZnO and TiO₂ nanoparticles.¹⁸ The typical method involves (1) the preparation of colloidal solution with the known concentration of NCs, (2) photodoping of the NCs to their maximum level, and (3) titration against one-electron oxidant under strictly air-free conditions. Briefly, the concentration of colloidal NCs was determined by the combination of the average amount of titanium per NC from size analysis from x-ray diffraction pattern

(see SI Figure 6.10 and Table 6.1-6.2) and the total titanium content in the stock solution from ICP-OES analysis (see SI Figure 6.11-12 and Table 6.3). Since there is one titanium per unit cell of SrTiO_3 , the concentration of NCs can be estimated from ratiometric analysis as summarized in Table 6.4.

For the photodoping experiments, NCs suspended in hexanes of known molarity were photodoped by prolonged exposure to UV irradiation in the presence of EtOH (1 mmol). The UV-Vis absorption measurements were taken periodically during the entire photodoping process. When absorption did not change over an exposure of a few hours, NCs were considered to have reached their maximum photodoping level. The average number of electrons added per NC (n_e) was estimated by stoichiometric chemical titration against TEMPO with known molarity using a customized air-tight glass setup (see SI Figure 6.13 and Table 6.5). In a typical titration experiment, aliquots of a freshly prepared toluene solution of TEMPO were added to the photodoped samples under constant stirring. The absorption measurements were taken after each aliquot addition, as shown in Figure 6.4A (see Figure 6.14-15 for replicate experiments). The complete elimination of absorption attributing to added electrons was considered an endpoint of the titration.

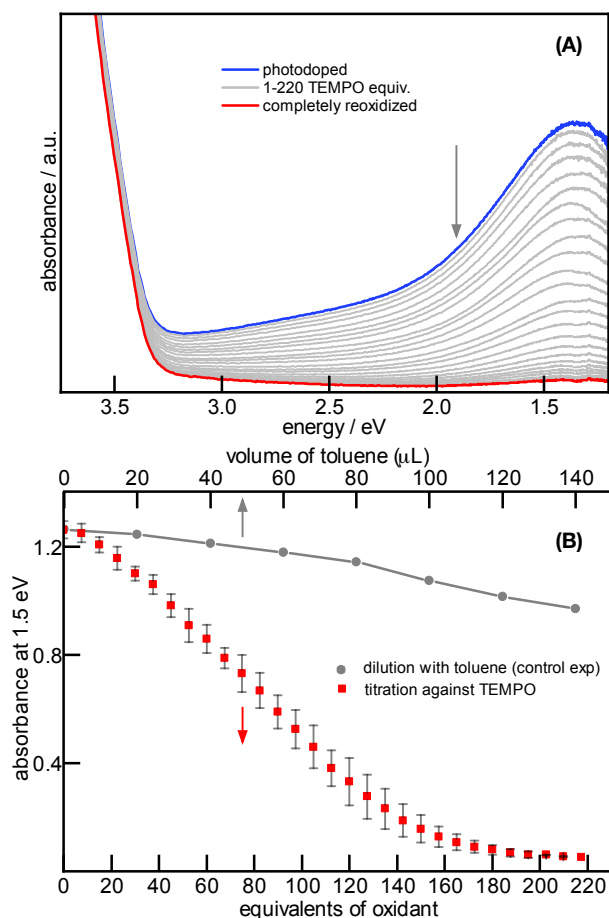


Figure 6.4 (A) Electronic absorption spectra of photodoped NCs with the sequential addition of TEMPO. (B) The change in absorbance at 1.5 eV plotted against the equivalent of TEMPO as an oxidizing agent. The error bars are shown in grey vertical capped lines and the top grey data shows the change in absorbance of photodoped NCs with the addition of blank toluene as a control experiment.

