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Molecular and Atomic Gas in the Local Group Galaxy M33

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ABSTRACT

We present high resolution large scale observations of the molecular and atomic gas in the Local Group Galaxy M33. The observations were carried out using the HIterodyne Receiver Array (HERA) at the 30m IRAM telescope in the CO(2–1) line achieving a resolution of 12”×2.6 km s⁻¹, enabling individual Giant Molecular Clouds (GMCs) to be resolved. The observed region is 650 square arcminutes mainly along the major axis and is cut to a radius of 8.5 kpc. and covers entirely the 2’ × 40’ radial strip observed with the HIFI and PACS spectrometers as part of the Herschel program. The achieved sensitivity in main beam temperature is 20–50 mK at 2.6 km s⁻¹ velocity resolution. The CO(2–1) luminosity of the observed region is 1.7 ± 0.1 × 10⁸ K km s⁻¹ pc² and is estimated to be 2.8 ± 0.3 × 10⁷ K km s⁻¹ pc² for the entire galaxy, corresponding to H₂ masses of 1.9 × 10⁸ M⊙ and 3.3 × 10⁸ M⊙ respectively (including He), calculated with a N(H₂)/I_{CO(1–0)} twice the Galactic value due to the half-solar metallicity of M33. H₁ 21 cm VLA archive observations were reduced and the mosaic was imaged and cleaned using the multi-scale task in the CASA software package, achieving a series of datacubes with resolutions ranging from 5” to 25”. The H₁ mass within a radius of 8.5 kpc is estimated to be 1.4 × 10⁹ M⊙. The azimuthally averaged CO surface brightness decreases exponentially with a scale length of 1.9 ± 0.1 kpc whereas the atomic gas surface density is constant at ΣHI = 6 ± 2 M⊙ pc⁻² exceeded to face-on. For a N(H₂)/I_{CO(1–0)} conversion factor twice that of the Milky Way, the central kiloparsec H₂ surface density is ΣHI = 8.5 ± 0.2 M⊙ pc⁻². The star formation rate per unit molecular gas (SF Efficiency, the rate of transformation of molecular gas into stars), as traced by the ratio of CO to H₂ and FIR brightness, is constant with radius. The SFE, with a N(H₂)/I_{CO(1–0)} factor twice galactic, appears 2–4 times greater than of large spiral galaxies. A morphological comparison of molecular and atomic gas with tracers of star formation is presented showing good agreement between these maps both in terms of peaks and holes. A few exceptions are noted. Several spectra, including those of a molecular cloud situated more than 8 kpc from the galaxy center, are presented.


1. Introduction

The Local Group galaxies span a broad range in mass, luminosity, morphology, and metallicity. Two large spirals (the Milky Way and M31) are the centers of two galaxy sub-groupings, each being surrounded by a large number of dwarf galaxies. In addition, M31 — the Andromeda Galaxy — has a small spiral companion, M33 (the Triangulum Galaxy); their separation is approximately 15 degrees, corresponding to 200 kpc (assuming a common distance of 840 kpc; Galleti et al. (2004)). Gaseous streams are observed between them, indicating tidal interaction (Putman et al. 2009). M33 provides a means of observing a galaxy morphologically similar to our own but with a mass only a tenth of the Milky Way and factor two lower metallicity (Rosolowsky & Simon 2008; Magrini et al. 2009). Further evidence for the difference between M33 and the Milky Way is the large gas fraction and blue stellar colors of the former relative to the latter. M33 thus represents an environment in which to study the interstellar medium (ISM) and star formation (SF) that cannot be replaced by Galactic observations and where individual GMCs can be resolved to probe their star formation activity. It may also be possible to apply what we learn by studying M33 to
the physics of early-universe objects, which share many of the characteristics of M33.

In this article we present sensitive and high-resolution mapping observations of the CO $J = 2 \rightarrow 1$ transition in M33 in order to study the morphology and dynamics of the molecular component. The total mapped area covers 650 square arcminutes, mainly along the major axis of the galaxy.

A $2' \times 40'$ wide strip along the major axis (see Fig. 1) was observed to a particularly low noise level of 25 mK at 2.6 km s$^{-1}$ velocity resolution to compare with the sensitive Cn Herschel/HIFI and Herschel/PACS spectroscopy observations which will be obtained as part of the HERM33ES Herschel Key Program (Kramer et al. 2010). While the most sensitive and among the highest resolution, these are not the first maps of M33 in the CO lines. Engargiola et al. (2003) observed the whole of the inner disk (up to about 5kpc along the major axis) with the BIMA array at 13'' resolution. Hever et al. (2004) observed the inner disk and a small major axis strip at 50'' resolution with FCRAO. Rosolowsky et al. (2005) combined the BIMA+FCRAO+NRO data to improve the sensitivity and resolution of the previous maps; and Gardan et al. (2007) observed a rectangle at high sensitivity and 15'' resolution extending from NGC 604 to the R$_{25}$ radius. Table 2 summarizes the characteristics of previous molecular and atomic gas surveys in M33. This work extends the Gardan et al. (2007) work further North and to the South at higher resolution and sensitivity. Earlier studies of GMCs in M33 include Wilson et al. (1997) and Rosolowsky et al. (2003) and studies similar to our own of other Local Group galaxies have been made by e.g. Fukui et al. (2008); Israel et al. (2003); Leroy et al. (2006); Nieten et al. (2006).

As mentioned by Blitz & Rosolowsky (2006) for M33 and Leroy et al. (2006) for IC 10 and discussed more extensively by Gardan et al. (2007), the Star Formation Rate (SFR) per unit H$_2$ mass or Star Formation Efficiency (SFE = SFR/M(H$_2$)) in yr$^{-1}$ was found to be up to an order of magnitude higher in these small galaxies than in large spirals. This appears to be the case in distant galaxies as well, given the factor 10 increase in (comoving) SFR density (e.g. Madau et al. 1996; Wilkins et al. 2008). Are there local universe analogs of these distant objects? Is M33 one of them?

One of the obvious questions is whether the H$_2$ mass has not been underestimated in these subsolar metallicity objects. The articles using the data presented here and as part of the HERM33ES project will attempt to answer that issue clearly.

Metallicities were lower in the past and H$_2$ production is believed to take place on grain surfaces – therefore H$_{2}$ to H$_{2}$ conversion is expected to be less efficient in low metallicity systems (Krumholz et al. 2009), and this is indeed observed in many systems (e.g. Leroy et al. 2007). The conversion of H$_{1}$ to H$_{2}$ in the intermediate redshift systems would have to be much more efficient than today, generating unusually high molecular fractions in these distant galaxies (contrary to expectations) with less stellar gravity, in order to have a similar efficiency in converting H$_{2}$ to stars.

Here, we will present new maps of CO and HI. In this first paper, we restrict ourselves to a study of the radial distribution of the molecular gas and infrared surface brightness, the molecular and atomic gas surface densities and the star formation efficiency and to a qualitative comparison between the maps of star formation rate tracers, i.e. the dust maps, and the gas maps. In addition, we will discuss CO spectra in a few selected regions. A series of articles will follow, focussing on at least (i) cloud populations, life cycle, and mass spectrum (ii) dynamics of the molecular gas and the role of spiral arms (iii) diffuse CO emission, after subtraction of the clouds identified and (iv) a more detailed comparison of the star formation rate – gas surface density relation.

\begin{table}[h]
\centering
\caption{Adopted parameters for M33}
\begin{tabular}{ll}
\hline
\hline
Parameter & Value \\
\hline
\(\alpha(12000)\) & 1$^{b}$33$^{m}$509 \\
\(\delta(12000)\) & +30$^{\circ}$39$^{\prime}$39$^{\prime\prime}$ \\
Distance & 840 kpc $^a$ \\
Optical Radius \(R_{25}\) & 30.8$^{b}$ $^b$ \\
Inclination & 56$^{b}$ $^b$ \\
Position angle & 22.5$^{b}$ $^b$ \\
Central Oxygen abundance & 12 + log(O/H) = 8.4 $^c$ \\
\hline
\end{tabular}
\end{table}

\textsuperscript{a} Galleti et al. (2004)
\textsuperscript{b} HIYPERLEDA (Paturel et al. 2003)
\textsuperscript{c} Magrini et al. (2009)
Table 2. Comparison of past survey characteristics.

<table>
<thead>
<tr>
<th>CO Survey</th>
<th>Telescope</th>
<th>Line</th>
<th>Angular resolution (&quot;)</th>
<th>Spectral resolution (km s(^{-1}))</th>
<th>Area (arcmin(^2))</th>
<th>(R_{\text{max}}) (kpc)</th>
<th>Sensitivity per channel (mK)</th>
</tr>
</thead>
<tbody>
<tr>
<td>This paper</td>
<td>IRAM 30m</td>
<td>CO(2–1)</td>
<td>12</td>
<td>2.6</td>
<td>643</td>
<td>8</td>
<td>15–30</td>
</tr>
<tr>
<td>Gardan et al. (2007)</td>
<td>IRAM 30m</td>
<td>CO(2–1)</td>
<td>15</td>
<td>2.6</td>
<td>253</td>
<td>7</td>
<td>15–30</td>
</tr>
<tr>
<td>Rosolowsky et al. (2007)</td>
<td>BIMA+NRO+FCRAO</td>
<td>CO(1–0)</td>
<td>20</td>
<td>2.6</td>
<td>170</td>
<td>2</td>
<td>60</td>
</tr>
<tr>
<td>Rosolowsky et al. (2007)</td>
<td>BIMA+FCRAO</td>
<td>CO(1–0)</td>
<td>13</td>
<td>2</td>
<td>900</td>
<td>5.5</td>
<td>240</td>
</tr>
<tr>
<td>Heyer et al. (2004)</td>
<td>FCRAO</td>
<td>CO(1–0)</td>
<td>45</td>
<td>1</td>
<td>900</td>
<td>5.5</td>
<td>53</td>
</tr>
<tr>
<td>Enzariola et al. (2003)</td>
<td>BIMA</td>
<td>CO(1–0)</td>
<td>13</td>
<td>2</td>
<td>900</td>
<td>5.5</td>
<td>240</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>HI Survey</th>
<th>Telescope</th>
<th>Line</th>
<th>Angular resolution (&quot;)</th>
<th>Spectral resolution (km s(^{-1}))</th>
<th>Area (arcmin(^2))</th>
<th>(R_{\text{max}}) (kpc)</th>
<th>Sensitivity per channel (mJy/beam)</th>
<th>Sensitivity per channel (K)</th>
</tr>
</thead>
<tbody>
<tr>
<td>This paper</td>
<td>VLA BCD</td>
<td>5–25(^{a})</td>
<td>1.3</td>
<td>9</td>
<td>1.1–2.8(^{a})</td>
<td>24–2.75(^{e})</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Putman et al. (2009)</td>
<td>Arecibo</td>
<td>204</td>
<td>5.15</td>
<td>22</td>
<td>55</td>
<td>0.3</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Corbelli &amp; Schneider (1997)</td>
<td>Arecibo</td>
<td>270(^{b})</td>
<td>4.1</td>
<td>20</td>
<td>...(^{c})</td>
<td>...(^{c})</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Deul &amp; van der Hulst (1987)</td>
<td>WRST+Effelsberg 100m</td>
<td>12 \times 24</td>
<td>8.2</td>
<td>9</td>
<td>...(^{c})</td>
<td>1.2</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

\(^{a}\) See Table 3 for details.
\(^{b}\) Undersampled hexagonal grid, the value corresponds to the grid spacing.
\(^{c}\) The authors give an integrated intensity sensitivity value of 1–2 Jy km s\(^{-1}\) or 5.5–11 K km s\(^{-1}\).