The changes in the absorbance at 1.5 eV with the addition equivalents on TEMPO along with the toluene plotted in Figure 6.4B shows a complete elimination at ~ 180 equivalents. Since TEMPO is a one-electron oxidant, this equates to $n_e = 180$ (see supporting information for detailed calculations). This number for SrTiO_3 here exceeds the previous reports on amorphous TiO_2 nanoparticles ($n_e = 72$) and ZnO NCs ($n_e = 5$).¹⁸ We suspect that the large size of SrTiO_3 NCs (~ 7 nm compared to 3 nm TiO_2) can lead to this difference in number of elections accumulated with photodoping. Brozek *et al.* recently reported a spike in electron accumulation ($n_e = 250$) in Fe-doped ZnO NCs.¹⁰ In these NCs, $\text{Fe}^{3+/2+}$ redox level is located below the

conduction band which acts as a trap site for photoexcited carriers facilitating the electron accumulation.²³ We have also recently seen a similar electron addition at $\text{Fe}^{3+/2+}$ dopant sites in colloidal SrTiO_3 and BaTiO_3 NCs.¹⁷ In this report, we presented a band structure of Ti(IV)-oxide nanocrystals where $\text{Fe}^{3+/2+}$ dopant level is located below the $\text{Ti}^{4+/3+}$ trap states. Although both redox levels are shallow trap centers, $\text{Fe}^{3+/2+}$ could not only serve a better trap center due to its relatively lower position than the $\text{Ti}^{4+/3+}$ level; it can also offer a tunable carrier density by simply depending on the dopant concentration. The experiments to explore the electron accumulation in doped SrTiO_3 would be an interesting future aspect for this work.

6.3 Supporting Information

Synthesis of colloidal SrTiO_3 NCs. The synthesis of colloidal SrTiO_3 NCs was carried out by a modified hydrothermal method reported earlier by Harrigan *et al.*²⁴ In a typical synthesis of 1.25 mmol of TALH and 1.25 mmol $\text{Sr}(\text{OH})_2 \cdot 8\text{H}_2\text{O}$ were dissolved in 30 mL of distilled water. The pH of the solution was then adjusted to 12.1 with an aqueous solution of NMe_4OH (10 M) followed by the addition of oleic acid (2.5 mmol) and hydrazine (5 mmol). The resulting solution was transferred to a 45 mL Teflon-lined autoclave (4744 General Purpose Acid Digestion Vessel, Parr Instrument Company) and heated to 200 °C in an oven for 24 hours. The resulting NCs were collected, washed with ethanol several times, suspended in hexanes, and sonicated for about 30 minutes to produce a cloudy suspension. Finally, a transparent hexane layer containing NCs was extracted by centrifuging the cloudy suspension (5 min at 4000 rpm). These NCs can be further purified by precipitation with ethanol.

Photodoping of NCs. The photodoping experiments were performed to introduce excess electrons in as-prepared NCs. Sample preparation was completed in the argon-filled glove box where NCs suspended in hexanes were transferred to a 1 cm air-tight cuvette before removal. Samples were photodoped by prolonged exposure to unfocused irradiation from 75 W xenon lamp in the presence of EtOH as a sacrificial reductant. The electronic absorption spectra were taken periodically during the entire photodoping process. It was estimated that the NCs were photodoped to their saturation limit when no further spectroscopic changes could be observed following prolonged exposure to UV irradiation.

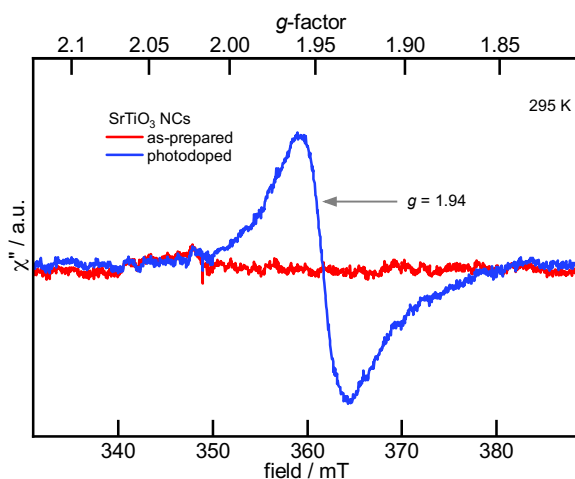


Figure 6.5 Room temperature EPR spectra of as-prepared and photodoped of SrTiO_3 NCs.