2. Observations and data reduction

2.1. IRAM CO(2–1) observations

The observations presented here are a follow-up to the Gardan et al. (2007) mapping of a large part of M33. All data were taken with the 30 meter telescope run by the Institut de RadioAstronomie Millimétrique (IRAM) on Pico Veleta near Granada, Spain. The observations were carried out starting in August 2008 and also include the data taken by Gardan et al. (2007).

The IRAM HERA instrument (HEterodyne arRAy), composed of 9 dual-polarization 1.3 mm receivers (Schuster et al. 2004), was used in On-The-Fly scanning mode. All data are presented in the main beam temperature scale and we have assumed forward and main beam efficiencies of \(\eta_{\text{for}} = 0.90\) and \(\eta_{\text{mb}} = 0.49 \pm 0.04\) for the HERA observations, the sensitivity is then 9.6 Jy/K (Schuster et al. 2004). The scanning speed was 3 arcseconds per second of time with dumping every second. Both the VESPA backend, at 1.25MHz resolution, and the WILMA correlator with 2MHz channel spacing, were used simultaneously. Reference positions were observed roughly every three minutes and these were taken to be holes in the H\(_{\text{I}}\) column density outside of the M33 disk with no visible FIR emission. The average beam size over the nine pixels is 11.7". The array was tilted by 18.5 degrees, obtaining 7.6 arcsecond spacing between the individual pixel tracks, and then shifted in position by 3.8 arcsec in order to obtain full sampling. This was first done for scanning parallel to RA, then parallel to Dec, the array was then rotated by 90 degrees and the procedure was repeated. The array rotation is to have different pixels cover the same regions to improve data homogeneity and reduce striping. For the HIIF low noise major axis strip, additional scans along and perpendicular to the strip direction where acquired. The repetition of independent scans over the same areas also randomizes any pointing errors. Thus, if there are pointing uncertainties of 2", these contribute negligibly to the overall beam size (through a negligible smearing). The part of M33 presented here comes from the observation of about 5 million spectra.

Reduction was carried out within Gildas\(^{E}\) CLASS and GREG software. The intensity of the CO lines is usually small so that emission is not seen in individual spectra, the baseline removal was therefore implemented using a windowing based on the H\(_{\text{I}}\) data at 17" presented in section 2.2. In a first step a linear baseline is fitted and subtracted from each individual spectrum not taking into account channels inside a window computed by finding the first channels at a 3\(\sigma\) level framing the peak channel of H\(_{\text{I}}\) emission and adding a 40 km s\(^{-1}\) buffer on each side of this. A filtering step is then applied to remove spectra with anomalous noise. The actual noise is computed for each spectrum and compared to the theoretical noise computed from the system temperature, integration time and spectral resolution. The spectra with excess noise attributable to poor baselines are removed; this corresponds to about 10% of the spectra taken. The remaining spectra are then convolved by a gaussian to obtain a regularly gridded cube of the desired resolution. We then fit and remove a third order polynomial to the spectra corresponding to each of the spatial pixels in the gridded cube this time reducing the velocity buffer described previously to 30 km s\(^{-1}\). With this method we have made CO cubes with spacial resolutions of 12, 15 and 25" and with velocity resolutions of 2.6 km s\(^{-1}\).

We compute the CO(2–1) integrated intensity map using a masking method (Gratier et al. 2010) taking into account H\(_{\text{I}}\) data, we developed in order to filter out some of the noise present in the observations and increase the sensitivity to low intensity possibly diffuse CO emission. Previous masking methods used masks created from spatially smoothed versions of the original CO data cubes to filter out regions dominated by noise (Adler et al. 1992; Diegel et al. 1996; Leonard et al. 1999). We use the 21 cm atomic hydrogen data at 17" \times 17" \times 1.27 km s\(^{-1}\).
resolution presented in section 2.2. To achieve the same goal, the underlying hypothesis being that molecular gas is unlikely to be present for low enough values of N(H\textsc{i}) so the corresponding velocity channels can be discarded when computing the integrated intensity CO map. For each pixel of the H\textsc{i} cube, we create a binary mask keeping only the velocity range for each pixel corresponding to a H\textsc{i} signal value above a defined threshold in the H\textsc{i} signal. The integrated moment map for the CO(2–1) data is then computed summing only velocity channels included in the H\textsc{i} mask. Figure 2 shows the integrated intensity map for the 12" cube using this method and Fig. 4 corresponding noise map at 12" computed from 30 channels free of emission from M33 (from −800 to 0 km s\(^{-1}\) \(V_{LSR}\)). The result is an increased S/N ratio as the channels contributing only noise to the sum are no longer taken into account. The value of the H\textsc{i} signal threshold was chosen at 10 K which corresponds to 1.9σ where σ is the rms noise level in the 17" H\textsc{i} cube. We tested masking values between 5 and 25 K (0.9 to 4.6σ) and the total CO intensity varied by only a few percent. Significantly above or below these values, CO signal is lost or more noise is included. At the 10 K brightness level, the H\textsc{i} lines are broad and sample several channels to either side of the rotation curve. Comparing with other techniques of looking for CO, we found no evidence of missing CO emission by this technique. As an example, the “Lonely Cloud” (\(\alpha = 1^\text{h}34^\text{m}16^\text{s}7, \delta = +30^\circ59'3'\), J2000) ([Gardan et al. 2007]) in an interarm region not near an H\textsc{i} maximum is included by this method.

2.2. VLA H\textsc{i} mosaicing

The H\textsc{i} observations are from archival VLA B, C, and D array data taken as part of projects AT206 and AT268 in 1997, 1998 and 2001. At the frequency of the H\textsc{i} line the VLA primary beam is 32". The mosaic comprises 20 D array pointings on a nearly square grid separated by 16' and 6 B and 6 C pointings in a diamond shape covering the entire star forming disk. The primary beam centers of the B/C and D array mosaic are shown in Fig. 4 overlaid on the Spitzer 70\(\mu\)m map – the D array data contains more pointings extending further from the center of M33. Since our goal is compare the distribution and kinematics of the atomic and molecular gas at high resolution, we focus here on the stellar disk where VLA B-, C- and D-array data are available. Data calibration and imaging were carried out using the 2.4 version of the CASA\footnote{http://casa.nrao.edu/} software package. The quasar 3C48 (\(\alpha = 1^\text{h}37^\text{m}41^\text{s}3, \delta = +33^\circ09'35'\)) was used as flux, bandpass and phase calibrator (with an adopted flux of 16.04 Jy). The calibration of the visibilities was done through the following tasks. First the correct flux was applied to the calibrator visibility model using setjy, then frequency bandpass and gain calibrations were computed using the bandpass and gaincal tasks.

The calibration solutions were finally applied to the entire dataset with the use of the applycal task, continuum emission averaged on channels not including line emission from M33 was subtracted in the \(\nu\)-space using the uvcontsub task. Flagging of bad data was done by hand. The imaging was achieved using the multi-scale clean MSCLEAN algorithm described in [Cornwell 2008]. This method has several advantages over the classical clean algorithms that use point like sources to model the observed emission. By cleaning several scales simultaneously, it is able to efficiently recover extended emission and eliminates problems of flux correction ([Jorsater & van Moorsel 1995]), negative bowls surrounding regions of extended emission, and the pedestal of low level uncleansed flux in the final residual map.

An extensive comparison of multiscale clean and its point source counterpart has been carried out in [Rich et al. 2008]. A series of four cubes with angular resolutions of 5, 12, 17, and 25 arc-seconds and 1.27 km s\(^{-1}\) channel width were computed. The 5" cube corresponds to the highest resolution achievable with the B array. The other cubes were obtained using gaussian tapers in the \(\nu\)-plane of respectively 4900 and 3150 and 1900 meters FWHM. In order to minimize sidelobes over the entire map, a robust weighting scheme [Briggs 1995] was applied to the data.

Fig. 2. IRAM CO(2–1) integrated intensity map of the galaxy M33 at a resolution of 12". This map was made using the VLA H\textsc{i} cube (17\arcsec resolution) as a mask, employing a 10 K cutoff per channel (see Sect 2.2). As a consequence, areas without HI emission above that level appear white. Grey ellipses are drawn every kpc using orientation parameters listed in Table 1.
Fig. 3. CO noise map of the 12″ cube computed from 30 channels free of M33 emission, from −80 to 0 km s\(^{-1}\) V\(_{\text{LSR}}\).

using a robustness factor of 0.5. The rms noise for the cubes was computed over the channels without emission from M33 within a galactocentric radius of 8.5 kpc. Table 3 summarizes the beam sizes and noise properties of the different cubes. A zeroth moment map (Fig. 5 with the CO 1 K km s\(^{-1}\) intensity contour superposed) was computed at 17″ resolution by masking out regions of the 17″ cube where emission from the 25″ cube was below 2.75K (1\(\sigma\)). The data presented here will be made available through the Centre de Données de Strasbourg (CDS)\(^d\).

Table 3. Beam and rms noise properties of the H\(_{\text{I}}\) 21 cm datacubes.

<table>
<thead>
<tr>
<th>Beam (″ × ″)</th>
<th>PA (°)</th>
<th>(\sigma_S) (mJy/beam)</th>
<th>(\sigma_T) (K)</th>
</tr>
</thead>
<tbody>
<tr>
<td>5.5 × 5.2</td>
<td>-95.1</td>
<td>1.1</td>
<td>24</td>
</tr>
<tr>
<td>12.0 × 11.6</td>
<td>-31.8</td>
<td>2.0</td>
<td>9</td>
</tr>
<tr>
<td>17.2 × 17.1</td>
<td>-45.8</td>
<td>2.5</td>
<td>5.4</td>
</tr>
<tr>
<td>25.9 × 24.2</td>
<td>-74.8</td>
<td>2.8</td>
<td>2.75</td>
</tr>
</tbody>
</table>

The rms noise, in mJy/beam and brightness temperature, was calculated over an ellipse of galactocentric radius 8.5 kpc after primary beam correction.

3. Results

3.1. Gas masses

In this section we present measurements of total integrated flux and derived total molecular and atomic gas masses.

3.1.1. Molecular Gas

The CO(2–1) luminosity measured over our observed area is \(1.7 \pm 0.1 \times 10^7\) K km s\(^{-1}\)pc\(^2\). Our CO mapped region covers 28% of the total galaxy area up to 8.5 kpc but 61 ± 2% of the infrared emission (see details in Sect. 3.3). The estimated CO(2–1) luminosity for the whole R\(_{25}\) disk is thus \(2.8 \pm 0.3 \times 10^7\) K km s\(^{-1}\)pc\(^2\). The molecular gas
Table 4. Comparison of results with previous studies of M33.

<table>
<thead>
<tr>
<th>CO surveys</th>
<th>Molecular gas mass(^b) ((M_\odot))</th>
<th>SFR(^g) ((M_\odot\ yr^{-1}))</th>
<th>Depletion time(^c) ((10^5\ yr))</th>
<th>Depletion time(^d) ((10^8\ yr))</th>
</tr>
</thead>
<tbody>
<tr>
<td>This study</td>
<td>(1.9 \times 10^8) ((3.3 \times 10^8)^e)</td>
<td>0.27(^h)</td>
<td>7</td>
<td>7</td>
</tr>
<tr>
<td>Gardan et al. (2007)</td>
<td>(3.2 \times 10^7)</td>
<td>0.05–0.13(^i)</td>
<td>2–6</td>
<td>1–3</td>
</tr>
<tr>
<td>Rosolowsky et al. (2007)</td>
<td>(3.6 \times 10^8)</td>
<td>0.45(^i)</td>
<td>8</td>
<td>…</td>
</tr>
<tr>
<td>Hever et al. (2004)</td>
<td>(3.5 \times 10^{8e})</td>
<td>0.45(^i)</td>
<td>8</td>
<td>1–3</td>
</tr>
<tr>
<td>Ensariola et al. (2003)</td>
<td>(9 \times 10^7)</td>
<td>0.24(^i)</td>
<td>4</td>
<td>2</td>
</tr>
<tr>
<td>Cortelli (2003)(^a)</td>
<td>(2.9 \times 10^8)</td>
<td>0.45(^i)</td>
<td>6</td>
<td>…</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>HI surveys</th>
<th>Atomic gas mass(^m) ((M_\odot))</th>
</tr>
</thead>
<tbody>
<tr>
<td>This study</td>
<td>(1.4 \times 10^9)</td>
</tr>
<tr>
<td>Putman et al. (2009)</td>
<td>(1.5 \times 10^9)</td>
</tr>
<tr>
<td>Cortelli &amp; Schneider (1997)</td>
<td>(2.6 \times 10^9)</td>
</tr>
<tr>
<td>Deul &amp; van der Hulst (1987)</td>
<td>(1.1 \times 10^9)</td>
</tr>
</tbody>
</table>