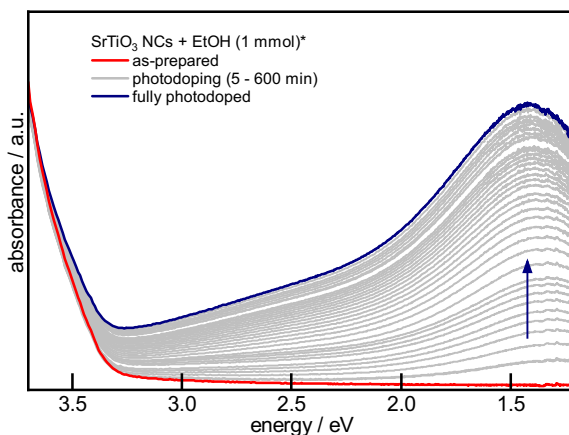


Figure 6.6 UV-vis spectra of SrTiO_3 NCs taken periodically during photodoping in the presence of EtOH (1 mmol).

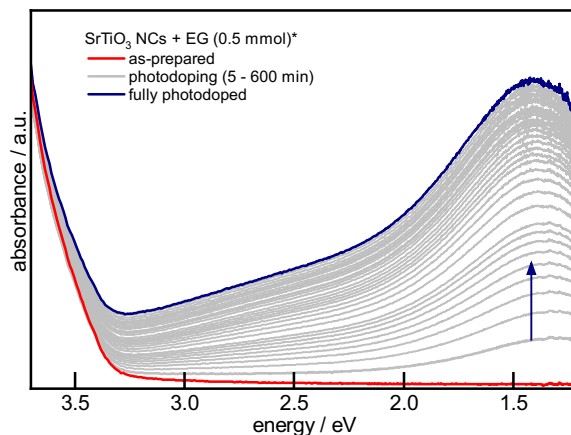


Figure 6.7 UV-vis spectra of SrTiO₃ NCs taken periodically during photodoping in the presence of ethylene glycol (0.5 mmol).

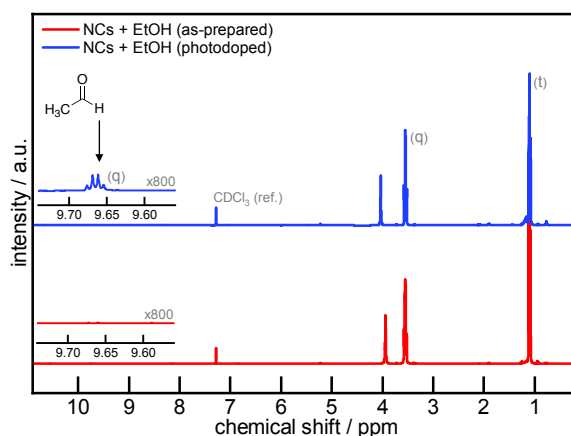


Figure 6.8 H-NMR spectra of colloidal solution of SrTiO₃ NCs suspended in deuterated chloroform before and after the photodoping in the presence of EtOH as hole quencher.

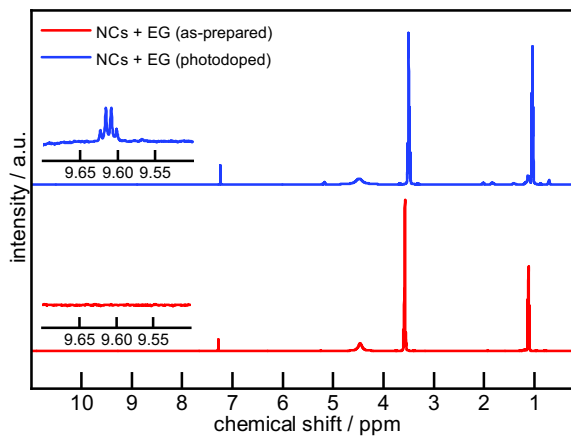


Figure 6.9 H-NMR spectra of colloidal solution of SrTiO₃ NCs suspended in deuterated chloroform before and after the photodoping in the presence of EG as hole quencher.

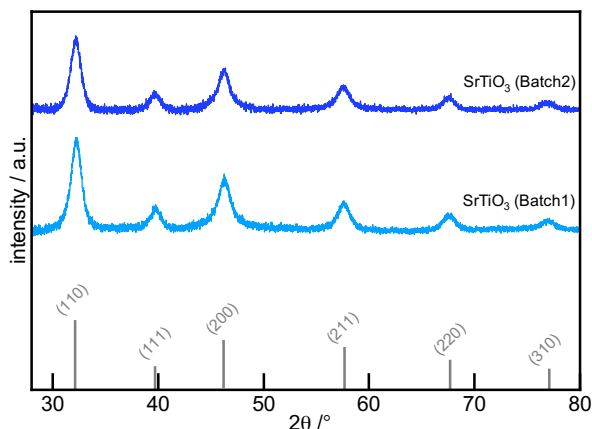


Figure 6.10 Powder X-ray diffraction patterns of SrTiO₃ NCs prepared under identical conditions for average particle size analysis.

Table 6.1 Determination of average particle size from powder diffraction analysis.

SrTiO ₃ NCs	Crystallite size from various reflections (nm)				Average NCs size
	(110)	(111)	(200)	(211)	
Batch-1	7.172	7.177	6.962	6.889	6.974 ± 0.25 nm
Batch-2	7.165	7.185	6.556	6.690	

Table 6.2 Determination of the average number of Ti per NC.

Average lattice parameter	3.924 Å
Average volume of unit cell	(3.924) ³ = 60.42 Å ³
Average size of NCs	69.74 Å
Average volume of NCs	(69.74) ³ = 339192.18 Å ³
Number of unit cells per NC	339192.18 / 60.42 = 5614
Number of Ti ⁴⁺ per NC	5614

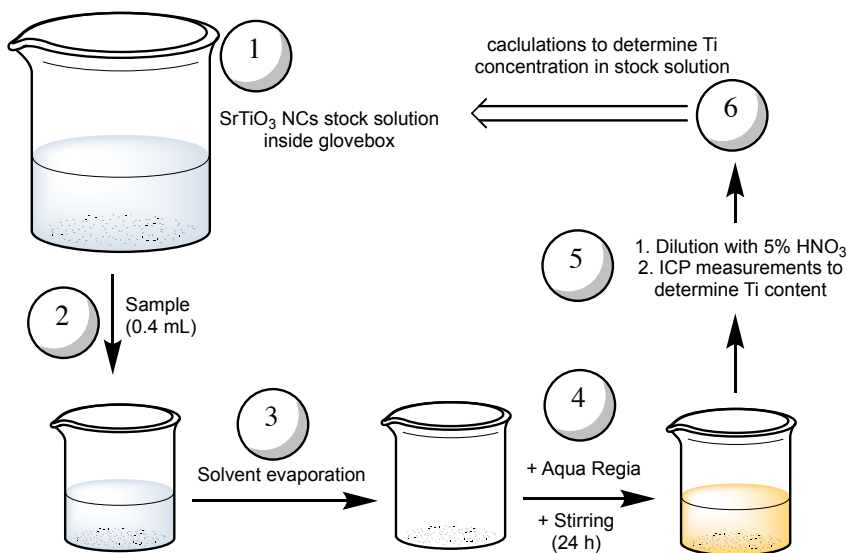


Figure 6.11 Flow chart elaborating the sample preparation of ICP-OES measurements to determine the total titanium content in stock solution.