\(\text{CO (1–0) data from} \) 
\(\text{H}_1 \text{ data from} \)
\(\text{FCRAO CO(1–0) data from} \)
\(\text{FCRAO CO(1–0) data from} \)
\(\text{FCRAO CO(1–0) data from} \)
\(\text{FCRAO CO(1–0) data from} \)
\(\text{FCRAO CO(1–0) data from} \)
\(\text{FCRAO CO(1–0) data from} \)

\(\text{CO surveys} \): This study, Gardan et al. (2007), Rosolowsky et al. (2007), Hever et al. (2004), Ensariola et al. (2003), Cortelli (2003).

\(\text{HI surveys} \): This study, Putman et al. (2009), Cortelli & Schneider (1997), Deul & van der Hulst (1987).

\(a\) FCRAO CO(1–0) data from Hever et al. (2004).

\(b\) All values converted to our adopted \(X_{\text{CO}} = 4 \times 10^{20} \text{cm}^{-2} / (\text{K km s}^{-1})\) (twice Galactic).

\(c\) The value in parentheses is the extrapolation to the whole galaxy for comparison with other surveys. See Sect. 3.1.1 for details.

\(d\) Original value used \(X_{\text{CO}} = 2 \times 10^{20} \text{cm}^{-2} / (\text{K km s}^{-1})\) (Galactic).

\(e\) Original value used \(X_{\text{CO}} = 3 \times 10^{20} \text{cm}^{-2} / (\text{K km s}^{-1})\).

\(f\) Original value used \(X_{\text{CO}} = 2.8 \times 10^{20} \text{cm}^{-2} / (\text{K km s}^{-1})\).

\(g\) Over the areas mapped in CO for each survey.

\(h\) Value computed by dividing the first by the second column of this table.

\(i\) Original value given by the cited papers.

\(j\) All values converted to our adopted distance \(D = 840 \text{kpc}\).

\(k\) Original value used a distance of 730 kpc. Inside a 3σ contour.

\(l\) Original value used a distance of 690 kpc. Over their whole mapped area extending beyond 8.5 kpc.

\(m\) All values converted to our adopted distance \(D = 840 \text{kpc}\).

\(n\) Original value used a distance of 730 kpc. Inside a 3σ contour. Depletion time is derived from the \(\text{H}_1\) integrated intensity using a conversion factor 1.8 \(\times 10^{18} \text{cm}^{-2} / (\text{K km s}^{-1})\) (Rohlfs & Wilson 1996). The total \(\text{H}_1\) mass over the same region is \(1.4 \times 10^9 M_\odot\), Putman et al. (2009) using the Arecibo telescope find give 1.5 \(\times 10^9 M_\odot\) inside a very similar contour which they define as the star forming disk.

3.2. Kinematics

We obtain velocity fields for both datasets by computing the first moment of the emission. In the case of \(\text{H}_1\), the 17’’ data was masked by the 25’’ cube with a threshold of 10K. For the CO data, we computed the first moment of the original non \(\text{H}_1\) masked datacube, masking out emission below 4σ. Both methods are thus completely independent of each other. Fig. 6 shows the CO and \(\text{H}_1\) velocity fields at respectively 12’’ and 17’’. The excellent agreement between the two velocity fields is a further argument in favor of using the \(\text{H}_1\) masking to obtain CO integrated intensity maps.

3.3. Radial distribution

The radial distribution of properties were derived taking azimuthal averages over elliptical rings with inclinations of i=56° and position angle PA=22.5° (measured toward East from North). The quantities are therefore deprojected from the sky plane onto the plane of M33. The radial step is 500pc which corresponds to about 2’. The bottom part of Fig. 7 shows the
cumulative fractions of the area (solid line) and 70 μm surface brightness (dashed line) in our CO map (black area in the inset) compared to the whole elliptical region of radius less than R (sum of grey and black areas of the inset). For galactocentric distances below 2 kpc our map samples more than 80% of the accessible disk at these radii both in area and 70 μm flux. The whole CO mapped area corresponds to 28% of the area up to 8.5 kpc but more than 60% of the corresponding 70 μm flux, this is explained by the rapid decrease of infrared emission with radius. The corresponding values for 24 and 160 μm are similar to a few percent at each radius. Tab. 5 summarizes the source and resolution of the ancillary data used.
3.3.1. Surface brightness

The top part of Fig. 7 presents the radial variations in average surface brightness of the CO and infrared data from Spitzer at 24 µm, 70 µm and 160 µm, the solid lines correspond to averages taking into account only our CO mapped area and the dashed lines to the whole elliptical annuli for a given radius. For each tracer, the similarity between these two averages is an indication that the area mapped in CO is representative of the M33 disk in its entirety. The 4 datasets have similar exponential decreases with radius i.e. $F(r) \propto e^{-r/L}$ with a corresponding scale length $L$. Table 6 summarizes the scale lengths computed by least square fitting for CO and IR data, and Table 7 scalelength values from surveys presented in Table 2. The infrared data scale lengths are in agreement with the values in Tabatabaei et al. (2007a) and Verley et al. (2009) albeit slightly smaller as the IR tracers seem to drop marginally more rapidly in our CO mapped region (solid lines in Fig. 7) compared to the whole galaxy (dashed lines in the same figure). The scale length of CO is higher than the ones of 24 µm and 70 µm and similar to the 160 µm.

3.3.2. Mass surface densities

We compute mass surface densities for atomic and molecular gas and study their radial distribution. For the molecular gas mass, we use the same hypotheses as in Sect. 3.1.1 including a constant CO(2–1)/CO(1–0) factor. Taking a linearly decreasing value of this factor would modify only slightly the radial distribution, raising the most extreme outer points by a factor 1.5 and lowering the inner ones by a factor 1.1, at the price of added uncertainties that cannot be simply estimated. The top panel of Fig. 8 shows the mass surface density of atomic and molecular gas as a function of radius. The atomic gas surface density is constant at an average value of $6 \pm 2 M_\odot$ pc$^{-2}$, the molecular gas mass surface density follows the same trend as the CO surface brightness presented in Fig. 7. The fraction of molecular to atomic gas in...
Table 5. Ancillary data used in this paper.

<table>
<thead>
<tr>
<th>Wavelength</th>
<th>Resolution</th>
<th>Telescope</th>
</tr>
</thead>
<tbody>
<tr>
<td>160 µm</td>
<td>40″</td>
<td>Spitzer MIPS a</td>
</tr>
<tr>
<td>70 µm</td>
<td>18″</td>
<td>Spitzer MIPS a</td>
</tr>
<tr>
<td>24 µm</td>
<td>6″</td>
<td>Spitzer MIPS a</td>
</tr>
<tr>
<td>8 µm</td>
<td>2″</td>
<td>Spitzer IRAC b</td>
</tr>
<tr>
<td>6570 Å (Hα)</td>
<td>2″ (pixel size)</td>
<td>KPNO c</td>
</tr>
<tr>
<td>1500 Å</td>
<td>5″</td>
<td>GALEX FUV d</td>
</tr>
</tbody>
</table>


Table 6. Exponential scale length L in kpc for Spitzer MIPS and CO emission computed over the area mapped in CO.

<table>
<thead>
<tr>
<th>Wavelength</th>
<th>0.5-3.5 kpc</th>
<th>3.5-7 kpc</th>
<th>0.5-7 kpc</th>
</tr>
</thead>
<tbody>
<tr>
<td>24 µm</td>
<td>1.36 ± 0.05</td>
<td>1.02 ± 0.07</td>
<td>1.40 ± 0.03</td>
</tr>
<tr>
<td>70 µm</td>
<td>1.51 ± 0.07</td>
<td>1.05 ± 0.06</td>
<td>1.48 ± 0.04</td>
</tr>
<tr>
<td>160 µm</td>
<td>2.26 ± 0.12</td>
<td>1.30 ± 0.06</td>
<td>1.83 ± 0.07</td>
</tr>
<tr>
<td>CO</td>
<td>2.0 ± 0.2</td>
<td>1.6 ± 0.3</td>
<td>1.9 ± 0.1</td>
</tr>
</tbody>
</table>

Blow ups of these 5 figures are available in the online version of this article.

3.3.3. Star Formation Efficiency

The star formation efficiency is usually defined as the ratio of the star formation rate over the mass of molecular gas available to form stars.

\[ \text{SFE} = \frac{\text{SFR}}{M_{\text{H}_2}} \]

Defined in this way, the SFE is the inverse of the molecular gas depletion time. Estimates from a variety of observations (Blitz & Rosolowsky 2006; Kennicutt 1998; Hippelein et al. 2003) give a value of the SFR which varies from 0.3 to 0.7 M⊙yr⁻¹ for the entire disk. A recent study by Verlev et al. (2009) using extinction corrected Hα, FUV, and infrared tracers gives a value of 0.45 ± 0.1 M⊙yr⁻¹ for total star formation rate. This is our adopted value.

Our CO map does not cover the whole disk but covers respectively 63%, 61% and 59% of the IR emission at 24 µm, 70 µm and 160 µm (see bottom panel of Fig. 7) so the SFR is ~ 0.27 M⊙ yr⁻¹ within the area covered by our CO map. With a N(H2)/I_{CO(1-0)} value of 4×10⁸ cm⁻²/(K km s⁻¹) and an average 2–1/1–0 line ratio of 0.73 (see Sect. 3.4.1 for details), the molecular gas mass within the CO map is 1.9×10⁹ M⊙. With these values, the SFE is 1.4×10⁻³ yr⁻¹ or a molecular gas depletion time 7×10⁹ yr. An SFE of 1.6×10⁻³ yr⁻¹ is a factor ~ 3 higher than that found in large spiral galaxies (sample in, Kennicutt 1998, Murgia et al. 2002, Leroy et al. 2008). Earlier studies used lower N(H2)/I_{CO(1-0)} values and thus obtained even higher SFEs (or shorter H2 depletion times) as is shown in Table 4.

In order to bring the SFE of M33 to the same level as in large spirals, the N(H2)/I_{CO(1-0)} factor would have to be about 3 times the value we use whereas previous studies used even smaller values (Wilson & Scoville 1990, Corbelli 2003, Heyer et al. 2004, Engargiola et al. 2003, Rosolowsky et al. 2007, Gardan et al. 2007). Through Virial mass estimates and 13/12 CO line ratios, Braine et al. (submitted to A&A) support a value of about N(H2)/I_{CO(1-0)} = 4×10⁸ cm⁻²/(K km s⁻¹) in the disk of M33, possibly closer to the Galactic value in the inner disk.