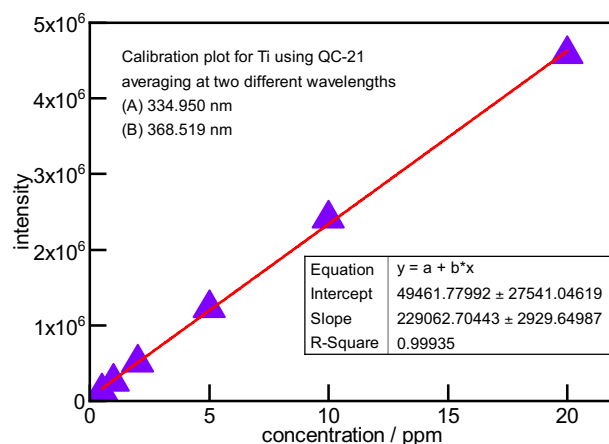


Figure 6.12 Calibration plot for ICP-OES analysis using standard solutions of different concentrations.

Table 6.3 Calculations to determine the total Ti content from ICP-OES data.

SrTiO ₃ NCs	Concentration of Ti (mg/L)			Molarity (mM)	Molarity of stock sol. (dilution factor = 20)
	334.940 nm	368.519 nm	Average		
Sample-1	26.82	27.51	27.26	0.569	11.88 mM
Sample-2	27.11	27.32			
Sample-3	27.23	27.57			

Table 6.4 Determination of concentration of NCs in stock solution.

Concentration of Ti in stock solution	11.88 mM
Number of Ti making up one NC	5614
Concentration of NCs in stock solution	0.00212 mM = 2.12 μ M

Table 6.5 Preparation of stock solution of TEMPO of known concentration.

Molecular weight	186.27
Amount taken	23.9 mg
Volume of solution	20 mL
Molarity of solution	6.4 mM

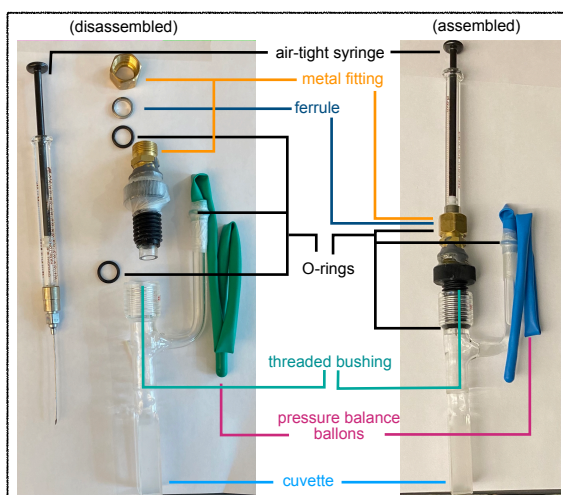


Figure 6.13 Various components of air-free custom setup used titration experiments.

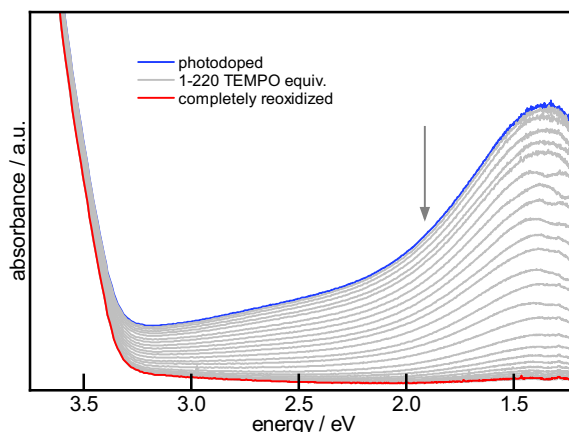


Figure 6.14 UV-Vis spectra of photodoped NCs collected periodically during titration against TEMPO under air-free conditions (trial 1).