Figure 8 presents the radial distributions of the SFE as calculated from the Hα, 24, and 70 µm emission. The units are arbitrary as we are interested in showing the radial trend. The SFR were taken as being directly proportional to the Hα luminosities. Hα and FIR emission tend to show opposite biases – the Hα suffers extinction where the FIR is strong and the FIR is weak where little dust is present but the Hα is unaffected. If one were to compute the star formation efficiency taking into account the total gas mass, the SFE would drop dramatically with increasing radius as opposed to the SFE computed using only molecular gas that does not appear to vary with radius. Table 4 summarizes the masses, star formation rates and depletion times for this study and previous surveys of M33.

3.4. The ISM and star formation tracers

In this section, we compare the CO emission with tracers of star formation. We present a series of figures (10 to 13) comparing the CO emission, as contours, overlaid on images of the 8, 24 and 70 µm FIR emission as well as the Hα and FUV emission (Blow ups of these 5 figures are available in the online version of this article). The 160 µm image is at lower resolution and brings little morphological information not present in the 70 µm image.
Fig. 9. Radial distribution of the Star Formation Efficiency in arbitrary units, all of the curves are normalized to the central value and shifted by 1 decade along the vertical axis.

Table 7. Comparison with previous surveys of CO emission exponential scale lengths \( L \) in kpc.

<table>
<thead>
<tr>
<th>Study</th>
<th>Range (kpc)</th>
<th>Scalelength (kpc)</th>
</tr>
</thead>
<tbody>
<tr>
<td>This paper</td>
<td>0–7</td>
<td>1.9 ± 0.1</td>
</tr>
<tr>
<td>Gardan et al. (2007)</td>
<td>2–6</td>
<td>1.4 ± 0.1</td>
</tr>
<tr>
<td>Engargiola et al. (2003)</td>
<td>0–7</td>
<td>1.4 ± 0.1</td>
</tr>
<tr>
<td>Corbelli (2003)</td>
<td>0–6</td>
<td>2.5</td>
</tr>
</tbody>
</table>

While the dynamic range in the CO image is much lower than in the images of star formation tracers, the correspondence is globally excellent. Figure 10 shows the scatter plot between the CO integrated intensity and 70 \( \mu \)m emission.

Figures 10–14 show that in almost all cases the CO emission is found where star formation is detected and in general the CO emission is stronger as the level of SF increases. The weaker CO contours outline the regions with a remarkable precision (see online figures 1 to 12 for the blowups). A partial exception is found near \((\alpha = 1^h 34^m 09^s, \delta = +30^\circ 49' 06'', J2000)\) where the CO emission is extremely strong while the SF tracers show rather weak emission. The CO spectrum in the direction of that particular cloud is shown in Fig. 17.

A true exception to the FIR–HI–CO correlation is the inter-arm cloud at \((\alpha = 1^h 34^m 16^s, \delta = +30^\circ 59' 3'\), J2000\) dubbed “Lonely Cloud” by Gardan et al. (2007). It is found in a region with very little if any star formation and relatively weak H\( \alpha \) emission, despite fairly strong CO (and even \(^{13}\)CO) emission. When H\( \alpha \) brightness thresholds higher than 10 K are used to create the masked zero-moment CO maps, the Lonely Cloud disappears. While this cloud was discovered in the first set of observations, it remains a unique object as the observations of the central and southern regions, despite considerably lower noise levels than in Gardan et al. (2007), have not revealed similar molecular clouds.

In the outer parts, the FIR emission is typically quite inconspicuous, even where CO has been detected, but the H\( \alpha \) emission is often quite strong. In the central regions, the alignment of the emission at the different wavelengths is nearly perfect, especially between CO and FIR. Further out, however, notable offsets become more common although there is a clear association between the CO and FIR/H\( \alpha \) emitting regions. The offsets are less than 100 pc, particularly in comparison with the FIR.

The CO emission follows HI-defined arm-like structures very closely, completely avoiding the H\( I \) holes. As can be seen in Fig. 15 and corresponding blow ups in the online version of the paper, the overall correlation between H\( I \) and CO emission is very good. CO is generally found in HI-rich regions although not all of the H\( I \) peaks show CO emission and not all CO emission falls on H\( I \). It is not clear what the source of the 70 \( \mu \)m emission is in for example the H\( I \) (and CO) hole around \((\alpha = 1^h 34^m 25^s, \delta = +30^\circ 49' 06'', J2000)\).
Fig. 11. KPNO Hα image with $I_\text{CO}$ main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s$^{-1}$. The beam size is shown as a white dot in the lower left corner. Blow ups are available in the online version.

$\delta = +30^\circ35'00'', J2000$) as the amount of neutral gas is very low. The emission in this region is lower in the 24 and 8 $\mu$m wavebands, so perhaps the 70 micron emission appears stronger than in reality due to the poorer spatial resolution. The upcoming Herschel HERM33ES observations should answer this.

3.5. Examples of interesting spectra

The spectra shown in Figures 17, 18 and 19 are just a few among a huge number but are intended to show “special” positions. Fig 17 shows a very strong CO profile towards a region without strong star formation (note good agreement with Engargiola et al. 2003). The line is very strong with a line width of about 11 km s$^{-1}$ at half-intensity. The other spectra shown are interesting because the CO lines are detected very far from the center of M33, despite the low metallicity and surface mass density. Fig. 18 shows the second molecular cloud detected at or beyond $R_{25}$ in M33. A secondary peak offset in velocity by about 8 km s$^{-1}$ may be present. Fig 19 shows further spectra near $R_{25}$ which clearly is composed of emission from two clouds, the cloud to the SE (of the figure) at -260 km s$^{-1}$ and the one to the NW at about -273 km s$^{-1}$. The spectral resolution of the data is barely sufficient to measure the linewidths of the clouds, which are typically less than 5 km s$^{-1}$ at half power. This was also true of the spectra of distant clouds shown in Braine et al. (2010) — they are much narrower than in the H$_2$-rich inner disk. It is not clear whether the broad linewidths in the inner disk are because the GMCs themselves are much more massive or due to a superposition of clouds (or a single GMC in a region with significant diffuse emission).