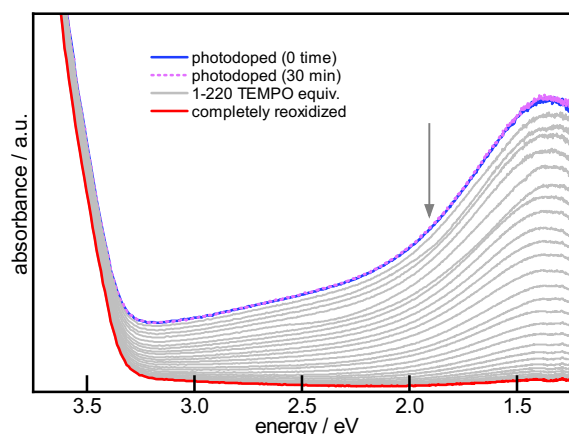


Figure 6.15 UV-Vis spectra of photodoped NCs collected periodically during titration against TEMPO under air-conditions (trial 2). The pink spectrum represents the stability of photodoped NCs during transition between glovebox and air-tight custom setup outside the glovebox.

Calculations to determine the number of electrons per NC:

Volume of photodoped NCs (0.00212 mM) = 2 mL (0.002L)

Number of moles of photodoped NCs = $0.00212 \text{ mmol} / \text{L} \times 2 \text{ mL} = 0.00212 \text{ } \mu\text{mol} / \text{mL} \times 2 \text{ mL} = 0.00424 \text{ } \mu\text{mol} = 0.00000424 \text{ mmol}$

Volume of TEMPO (6.4 mM) added incrementally = 140 μL = 0.140 mL

Number of moles added = $6.4 \text{ mM} \times (0.140/1000) = 0.000896 \text{ mmol}$

Equivalents of TEMPO added against NCs = $0.000896/0.00000424 = 211$

Equivalents of TEMPO added against NCs per increment = total equivalent / no. of data points = $211/28 = \sim 7.5$

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CHAPTER

CONCLUSION AND FUTURE DIRECTIONS

This thesis explores the defect chemistry and dopant-carrier interactions in aliovalent-doped colloidal SrTiO_3 , BaTiO_3 and TiO_2 NCs. In chapter 2, we have presented a systemic study to control the surface defect formation during the hydrothermal synthesis of colloidal SrTiO_3 and BaTiO_3 NCs. We found that an electron transfer step presumably from lactate anions to dissolved O_2 is the critical step that leads to the formation of the surface-adsorbed $\text{O}_2^{\bullet-}$ ions. These surface $\text{O}_2^{\bullet-}$ defects are EPR active and impart a broad feature throughout the visible region that produces a reddish-brown color to the NCs. We also provided evidence that the removal of dissolved O_2 by either addition of hydrazine, purging with argon or using lactate-free Ti(IV) precursors inhibit $\text{O}_2^{\bullet-}$ formation. These results shed light on the role of additives in the chemical synthesis of colloidal ATiO_3 NCs. A potential direction for this work could be exploring the existence of such defect formation and their control in other oxide NCs. We also speculate that these reactive defects can play an integral role in surface-based chemistry such as photocatalysis and there is significant work that can be done to make use of these $\text{O}_2^{\bullet-}$ radicals for various applications. In addition, the brown color that these defects impart to the SrTiO_3 NCs can undoubtedly enhance the optical absorbance for visible light photocatalysis.