The spectra shown also illustrate the dynamic range of the data acquired, roughly 100 in integrated intensity.
4. Conclusions

This work presents high-resolution maps of the atomic and molecular gas in the disk of M33 via observations of the H\textsc{i} 21 cm and CO(2–1) lines. The whole disk out to 8.5 kpc is covered in the H\textsc{i} line and about 60% of the emission in CO.

Assuming the N(H\textsubscript{2})/I\textsubscript{CO(1–0)} factor to be twice that of the Galaxy, because of the sub-solar metallicity of M33, we estimate a molecular gas mass of $3.3 \times 10^8$ M\textsubscript{☉} roughly 20% of the $1.4 \times 10^9$ M\textsubscript{☉} detected in the inner 8.5 kpc in H\textsc{i}. Azimuthally averaging, the H\textsc{i} surface density is close to constant with radius but the H\textsubscript{2} decreases exponentially with a scale length of 1.9kpc. The H\textsubscript{2}/H\textsc{i} mass ratio decreases from about unity to 1%.

The correspondence between the peaks and holes in the distributions of molecular and atomic gas is excellent and follows the peaks and troughs in the FIR, MIR, and H\textalpha{} images. The SFE is approximately constant with radius, suggesting that molecular gas is transformed into stars at a similar rate (assuming a similar IMF) at all galactocentric radii. However, the SFE in the small, gas-rich, low-metallicity, blue spiral M33 appears 2-4 times higher than what is observed in large nearby spirals.

The sensitivity of the survey is such that CO emission is detected far out in the disk of M33 although few clouds are present and the lines are much weaker in intensity.

Acknowledgements. We thank the IRAM staff in Granada for their help with the observations.
Fig. 15. VLA H\textsubscript{i} zeroth moment image with I\textsubscript{CO} main beam contours of 1 (green), 2 (red), 4 (black), 8 (blue) K km s\textsuperscript{-1}. The beam size is shown as a white dot in the lower left corner. Note how the CO emission follows the bright H\textsubscript{i} and outlines the H\textsubscript{i} holes. Blow ups are available in the online version.

Fig. 16. Scatter plot between CO integrated intensity and 70 µm emission.

Fig. 17. Spectrum of a bright cloud at position (\(\alpha = 1\text{h}34\text{m}09.4\text{s}, \delta = +30^\circ49'06''\), J2000), unlike other strong peak of CO integrated intensity, it lacks strong emission in star formation tracers (see Sect. 3.4).
Fig. 18. Spectra of the outer disk cloud ($\alpha = 1^{h}34^{m}50^{s}.8$, $\delta = +31^\circ08'26''$, J2000) shown with 6'' spacing between panels. If the secondary peak at -266 km s$^{-1}$, seen in the second spectrum from the bottom in the middle column, is real then the cloud is small as it is only seen within a single beam.

Fig. 19. Spectra of two more outer disk clouds at different velocities. The lower left panel is at ($\alpha = 1^{h}34^{m}21^{s}.7$, $\delta = +31^\circ04'13''$, J2000) and the panel spacing is 6''.
References

Israel, F. 2000, in Molecular Hydrogen in Space, ed. F. Combes & G. Pineau Des Forêts, 293+–
Paturel, G., Petit, C., Prugniel, P., et al. 2003, VizieR Online Data Catalog, 7237, 0

Appendix A: Online Figures
Fig. A.1. Northern part of KPNO Hα image with I_{CO} main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s^{-1}. The beam size is shown as a white dot in the lower left corner.
Fig. A.2. Center part of KPNO Hα image with I_{CO} main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s^{-1}. The beam size is shown as a white dot in the lower left corner.
Fig. A.3. Southern part of KPNO Hα image with $I_{\text{CO}}$ main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s$^{-1}$. The beam size is shown as a white dot in the lower left corner.
Fig. A.4. Northern part of Spitzer 8 μm image with I_{CO} main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s^{-1}. The beam size is shown as a white dot in the lower left corner.
Fig. A.5. Center part of Spitzer 8 µm image with I$_{CO}$ main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s$^{-1}$. The beam size is shown as a white dot in the lower left corner.
Fig. A.6. Southern part of Spitzer 8 μm image with I_{CO} main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s^{-1}. The beam size is shown as a white dot in the lower left corner.
Fig. A.7. Northern part of Spitzer 24 µm image with I_{CO} main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s^{-1}. The beam size is shown as a white dot in the lower left corner.
Fig. A.8. Center part of Spitzer 24 μm image with $I_{CO}$ main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s$^{-1}$. The beam size is shown as a white dot in the lower left corner.
Fig. A.10. Northern part of Spitzer 70 μm image with I_{CO} main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s\(^{-1}\). The beam size is shown as a white dot in the lower left corner.
Fig. A.11. Center part of Spitzer 70 µm image with $I_{\text{CO}}$ main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s$^{-1}$. The beam size is shown as a white dot in the lower left corner.
Fig. A.12. Southern part of Spitzer 70 µm image with $I_{\text{CO}}$ main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s$^{-1}$. The beam size is shown as a white dot in the lower left corner.
Fig. A.13. Northern part of VLA H I 21 cm integrated intensity image with $I_{CO}$ main beam contours of 1 (green), 2 (red), 4 (black), 8 (blue) K km s$^{-1}$. The beam size is shown as a white dot in the lower left corner.
Fig. A.14. Center part of VLA H\textsc{i} 21 cm integrated intensity image with $I_{\text{CO}}$ main beam contours of 1 (green), 2 (red), 4 (black), 8 (blue) K km s$^{-1}$. The beam size is shown as a white dot in the lower left corner.
Fig. A.15. Southern part of VLA H I 21 cm integrated intensity image with I_{CO} main beam contours of 1 (green), 2 (red), 4 (black), 8 (blue) K km s^{-1}. The beam size is shown as a white dot in the lower left corner.
Fig. 9. Southern part of Spitzer 24 μm image with I_{CO} main beam contours of 1 (white), 2, 4 (black), 8 (blue) K km s^{-1}. The beam size is shown as a white dot in the lower left corner.