Chapter 3 discussed the synthesis and characterization of Fe-doped SrTiO_3 and BaTiO_3 colloidal NCs and investigated their dopant-defect interaction using various spectroscopies. The results presented here provide convincing evidence that the valance state of Fe dopant can be modulated between 3+ and 2+ by post-synthetic

manipulation of carrier densities in these NCs. The introduction of excess charge carriers at room temperature using photodoping methods into $\text{Fe}^{3+}:\text{SrTiO}_3$ and $\text{Fe}^{3+}:\text{BaTiO}_3$ colloidal NCs results in the reduction of Fe^{3+} to Fe^{2+} . Both the dopant and lattice electron storage capacity makes these colloidal NCs potentially suitable for non-aqueous redox flow batteries and charge storage applications. Direct verification of Fe^{2+} is challenging due to the air-sensitive nature of the photodoped NCs. Therefore, the alternative methods to achieve air-stable n-type defects in related d^0 colloidal semiconductor NCs are discussed in chapter 4. In addition, we have seen a significant difference in the microenvironment of Ti^{3+} defects between SrTiO_3 and BaTiO_3 NCs. More sophisticated spectroscopic methods would be required to probe the speciation (surfaces, sub-surfaces and/or cores) of added electrons in these colloidal NCs.

Chapter 4 presented the synthesis and characterization of Fe^{3+} -doped TiO_2 colloidal NCs using a solvothermal method and reversible control of dopant valence state under ambient conditions. The spectroscopic analysis of as-prepared NCs shows that Fe^{3+} is substituted at the Ti^{4+} site, and there is no evidence of any surface dopants. The introduction of excess electrons using the photodoping method can reversibly manipulate the oxidation state of the dopant between Fe^{3+} and Fe^{2+} . We further presented an alternative approach to control the oxidation of Fe^{3+} dopant by incorporating the aliovalent Nb^{5+} as a co-dopant. The substitution of Nb^{5+} at Ti^{4+} adds air-stable excess electrons, majority of which are identified as free carriers in the conduction band. This method reduces the Fe^{3+} dopants to Fe^{2+} , offering control over dopant oxidation under ambient conditions. This latter method opens the avenue to

directly observe the Fe^{2+} dopants and their possible spin interactions with Fe^{3+} and Ti^{3+} defects. There is also a possibility of variable Nb^{5+} doping concentration to achieve the mixed-valence states of Fe^{3+} and the other redox-active dopants in oxide lattices.

Chapter 5 is motivated by previous work on Cr-doped SrTiO_3 NCs, where the effect of excess carriers on substitutional Cr^{3+} dopants in additional Ti-based oxide NCs is discussed. We prepared and characterized the Cr-doped colloidal BaTiO_3 and TiO_2 NCs. The synthetic methods employed in this work allowed the incorporation of Cr^{3+} dopants at substitutional Ti^{4+} sites. The introduction of excess electrons using photodoping methods shows the disappearance of the Cr^{3+} -related EPR signal in both the BaTiO_3 and TiO_2 NCs. We attribute this signal disappearance to changes in spin relaxation times (T_1 and T_2) of Cr^{3+} dopants. This work expands on the interactions we previously found with Cr-doped SrTiO_3 NCs and opens the avenue of near-resonant cross-relaxation as a new tool to control the spin properties of paramagnetic centers. Direct measurements of Cr^{3+} relaxation times, its dependence on the concentration of Ti^{3+} defects, and NCs size and surface chemistry are some of the future directions for this project.

In the last chapter of this thesis, a method to estimate the carrier concentration in photodoped colloidal SrTiO_3 NCs is discussed. We explored the kinetics and various competing processes for electron addition during photodoping. The controlled photodoping experiments show the rate of electron addition is significantly faster at the earlier stages of photodoping time, which can be further enhanced by employing a better sacrificial reductant. In the latter part of this work, we combined x-ray

diffraction and ICP-OES analysis to determine the concentration of NCs that led to streamlining the chemical titration of photodoped NCs. The chemical titration of photodoped NCs against one-electron oxidant TEMPO was performed to estimate the maximum number of electrons stored per NC of SrTiO_3 , which is significantly higher than the previously reported number for TiO_2 and ZnO NCs. This work can lead to using these NCs as solution-based charge storage. There is also a great potential to expand this work on doped SrTiO_3 NCs to tune the kinetics of electron addition and NCs carrier storage capacity. Effect of NCs size, surface chemistry and morphology are some of the challenges to navigate this project in the future.

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